

Extraction of level density of 2^+ states in ^{208}Pb and ^{120}Sn nuclei from high energy-resolution (p,p') experiments

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<https://www.wits.ac.za/physics/research-areas/nuclear-and-radiation-physics/>



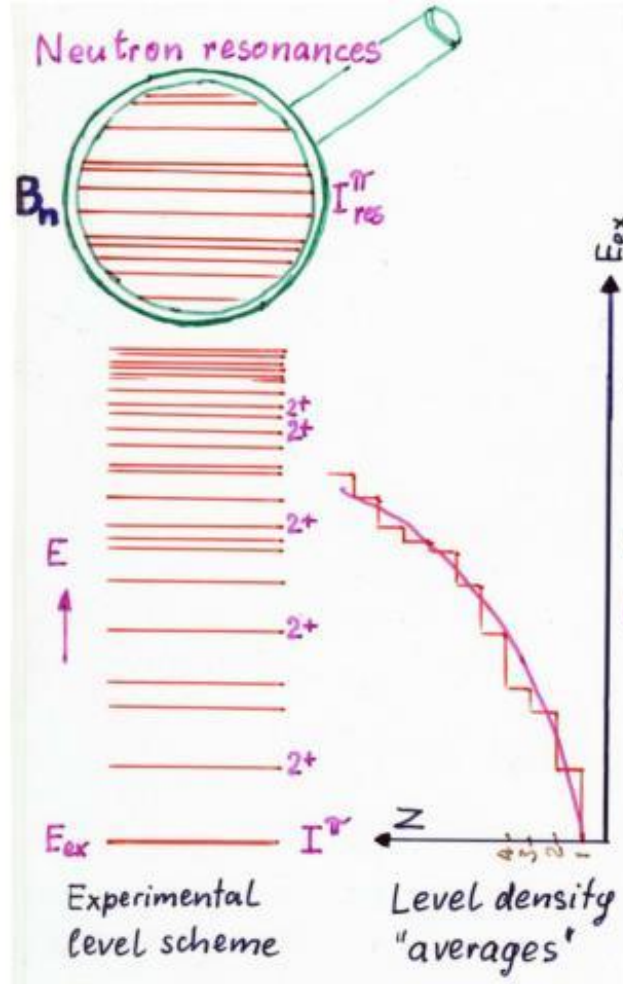
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Outline

- Motivation
- Experiments
- Theoretical Models
- Fluctuation Analysis and Autocorrelation functions
- Results
- Conclusion and Outlook

Motivation

Energy Levels of nuclei



Level densities: averages

Average level density $\rho(E)$:

$$\rho(E) = dN/dE = 1/D(E)$$

Cumulative number $N(E)$

Average level spacing D

Level spacing $S_i = E_{i+1} - E_i$

$D(E)$ determined by fit to individual level spacings S_i

Level spacing correlation:

Chaotic properties determine fluctuations about the averages and the errors of the LD parameters.

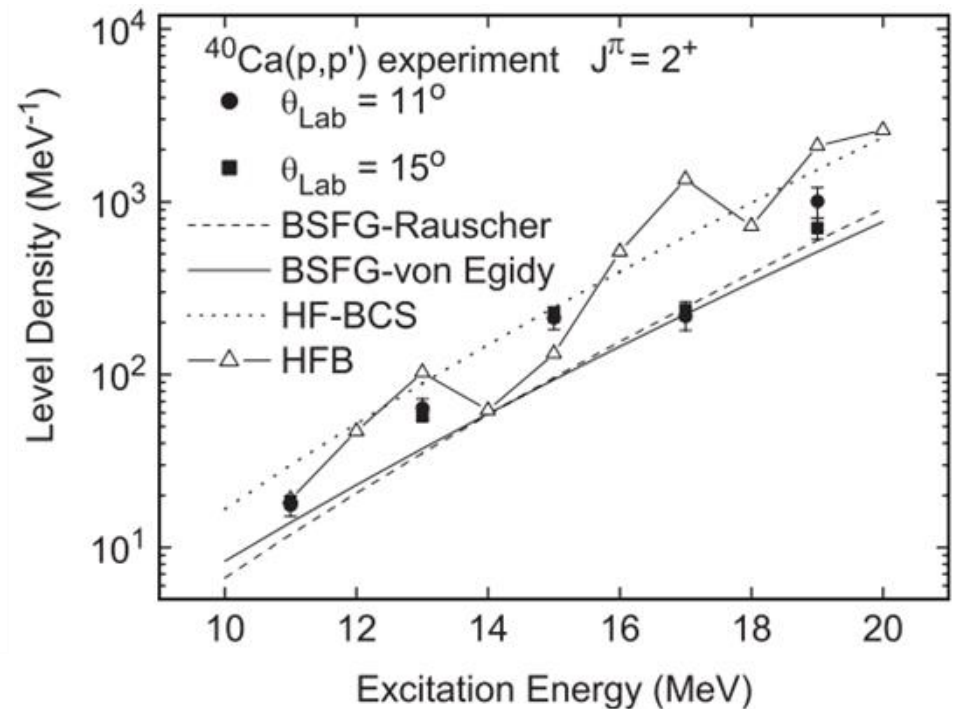
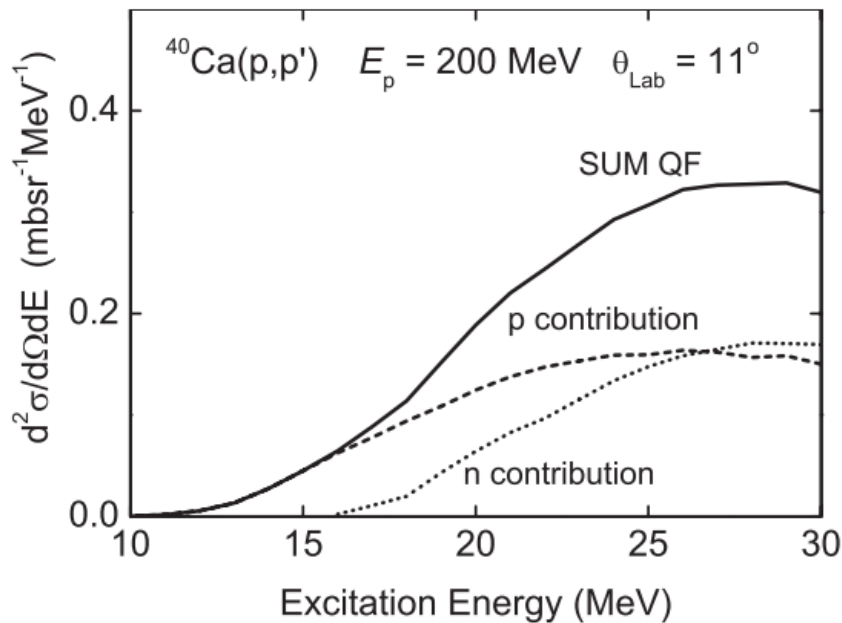
- High density of unresolved states above S_n
- Is there a direct measurement covering both low-lying and giant resonance regions?

Motivation

PHYSICAL REVIEW C **84**, 054322 (2011)

Level density of 2^+ states in ^{40}Ca from high-energy-resolution (p,p') experiments

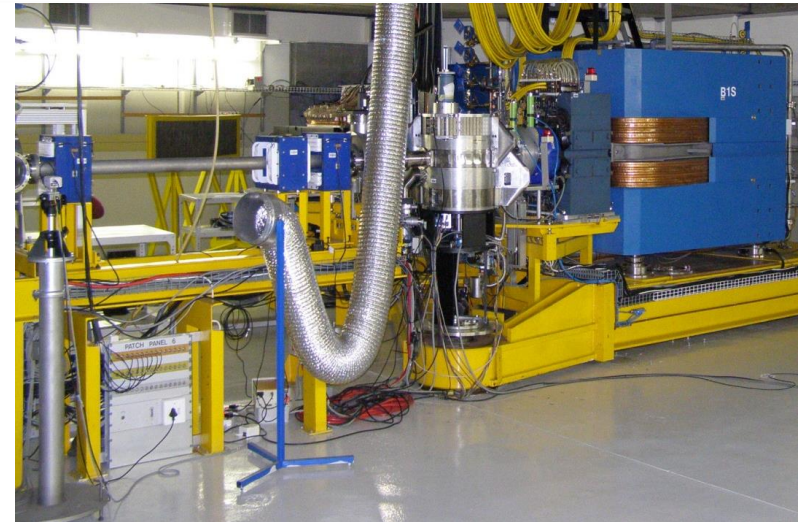
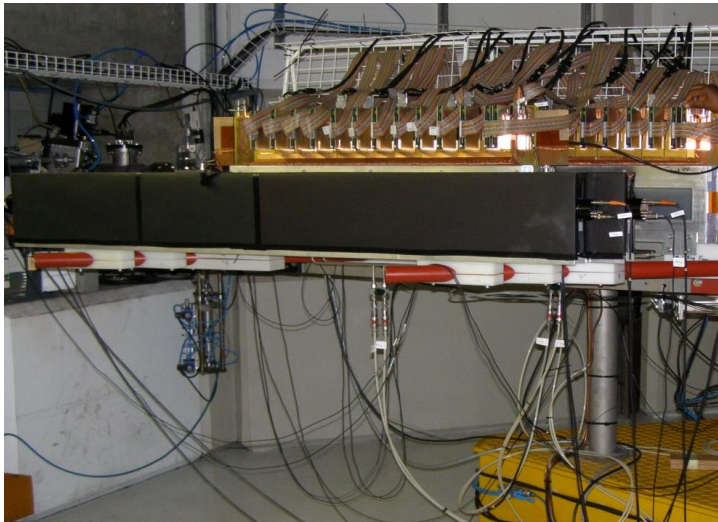
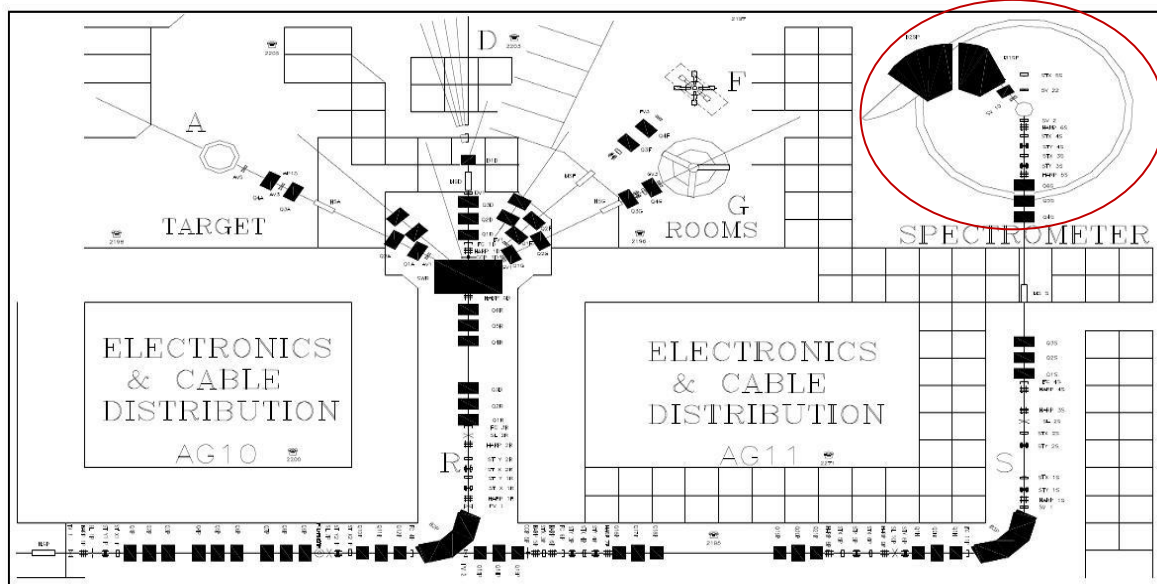
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F. D. Smit,¹ and J. Wambach⁵



Backgrounds determined using Quasi-free + DWT

Experiments

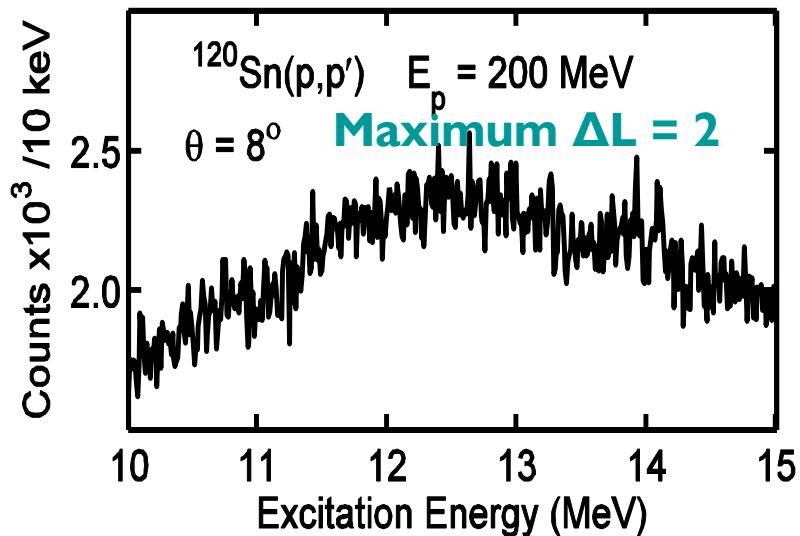
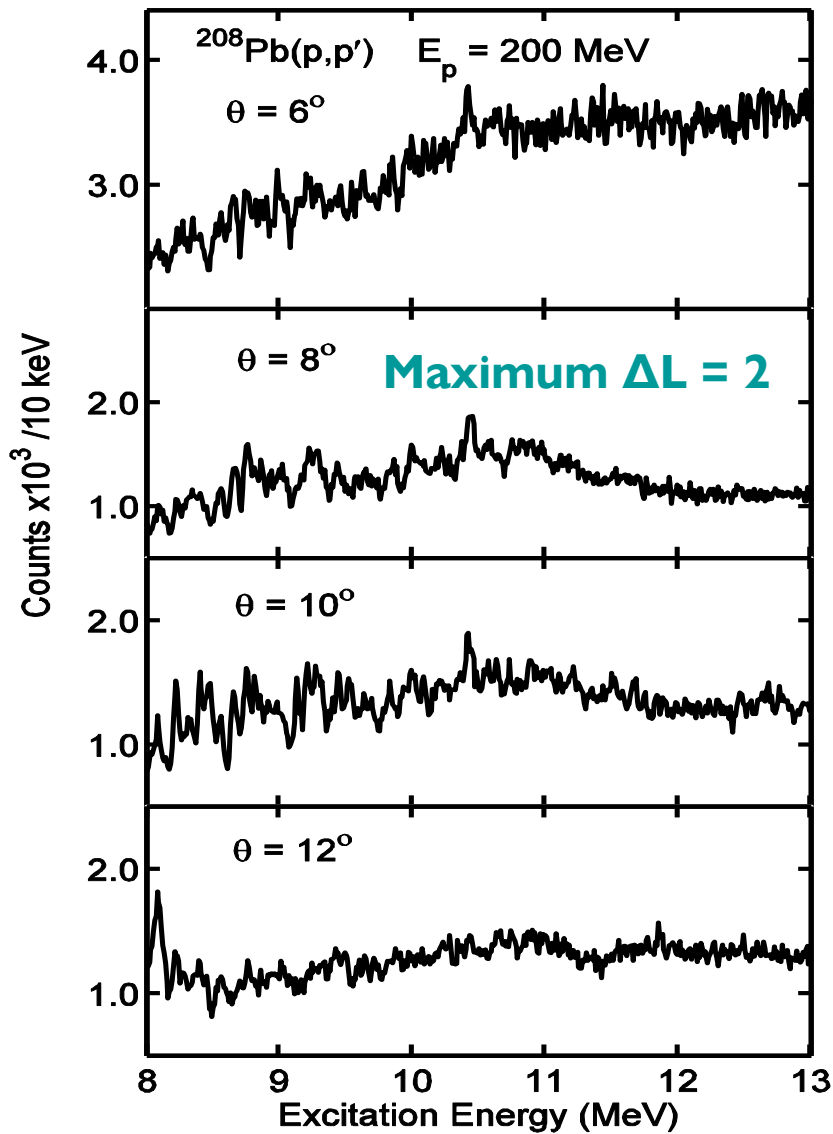
Nuclear Physics Beamlines at iThemba LABS



✓ High energy-resolution → 25 – 45 keV @ $E_p = 200$ MeV

✓ Dispersion matching technique

Energy Spectra: ^{208}Pb , ^{120}Sn



A. Shevchenko, *et al.*, *Phys. Rev. C* **77**, (2008) 024302.

A. Shevchenko *et al.*, *Phys. Rev. Lett.* **93**, (2004) 122501.

Theoretical Models

1) Back-Shifted Fermi Gas (BSFG)

Phenomenological Approach

T. Rauscher, F. K. Thielemann, and K.-L. Kratz, Phys. Rev. C 56, (1997) 1613.

T. von Egidy and D. Bucurescu, Phys. Rev. C 80, (2009) 054310.

$$\rho(E_x, J) = \frac{2J+1}{2\sigma^2} \frac{\exp\left(2\sqrt{a(E_x - \delta)} - J\left(J + \frac{1}{2}\right)^2 / 2\sigma^2\right)}{12\sqrt{2\sigma}a^{1/4}(E_x - \delta)^{5/4}}$$

with

$$\delta = \Delta(Z, N)$$

$\rho(E_x)$ = Level density at energy E_x

σ = Spin-cutoff parameter

a = Level density parameter

δ = Backshift energy

Δ = Pairing energy

Theoretical Model Parameters

	^{208}Pb	^{120}Sn
mass number A	208	120
number of protons Z	82	50
number of neutrons N	126	70
spin of states	2	2
parity of states	+	+
back shift in MeV	1.9089	1.9620
microscopic energy correction in MeV	-12.84	0.19
mass of the (Z,N) nucleus in MeV	193729	111688
mass of the (Z+1,N+1) nucleus in MeV;	195598	113553
mass of the (Z-1,N-1) nucleus in MeV	191865	109829
pairing energy in MeV	-3.22602	-3.3248
Proton separation energy in MeV	8.00376	10.6895
Neutron separation energy in MeV	7.36786	9.10802
Alpha separation energy in MeV	-0.51687	4.81059
local level density parameter BSFG in MeV^{-1}	10.01	13.14
local energy shift BSFG in MeV	1.17	0.85

2) Hartree-Fock-Bogoliubov (HFB) and HF-BCS

Microscopic Single-Particle Levels Approach

P. Demetriou and S. Goriely, Nucl. Phys. A 695, (2001) 95.

S. Goriely, S. Hilaire, and A. J. Koning, Phys. Rev. C 78, (2008) 064307.

➤ Calculations using basic nuclear structure properties:

Single-particle energies \mathcal{E}_i^k

Pairing strength Δ_i^k

Quadrupole deformation parameter β_2

Deformation energy E_{def}

➤ For spherical nuclei:

$$\rho_{sph}(E, J, \pi) = \rho_i(E, M = J, \pi) - \rho_i(E, M = J + 1, \pi)$$

➤ For deformed nuclei:

$$\rho_{def}(E, J, \pi) = \frac{1}{2} \left[\sum_{k=-J, k \neq 0}^J \rho_i(E - E_{rot}^{J,k}, K, \pi) \right] + \delta_{(J \text{ even})} \delta_{(\pi=+)} \rho_i(E - E_{rot}^{J,0}, 0, \pi) \\ + \delta_{(J \text{ odd})} \delta_{(\pi=-)} \rho_i(E - E_{rot}^{J,0}, 0, \pi)$$

Fluctuation Analysis

Measure of cross section fluctuations with respect to a stationary mean value.

Assumptions:

$$\langle \Gamma \rangle \leq \langle D \rangle \leq \Delta E$$

$$\alpha = \alpha_w + \alpha_{PT}$$

$\langle \Gamma \rangle$ Mean level width

$\langle D \rangle$ Mean level spacing

ΔE Energy resolution

α Sum of normalised variances

Procedure:

- Background subtraction from the experimental spectrum
- Smoothing by convolution with a Gaussian function of width larger than $\Delta E \longrightarrow g_>(E_x)$
- Folding with a Gaussian function with a width smaller than $\Delta E \longrightarrow g(E_x)$
- Create a stationary spectrum $\longrightarrow \langle d(E_x) \rangle = \left\langle \frac{g(E_x)}{g_>(E_x)} \right\rangle = 1$

Autocorrelation Function

➤ Mean level spacing $\langle D \rangle = \frac{1}{\langle \rho \rangle}$

proportional to the variance of $\langle d(E_x) \rangle$

➤ Intensity fluctuations in $\langle d(E_x) \rangle$ can be autocorrelated at energies E and $E + \varepsilon$

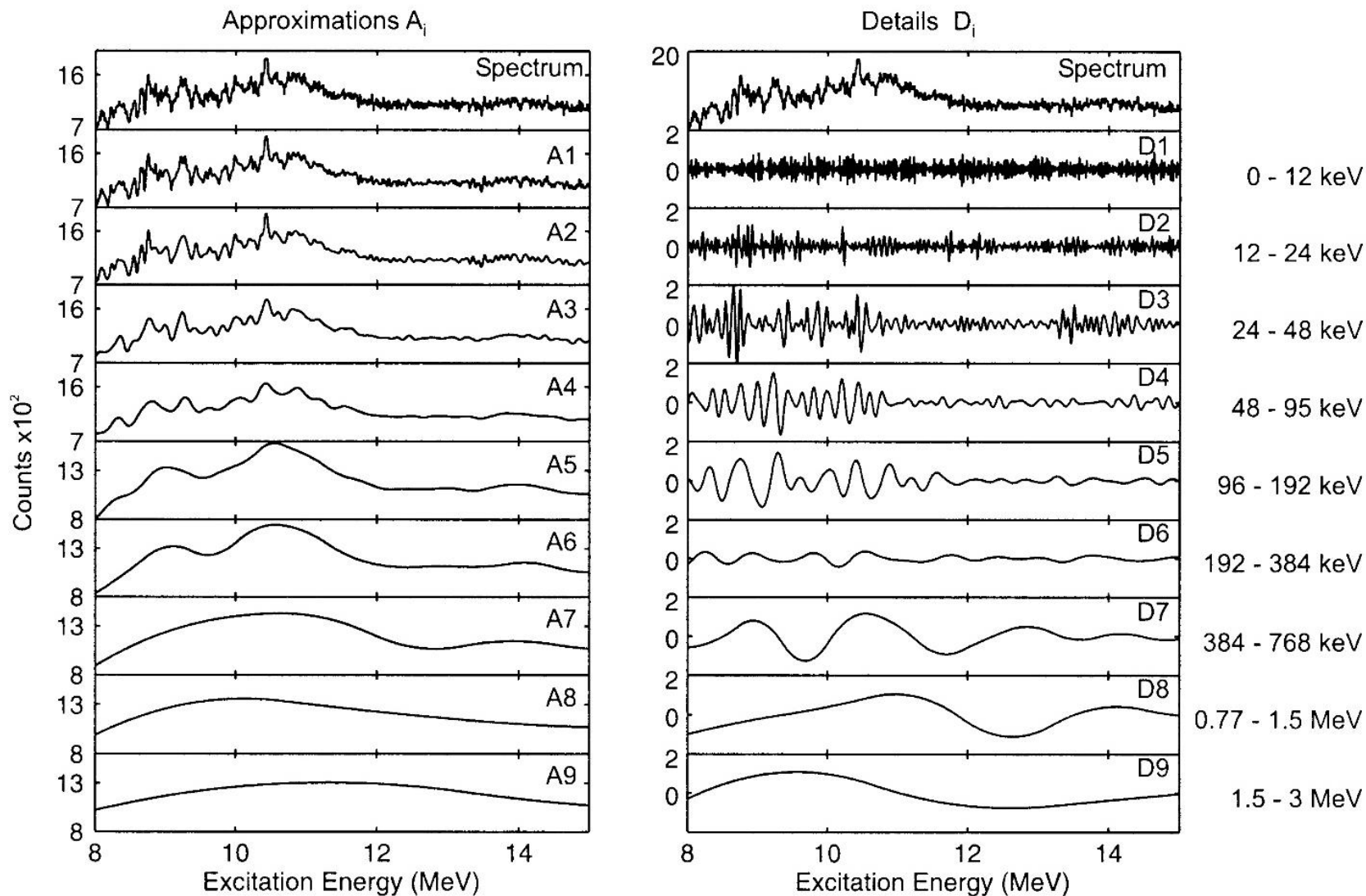
$$C(\varepsilon) = \frac{\langle d(E_x) d(E_x + \varepsilon) \rangle}{\langle d(E_x) \rangle \langle d(E_x + \varepsilon) \rangle}$$

➤ $\langle D \rangle$ can be extracted from

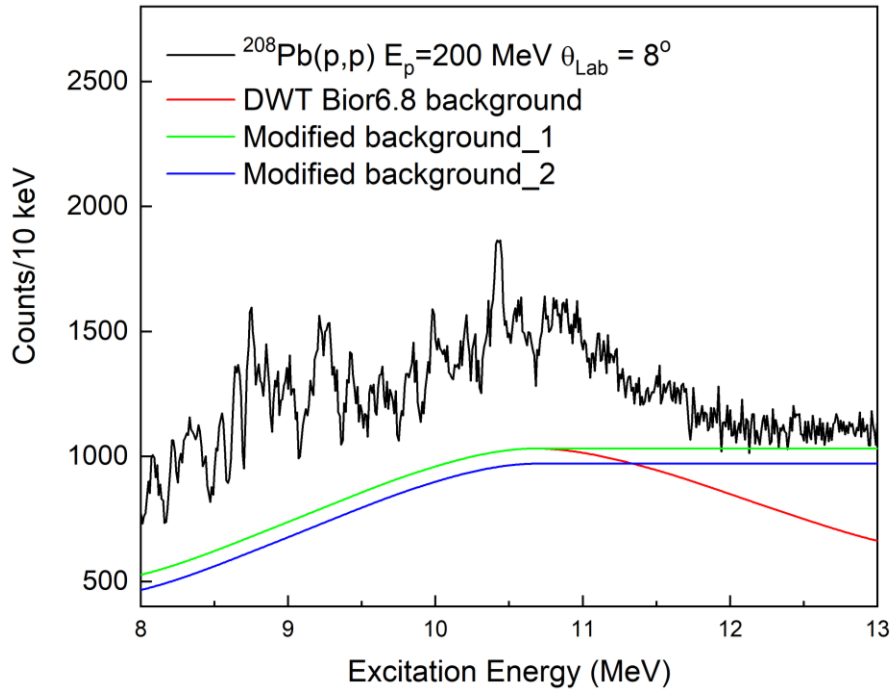
$$C(\varepsilon) - 1 = \frac{\alpha \langle D \rangle}{2\Delta E \sqrt{\pi}} f(\varepsilon, \Delta E)$$

$\varepsilon = \text{energy increment}$; $C(\varepsilon) = \text{Autocorrelation function}$

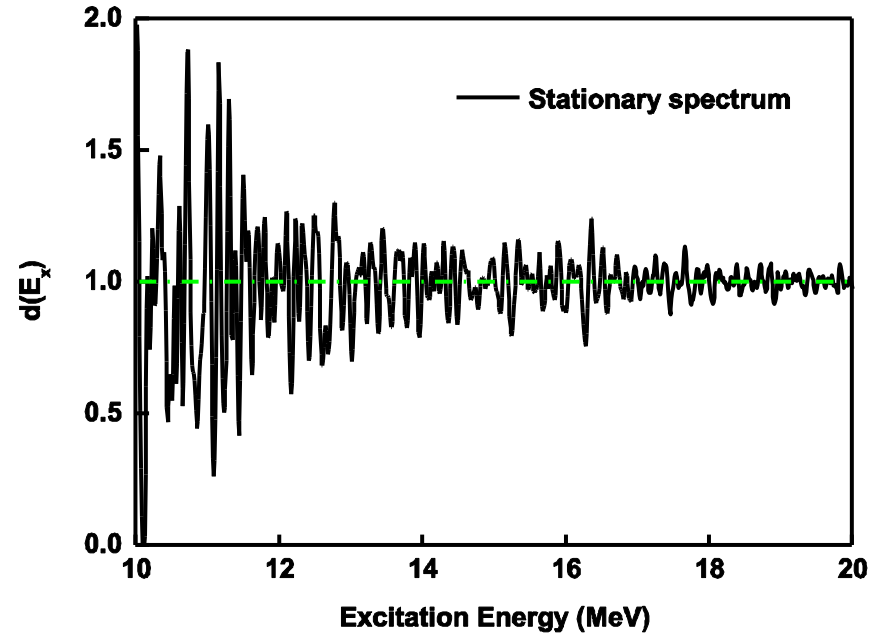
Discrete Wavelet Transform (DWT) of $^{208}\text{Pb}(p,p')$ data



Background subtraction, Fluctuation Analysis, and Autocorrelation function

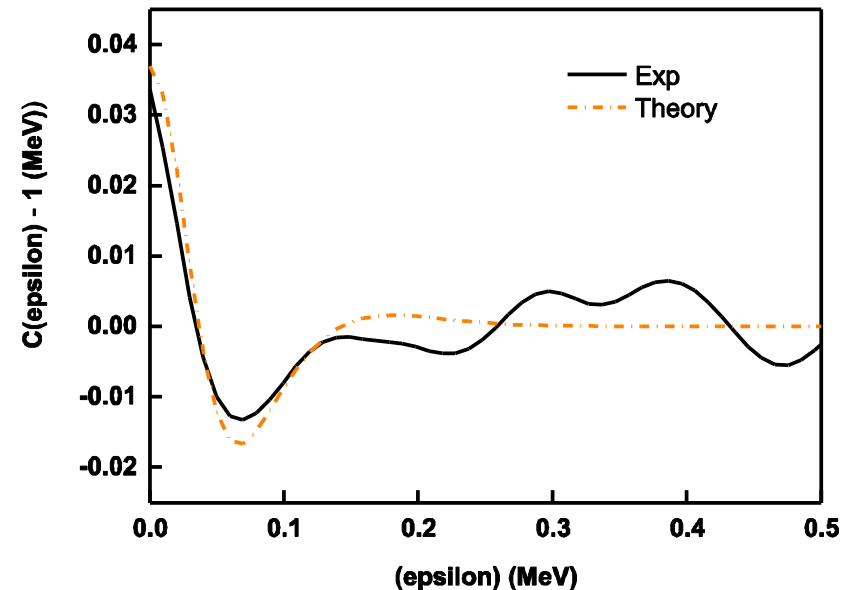


Top Left: Data plus background

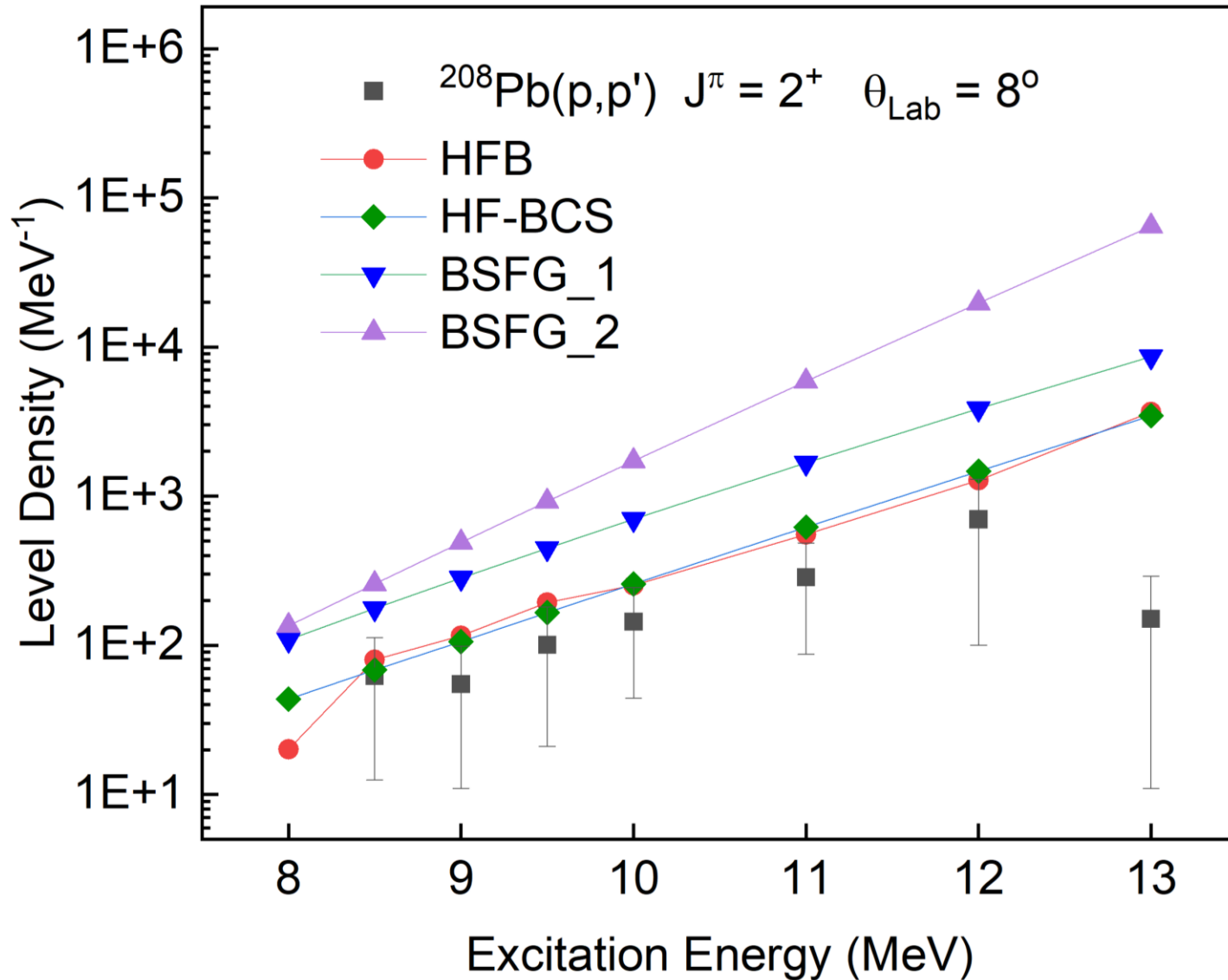


Top Right: Stationary Spectrum averaged around 1

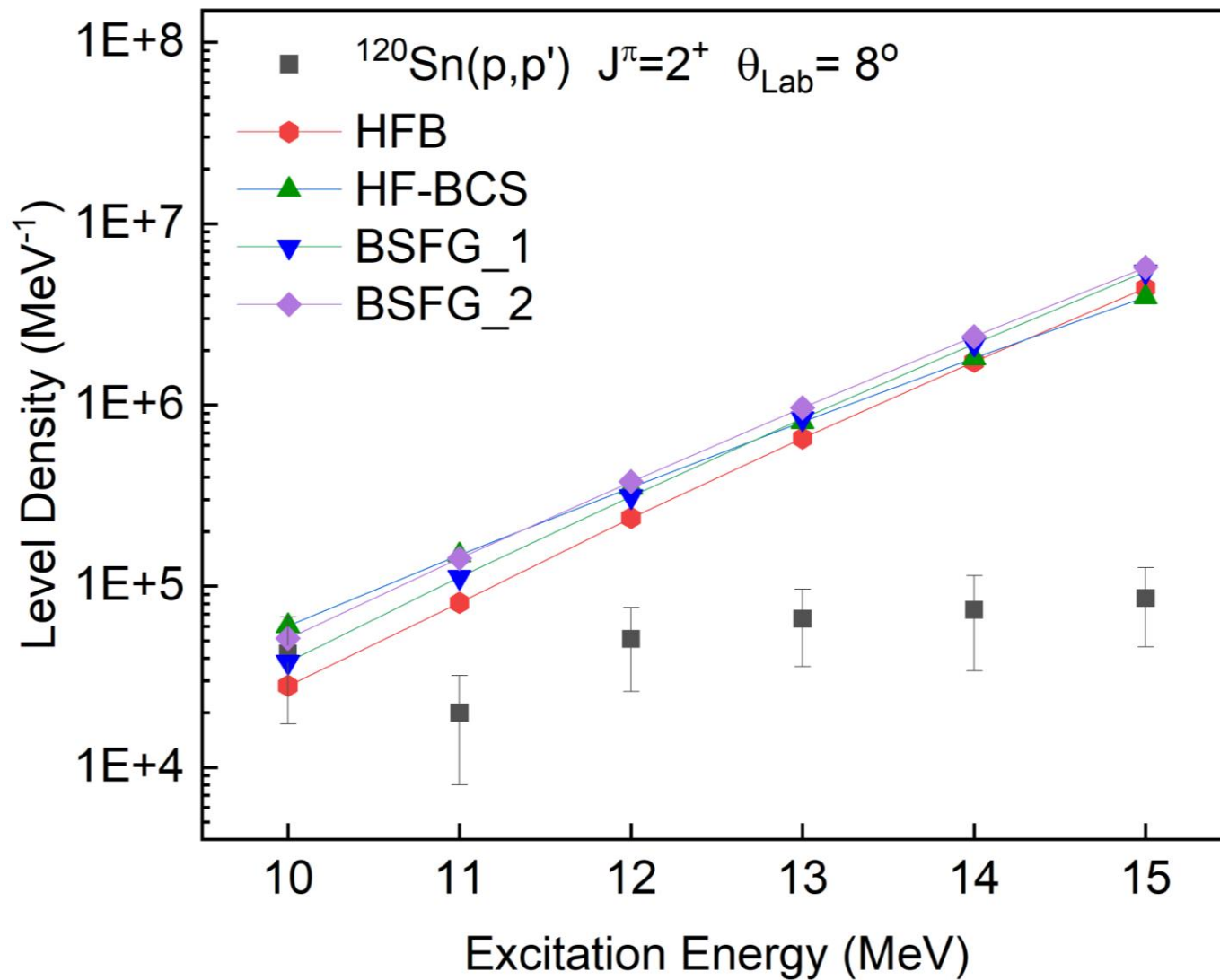
Bottom Right: Autocorrelation function



Experimental and Theoretical Level Density: ^{208}Pb



Experimental and Theoretical Level Density: ^{120}Sn



Summary

- Experimental level density in ^{208}Pb and ^{120}Sn using Autocorrelation function analysis provides fair agreements compared to all models.
- The selectivity of the ISGQR and high energy-resolution data obtained at iThemba LABS makes it an ideal choice for level density extraction above neutron threshold
- Parity dependence results was found to be important.

Conclusions *and* Outlook

- Extracted level densities in the giant resonance region of heavy magic nuclei (^{208}Pb and ^{120}Sn) are considered very low.
- ISGQR data in heavy nuclei is prone to very low peak to background ratio.
- Uncertainties can be due to the assumed background which might not include background due to quasi-free continuum.
- Calculations on ^{90}Zr and ^{58}Ni in progress.
- Comparison with $M2(2^-)$ data from electron spectrometer to be investigated.

Collaborators

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