



CERN



Dilia María PORTILLO on behalf of the ATLAS Collaboration

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Beyond the Standard Model (BSM) physics

•The Standard Model (SM) has been experimentally validated and proven to be impressively precise. However, it cannot explain some of the important questions in physics:



Several BSM models extends the SM by introducing new symmetries and particles/fields.

Some of the most popular in BSM searches at LHC:

 Minimal Supersymmetric Standard Model (MSSM) solution to "hierarchy problem" and has DM candidate
 Two Higgs Doublet Models (2HDM) extend beyond the SM Higgs sector to include two complex Higgs Doublets (subset of MSSM, could introduce baryon asymmetry or a DM candidate or neutrino masses)
 Heavy Vector Triplet model: Simplified model used to capture the phenomenology of a new resonance
 Radion model that includes Kaluza-Klein gauge bosons

Effective Field Theory (EFT): Contact interactions for a more model-independent description
 Simplified Models for DM: Usually adding a massive mediator (e.g. Z')

LHC operation and ATLAS detector



SM is successful for particle collisions

■Discrepancies may indicate new physics = new particles/fields



Limits for BSM searches

ATL-PHYS-PUB-2022-034



†Small-radius (large-radius) jets are denoted by the letter j (J).

~100 decay channels studied for various models that predict certain production rate (extra dimensions, gauge bosons, contact interactions, dark matter, heavy quarks, excited fermions, leptoquarks, etc.)
 Commonly excluded masses ~ 0.4 – 12 TeV

Limits for SUSY searches here: <u>ATL-PHYS-PUB-2022-013</u> (and in bonus slides)

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Some of the latest results

Using full Run 2 date-set!

ArXiv 2207.00230	July 2022
ATLAS-CONF-2022-028	May, 2022
ATL-PHYS-PUB-2022-04	Sep. 2022 3
}- ArXiv 2211.02617	Nov. 2022
s <u>ArXiv. 2211.08945</u>	Nov. 2022
d <u>ArXiv 2211.08028</u>	Nov. 2022
ATLAS-CONF-2022-066	Nov. 2022
ArXiv 2207.03925	Sep 2022
<u>ArXiv 2210.05415</u>	July, 2022
ATLAS-CONF-2022-039	June, 2022
	1 0000
	ArXiv 2207.00230 ATLAS-CONF-2022-028 ATL-PHYS-PUB-2022-043 ATL-PHYS-PUB-2022-043 ArXiv 2211.02617 S ArXiv. 2211.08945 ArXiv 2211.08945 ArXiv 2211.08028 ATLAS-CONF-2022-066 ArXiv 2207.03925 ArXiv 2210.05415 ATLAS-CONF-2022-039

More papers and conference notes to be found in ATLAS publications page for Exotics: <u>EXOT</u>, Higgs/Diboson searches: <u>HDBS</u>, and Supersymmetry: <u>SUSY</u>

$A/V' \rightarrow Zh(\rightarrow bb)$: Search for heavy resonances decaying into a Z or W and a Higgs boson in final states with leptons ArXiv 2207.00230

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•Search for new resonances decaying into Z or W and a SM Higgs boson • $\nu\nu b\bar{b}$, $\ell\ell b\bar{b}$, $\ell\nu b\bar{b}$ final states **bbA**

Theory Models

Heavy CP-odd scalar boson A from a generic two Higgs Doublet Model (<u>2HDM</u>)
 Z' or W' from Heavy Vector Triplet (HVT)

Selection

•Events with exactly 0, 1 or 2 charged leptons are selected with $E_{\rm T}^{\rm miss}$ or a combination of single-lepton triggers



 Higgs (bb) is reconstructed either as two small-R jets (resolved category) or as a single large-R jet (merged category)



$A/V' \rightarrow Zh(\rightarrow bb)$: Search for heavy resonances decaying into a Z or W and a Higgs boson in final states with leptons ArXiv 2207.00230



 Largest deviation is found in the Z' and ggA searches for a mass of 500 GeV with a local (global) significance of 2.1 σ (1.1 σ)

Results: Exclusion in HVT parameter space



Results: Exclusion in 2HDM parameter space



Combination of searches for heavy resonances

ATLAS-CONF-2022-028

Direct searches for new heavy resonances are a staple of the LHC physics program
 Combine 13 ATLAS publications during 2018 - 2022!

Final states:

- VV: $WZ \rightarrow qqqq$, vvqq, lvqq, llqq, lvll, $WW \rightarrow qqqq$,
- •VH: $WH \rightarrow \underline{qqbb}$, $\underline{\ell vbb}$, $\underline{ZH} \rightarrow vvbb$, $\underline{\ell \ell bb}$
- •Leptonic \underline{ll} , \underline{lv} , $\underline{\tau v}$

Analyses generally search for narrow-width resonances in the final-state mass distribution

Theory Models

Spin-1 mass-degenerate Heavy Vector Triplet model (HVT) [arXiv: 1402.4431]

•Leads to W'±, and Z', collectively denoted V'

• Model A: Weakly-coupled model, couplings gH = -0.56, gf = -0.55

• model B: Strongly-coupled model, couplings gH = -2.9, gf = 0.14

Selection

Analysis	leptons	$E_{T_{miss}}$	jets	b-tags	Discr.
$WW/WZ \rightarrow qqqq$	0	Veto	≥2J	-	m_{VV}
$WZ \rightarrow \nu \nu q q$	0	Yes	≥1J	0	m_{VV}
$WZ \rightarrow \ell \nu q q$	1e, 1µ	Yes	≥2j, ≥1J	0, 1, 2	m_{VV}
$WZ \rightarrow \ell \ell q q$	2e, 2µ	-	≥2j, ≥1J	0	m_{VV}
$WZ \to \ell \nu \ell \ell$	$3 \subset (e, \mu)$	Yes	-	0	m_{VV}
$WH \rightarrow qqbb$	0	Veto	≥2J	1, 2	m_{VH}
$ZH \rightarrow \nu \nu bb$	0	Yes	≥2j, ≥1J	1, 2	m_{VH}
$WH \rightarrow \ell \nu bb$	1e, 1µ	Yes	≥2j, ≥1J	1, 2	m_{VH}
$ZH \rightarrow \ell\ell bb$	2e, 2µ	Veto	≥2j, ≥1J	1, 2	m_{VH}
ℓv	1e, 1µ	Yes	-	-	m_T
τν	1τ	Yes	-	-	m_T
ll	$\geq 2e, \geq 2\mu$	-	-	-	$m_{\ell\ell}$



Combination of searches for heavy resonances

ATLAS-CONF-2022-028



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Summary for beyond Standard Model Higgs boson benchmarks





Summary for beyond Standard Model Higgs boson benchmarks

Benchmark scenario of type-I 2HDM

• assumptions:
$$m_H = m_A = m_{H^{\pm}}$$
, $m_h = 125 \text{ GeV}$, $m_{12}^2 = m_A^2 \sin \beta \cos \beta$

Assume a narrow width (< 5%) Higgs boson.</p>



 $\cos(\beta - \alpha) = -0.1$

 $\cos(\beta - \alpha) = +0.1$

Search for new phenomena in multi-body invariant masses in events with at least one lepton and two jets ArXiv. 2211.08945



Theory Models

Sequential Standard Model, SSM

Includes new heavy gauge bosons W' and Z'
 Simplified dark matter model with a Z' mediator

Radion model that includes Kaluza-Klein gauge bosons coupling to a radion that decays to gluons
 Composite lepton model contains two composite SU(2) fermion doublets that mix to create a new vector boson (Z') and lepton (E)

Strategy

Search for an observable excess of events in multibody invariant masses $(m_{jj\ell}, m_{jj\ell\ell})$ on a smooth, decreasing fitted background.



SSM

DM model

 \mathcal{X}'

W/Z

 v_e/v_μ or e/μ

e/µ

 v_e/v_u

Radion model

Composite lepton model

 \mathcal{N}

мкк 2000 ж.К. 2000 ж.К. 2000 ж.К. 2000 ж.К. 2000 ж.К. 2000 ж.К.

Z/h

e/µ

 v_e/v_μ

Search for new phenomena in multi-body invariant masses in events with at least one lepton and two jets ArXiv. 2211.08945

Results: Limits





WH(→WW): Search for heavy Higgs bosons decaying into vectorbosons in same-sign two-lepton final statesArXiv 2211.02617

Results: Exclusion contours



Search for SUSY in final states with missing transverse momentum

and three or more *b*-jets

ArXiv 2211.08028

Theory Model

- Search for stops, sbottoms and neutralinos in pair-produced gluino events
- •Lightest Stable Particle: χ_1^{0}

•Gtt/Gbb: Simplified models with off-shell stops and sbottoms

Strategy

- Targeting final states with large amount of MET and many b-jets
- Neural Network (NN) trained on four-vector inputs
 - NN defines the probability of an event to be a Gtt or a Gbb
 - Maximise model-dependent sensitivity to models
- •Gbb: 4 Signal Regions, 4 Control Regions (CRs) to constraint ttbar & 4 CR for Z+Jets

Gtt: 4 Signal Regions, 4 CRs to constraint ttbar



Control Regions (CRs)



Signal Regions (SRs)

Search for SUSY in final states with missing transverse momentum and three or more *b*-jets Image: ArXiv 2211.08028 Image: Image:



 Gluinos with masses below 2.44 TeV (2.35 TeV) are excluded at 95% CL for massless neutralinos in the Gtt (Gbb) models.

•These limits represent a substantial increase in over previous ATLAS analysis!

• A cut-and count analysis was also performed considering the processes where both gluinos can have different decays in a event (Gtb) in combination of Gbb and Gtt processes. Exclusion limits for gluino branching ratios to $b\bar{b}_{\chi_1^0}^0$, $t\bar{t}_{\chi_1^0}^0$ and $t\bar{b}_{\chi_1^-}^-$, $t\bar{b}_{\chi_1^-}^+$.

Conclusions

ATLAS continues exploring the energy frontier: Extensive program for searches for new heavy particles!

•No evidence yet but no shortage of models predicting exotic heavy particles

Strong constraints have been placed on the production of such new heavy particles and on BSM parameter space.

More results with full Run 2 data are still being released.

Run 3 already started!

Opportunities at Run 3:

- ■13 TeV \rightarrow 13.6 TeV CM energy
- Increase in luminosity (x 2, statistics!)Cutting edge analysis techniques



Thank you!

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Limits for SUSY searches

• <u>ATL-PHYS-PUB-2022-013</u>

ATLAS SUSY Searches* - 95% CL Lower Limits March 2022

ATLAS SUSY Searches* - 95% CL Lower Limits March 2022							ATLAS Preliminary $\sqrt{s} = 13$ TeV	
	Model	Sig	nature	∫ <i>L dt</i> [1	fb ⁻¹]	Mass limit		Reference
Si	$\tilde{q}\tilde{q},\tilde{q}{\rightarrow}q\tilde{\chi}^0_1$	0 e,µ mono-jet	2-6 jets 1 1-3 jets 1	T_T^{miss} 139 T_T^{miss} 139		[1x, 8x Degen.] 1.0 [8x Degen.] 0.9	1.85 m($\tilde{\chi}_1^0$)<400 GeV m(\tilde{q})-m($\tilde{\chi}_1^0$)=5 GeV	2010.14293 2102.10874
Inclusive Searche	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q \bar{q} \tilde{\chi}_1^0$	0 <i>e</i> , <i>µ</i>	2-6 jets 1	T_T^{miss} 139		Forbidden	2.3 m(t̃ ⁰ ₁)=0 GeV 1.15-1.95 m(t̃ ⁰ ₁)=1000 GeV	2010.14293 2010.14293
	$\begin{array}{l} \widetilde{g}\widetilde{g}, \ \widetilde{g} ightarrow q \overline{q} W \widetilde{\chi}_1^0 \\ \widetilde{g}\widetilde{g}, \ \widetilde{g} ightarrow q \overline{q} (\ell \ell) \widetilde{\chi}_1^0 \\ \widetilde{g}\widetilde{g}, \ \widetilde{g} ightarrow q W Z \widetilde{\chi}_1^0 \end{array}$	1 e,μ 5 ee,μμ 0 e,μ 7	2-6 jets 2 jets 1 7-11 jets 1	139 T 139 T 139 T 139			$\begin{array}{ccc} \textbf{2.2} & m(\tilde{\chi}_1^0) < 600 \ \text{GeV} \\ \textbf{2.2} & m(\tilde{\chi}_1^0) < 700 \ \text{GeV} \\ \textbf{1.97} & m(\tilde{\chi}_1^0) < 600 \ \text{GeV} \end{array}$	2101.01629 CERN-EP-2022-014 2008.06032
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t t \tilde{\chi}_1^0$	SS e,μ 0-1 e,μ SS e,μ	6 jets 3 <i>b 1</i> 6 jets	139 T 79.8 139		1.15	m($ ilde{g}$)-m($ ilde{Y}_1^{'}$)=200 GeV 2.25 m($ ilde{X}_1^{(0)}$)<200 GeV 5 m($ ilde{g}$)-m($ ilde{X}_1^{(0)}$)=300 GeV	1909.08457 ATLAS-CONF-2018-041 1909.08457
	$\tilde{b}_1 \tilde{b}_1$	0 <i>e</i> , <i>µ</i>	2 b 1	T_T^{miss} 139	,	0.68	5 m($\tilde{\chi}_1^0$)<400 GeV 10 GeV< ∆m $\tilde{\chi}_1^0$ >20 GeV	2101.12527 2101.12527
3 rd gen. squarks direct production	$\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 {\rightarrow} b \tilde{\chi}^0_2 {\rightarrow} b h \tilde{\chi}^0_1$	0 <i>e</i> ,μ 2 τ	6 <i>b</i> 1 2 <i>b</i> 1	$T_T^{miss} = 139$:	Forbidden 0.23-1.	.35 $\Delta m(\tilde{k}_2^0, \tilde{k}_1^0) = 130 \text{ GeV}, m(\tilde{k}_1^0) = 100 \text{ GeV}$ $\Delta m(\tilde{k}_2^0, \tilde{k}_1^0) = 130 \text{ GeV}, m(\tilde{k}_1^0) = 0 \text{ GeV}$	1908.03122 2103.08189
	$ \begin{array}{l} \tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{t}_1^0 \\ \tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0 \\ \tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 bv, \tilde{\tau}_1 \rightarrow \tau \tilde{G} \\ \tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{c} \tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow \tilde{c} \tilde{\chi}_1^0 \end{array} $	0-1 <i>e</i> , μ 1 <i>e</i> , μ 3 1-2 τ 2 0 <i>e</i> , μ 0 <i>e</i> , μ π	≥ 1 jet β jets/1 b β jets/1 b 2 c mono-jet b	${ \begin{array}{ccc} miss \\ T \\$		1.25 Forbidden 0.65 Forbidden 0.85 0.55	5 $m(\xi^0)=1 \text{ GeV}$ $m(\xi^0)=500 \text{ GeV}$ 1.4 $m(\tau_1)=800 \text{ GeV}$ $m(\tau_1^0)=00 \text{ GeV}$ $m(\xi^0_1)=0 \text{ GeV}$	2004.14060.2012.03799 2012.03799 2108.07665 1805.01649 2102.10874
	$ \begin{array}{l} \tilde{t}_1 \tilde{t}_1, \tilde{t}_1 {\rightarrow} t \tilde{\chi}_2^0, \tilde{\chi}_2^0 {\rightarrow} Z/h \tilde{\chi}_1^0 \\ \tilde{t}_2 \tilde{t}_2, \tilde{t}_2 {\rightarrow} \tilde{t}_1 + Z \end{array} $	1-2 <i>e</i> , μ 3 <i>e</i> , μ	1-4 <i>b l</i> 1 <i>b l</i>	T_T^{miss} 139 T_T^{miss} 139		1 0.067-1.18 2 Forbidden 0.86	$m(\tilde{\chi}_{1}^{0})=500 \text{ GeV}$ $m(\tilde{\chi}_{1}^{0})=360 \text{ GeV}, m(\tilde{r}_{1})-m(\tilde{\chi}_{1}^{0})=40 \text{ GeV}$	2006.05880 2006.05880
EW direct	${ ilde \chi}_1^{\pm} { ilde \chi}_2^0$ via WZ	Multiple ℓ /jets $ee, \mu\mu$	∠ljet l	T_T^{miss} 139 T_T^{miss} 139		$\frac{1}{4}/\tilde{\chi}_{2}^{0}$ 0.96	$m(\tilde{\chi}_1^0)=0$, wino-bino $m(\tilde{\chi}_1^{\pm})$ - $m(\tilde{\chi}_1^0)=5$ GeV, wino-bino	2106.01676, 2108.07586 1911.12606
	$ \begin{split} \tilde{\chi}_1^{\dagger} \tilde{\chi}_1^{\dagger} & \text{via } WW \\ \tilde{\chi}_1^{\dagger} \tilde{\chi}_2^0 & \text{via } Wh \\ \tilde{\chi}_1^{\dagger} \tilde{\chi}_1^{\dagger} & \text{via } \tilde{\ell}_L/ \wp \\ \tilde{\tau}_1^{\dagger} \tilde{\tau}_1^{\dagger} & \text{via } \tilde{\ell}_L R_1^0 \\ \tilde{\ell}_{\text{LR}} \tilde{\ell}_{\text{LR}}, \tilde{\ell} \to \ell \tilde{\chi}_1^0 \end{split} $	$\begin{array}{c} 2 \ e, \mu \\ \text{Multiple } \ell/\text{jets} \\ 2 \ e, \mu \\ 2 \ \tau \\ 2 \ e, \mu \\ ee, \mu \mu \end{array}$	0 jets 1 2 ≥ 1 jet 1	${ { { { { { T} } { T } } } } { { T } } } { { T } } { {$		t 0.42 / k ² Forbidden 1.06 t (⁷ L, ⁷ R,L) 0.16-0.3 0.12-0.39 0.256 0.7	$\begin{split} & m(\tilde{\xi}_1^0){=}0, \text{ wino-bino} \\ & m(\tilde{\xi}_1^0){=}70 \text{ GeV, wino-bino} \\ & m(\tilde{\ell},\tilde{v}){=}0.5(m(\tilde{\xi}_1^0){+}m(\tilde{\xi}_1^0)) \\ & m(\tilde{\ell},\tilde{v}){=}0 \\ & m(\tilde{\xi}_1^0){=}0 \\ & m(\tilde{\xi}_1^0){=}10 \text{ GeV} \end{split}$	1908.08215 2004.10894,2108.07586 1908.08215 1911.06660 1908.08215 1911.12606
	ĤĤ, Ĥ→hĜ/ZĜ	$\begin{array}{l} 0 \ e, \mu \\ 4 \ e, \mu \\ 0 \ e, \mu \end{array} \ge 2$	$\geq 3 b$ 1 0 jets 1 2 large jets 1	\sum_{T}^{miss} 36.1 \sum_{T}^{miss} 139 T 139		r 0.13-0.23 0.29-0.88 r 0.55 r 0.45-0.93	$\begin{array}{l} BR(\tilde{k}_{1}^{O} \to h\bar{G}) \!=\! 1 \\ BR(\tilde{k}_{1}^{O} \to Z\bar{G}) \!=\! 1 \\ BR(\tilde{k}_{1}^{O} \to Z\bar{G}) \!=\! 1 \end{array}$	1806.04030 2103.11684 2108.07586
٦	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet 1	Σ_T^{miss} 139		± 0.66	Pure Wino Pure higgsino	2201.02472 2201.02472
Long-lived particles	Stable \tilde{g} R-hadron Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$ $\tilde{\ell}\tilde{\ell}, \tilde{\ell} \rightarrow \ell \tilde{G}$	pixel dE/dx pixel dE/dx Displ. lep pixel dE/dx	E I I	$\begin{array}{ccc} T_T & 139\\ T_T & 139\end{array}$		[r(ĝ) =10 ns] ,μ 0.34 0.7 0.36	2.05 2.2 $m(\tilde{t}^0_1)=100 \text{ GeV}$ $\tau(\tilde{t})=0.1 \text{ ns}$ $\tau(\tilde{t})=0.1 \text{ ns}$ $\tau(\tilde{t})=10 \text{ ns}$	CERN-EP-2022-029 CERN-EP-2022-029 2011.07812 2011.07812 CERN-EP-2022-029
RPV	$\begin{array}{c} \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{+},\tilde{\chi}_{1}^{0},\tilde{\chi}_{1}^{+}\rightarrow\mathcal{Z}\ell\rightarrow\ell\ell\ell\ell\\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{+},\tilde{\chi}_{2}^{0}\rightarrow WW/Z\ell\ell\ell\ell\nu\nu\\ \tilde{\chi}_{2}^{0},\tilde{\chi}_{1}^{0}\rightarrow WW/Z\ell\ell\ell\ell\nu\nu\\ \tilde{\chi}_{3}^{0},\tilde{\chi}_{1}^{0}\rightarrow Uhs\\ \tilde{\pi}_{1},\tilde{\pi}_{1}\rightarrow \tilde{\chi}_{1}^{0},\tilde{\chi}_{1}^{0}\rightarrow Uhs\\ \tilde{\pi}_{1},\tilde{\pi}_{1}\rightarrow bs\\ \tilde{\eta}_{1}\tilde{\eta}_{1},\tilde{\eta}\rightarrow bs\\ \tilde{\eta}_{1}\tilde{\eta}_{1},\tilde{\eta}\rightarrow d\ell\\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{2}^{0},\tilde{\chi}_{1}^{0},\tilde{\chi}_{1}^{0},\sigma\rightarrow ths,\tilde{\chi}_{1}^{+}\rightarrow bbs\end{array}$	$3 e, \mu$ $4 e, \mu$ 4.5 1 2] $2 e, \mu$ 1μ $1.2 e, \mu$	0 jets b i large jets Multiple $\geq 4b$ jets + 2 b 2 b DV ≥ 6 jets	139 T 139 36.1 36.1 139 36.7 36.1 139 36.7 36.1 136		$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	$\begin{array}{c c} & & & & & \\ \textbf{1.55} & & & & & & \\ \textbf{1.3} & \textbf{1.9} & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ \textbf{1.4.5} & & & & & \\ \textbf{1.6} & & & & & \\ & & & & & & \\ & & & & & & $	2011.10543 2103.11684 1804.03568 ATLAS-CONF-2018-003 2010.01015 1710.07171 1710.05544 2003.11956 2106.09609
Only	a selection of the available ma	ass limits on ne	w states (or	10		Mass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

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Search for heavy resonances in the decay channel WW $\!\!\rightarrow\!\!ev\mu v$

 Search for the neutral heavy resonances, e.g. heavy CP even Higgs (ggF + VBF) and neutral GM Higgs (VBF)

Interpretations

Interpreted in 5 models

GM

2HDM

HVT

Radion in Randall-Sundrum modelSpin-2 graviton



ATLAS-CONF-2022-066

Selection

3 signal categories (ggF, VBF+1jet, VBF+2jets)

Strategy

•Discriminating variable: $m_T = \sqrt{(E_T^{\ell \ell} + E_T^{\text{miss}})^2 - |\vec{p}_T^{\ell \ell} + \vec{E}_T^{\text{miss}}|^2}$



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Search for heavy resonances in the decay channel WW \rightarrow evµv







$H \pm \rightarrow WZ \rightarrow l\nu l' l'$: Search for resonant $WZ \rightarrow \ell \nu \ell' \ell'$ production

Drell-Yan

ArXiv 2207.03925

(c)



Interpretation

- Heavy vector triplet model HVT
- <u>Georgi-Machacek (GM)</u> fermiophobic charged Higgs boson from the fiveplet: $(H_5^{++}, H_5^+, H_5^0, H_5^-, H_5^{--})$

Selection

2 signal regions: Drell-Yan and VBF

Strategy

VBF category: NN is adopted to enhance signal



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Vector Boson Fusion (VBF)



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VH(→hh): Search for Higgs boson pair production in association with a vector boson ArXiv:2210.05415

Search for a neutral heavy CP-even scalar H→hh→bbbb in association with a vector boson (W/Z)
 First analysis targeting Vhh final state.

Interpretation

Two signal scenarios are considered:

1) Generic neutral CP-even scalar H

2) CP-odd scalar A decaying into Z and a CP-even scalar *H*



Selection

 Three leptonic channels (0L, 1L, 2L)targeting vvbbbb, lvbbbb, llbbbb final states

Invariant mass resolution improved by scaling the b-jet momenta with the ratio of the measured di-b-jet mass to 125 GeV.



•The relative mass resolution improves by a factor of \sim 3!

Strategy

BDTs used as final discriminants trained with jets kinematic information in each SR



$VH(\rightarrow hh)$: Search for Higgs boson pair production in association with a vector boson ArXiv:2210.05415



Largest deviation

Narrow width A boson

(mA, mH) = (800, 300) Local: 3.6 σ Global: 1.6 σ Large-width (20%) A boson

(mA, mH) = (420, 320) Local: 3.8 σ Global: 2.8 σ

Constraints in the 2HDM type-I parameter space





Search for heavy flavour-violating Higgs bosons in multilepton plus b-jets final states ATLAS-CONF-2022-039

Interpretation

•General two Higgs doublet model (g2HDM) without Z2symmetry, featuring FCNH (Flavour Changing Neutral Higgs) couplings.

Couplings of H involving top-quark: ptt, ptc, and ptu.

First analysis to target BSM production leading to three-top final states and the first to probe g2HDM. Sensitive probe of new physics: same-sign (SS) top, 3-top, 4-top final states.



Strategy



Scan ptt, ptc, ptu couplings and heavy Higgs masses 200 GeV < mH < 1000 GeV

Most significant deviation observed at mH=1000 GeV with local significance of 2.810 local

0.5

Results are also interpreted in **RPV SUSY model.**

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Search for type-III seesaw heavy leptons in leptonic final states

ArXiv 2202.02039

Interpretation

•At least one extra fermionic SU(2)L triplet predicted by the type-III seesaw mechanism: •Pair production of new heavy charged (L^{\pm}) and neutral (N^{0}) lepton particles

Selection and strategy



Search for electroweakinos production in final states with one lepton, jets and MET ATLAS-CONF-2022-059

First search targeting these processes in the 1Lepton channels

Event Selection

Variable	C1C1-WW model			C1N2-WZ model			
	SRLM	SRMM	SRHM	SRLM	SRMM	SRHM	
$N_{\text{lep}} (p_{\text{T}} > 25 \text{ GeV})$			1	1			
$N_{\rm jet} (p_{\rm T} > 30 {\rm GeV})$			1 -	- 3			
$N_{\text{large-Rjet}} (p_{\text{T}} > 250 \text{ GeV})$							
$E_{\rm T}^{\rm miss}$ [GeV]	> 200						
$\Delta \dot{\phi}(\ell, \mathrm{E}_{\mathrm{T}}^{\mathrm{miss}})$	< 2.6						
large-R jet type	W-tagged			Z-tagged			
$m_{\rm T}$ [GeV]	120-200	200-300	> 300	120-200	200-300	> 300	
	Exclusion SR						
$m_{\rm eff}$ [GeV] (excl.)	[60	0-850, > 85	50]	[600-850, > 850]			
$m_{\rm ii}[{\rm GeV}]$ (excl.)	[70–90, -]			[80–100, -]			
$\sigma_{E_{T}^{\text{miss}}}$ (excl.)	[> 12, > 15]			[> 12, > 12]			
i	Discovery SR						
$m_{\rm eff}$ [GeV] (disc.)	> 600	> 600	> 850	> 600	> 850	> 850	
m _{ii} [GeV] (disc.)	-	-	-	80-100	-	-	
$\bar{\sigma}_{E_{\mathrm{T}}^{\mathrm{miss}}}$ (disc.)	> 15	> 15	> 15	> 12	> 12	> 12	

Strategy





Results $\widetilde{\chi}_{*}^{\pm}\widetilde{\chi}_{*}^{\mp} \rightarrow WW \ \widetilde{\chi}_{*}^{0}\widetilde{\chi}_{*}^{0}, W \rightarrow Iv, W \rightarrow q\overline{q}$ 400 $n(\widetilde{\chi}_{1}^{0})$ [GeV] ATLAS Preliminary 350 √s=13 TeV, 139 fb⁻¹, All limits at 95% CL Expected Limit $(\pm 1 \sigma_{exp})$ 300 Observed Limit (±1 of theory) 250 200 150 100 50 0[200 300 400 500 600 700 800 900 1000 $m(\tilde{\chi}_{1}^{\pm})$ [GeV]

- •For the $\chi^{-1} + \chi^{-1}$ model, chargino masses ranging from 260 to 520 GeV can be excluded for a massless χ^{-1} 0
- x⁻¹± x⁻²⁰ process, degenerate chargino/ neutralino masses ranging from 260 to 420 GeV can be excluded for a massless x⁻¹⁰