

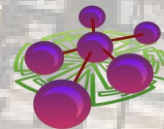
Reaction dynamics in the $^{58}\text{Ni}+^{58}\text{Ni}$ system at intermediate energies.

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The African Nuclear Physics Conference 2025
24-28 November 2025



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Physics case

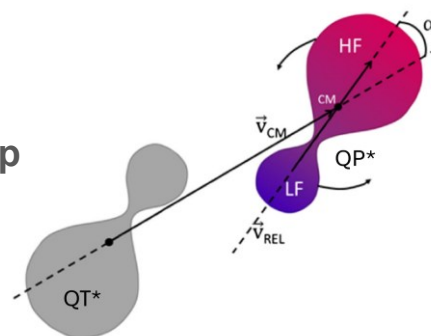
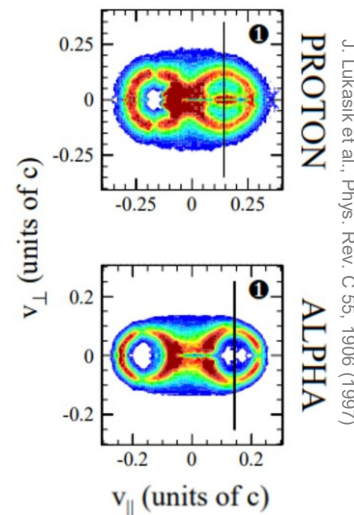
Heavy-Ion collision in the Fermi energy domain

Semiperipheral collisions at intermediate energies
($20 < E_{\text{beam}} < 100$ A MeV):

- **The dominant channel is the binary one:** formation of two excited fragments, which preserve memory of the entrance channel.

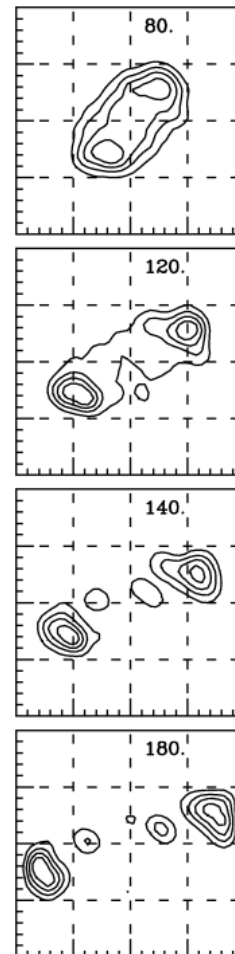
In the ejectiles we observe the superposition of two effects:

- **Statistical:** decay from a thermodynamically equilibrated source
- **Dynamical:**
 - emission towards midvelocity \rightarrow neck
 - Neutron enrichment of the midvelocity
- Another possible exit channel is the **dynamical breakup** of the QP and/or QT.



A. Jedyte et al., Phys. Rev. Lett. 118, 062501 (2017)

BNV calculation of density contour plots



V. Baran et al., Nucl. Phys. A 703, 603 (2002).

Physics case

Heavy-Ion collision and the isospin drift

An experimental observation is the **neutron enrichment of the neck**, which can be interpreted in the framework of the **nuclear Equation of State**

V. Baran et al. / Physics Reports 410 (2005) 335 – 466

$$\frac{E}{A}(\rho, I) = \frac{E}{A}(\rho) + \frac{E_{sym}}{A}(\rho)I^2$$

$$I = \frac{N - Z}{A}$$

$$\rho = \rho_n + \rho_z$$

Symmetric matter

Asymmetric matter

$$\frac{E}{A}(\rho) = E_{sat} + \frac{1}{2}K_{sat}x^2 + \dots \quad x = \frac{\rho - \rho_0}{\rho_0}$$

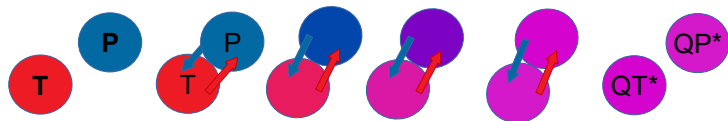
$$\frac{E_{sym}}{A}(\rho) = E_{sym} + L_{sym}x + \frac{1}{2}K_{sym}x^2 + \dots$$

The symmetry energy produces a difference of the proton and neutron currents:

$$\dot{j}_n - \dot{j}_p \propto \frac{E_{sym}}{A}(\rho) \nabla I + \frac{\partial E_{sym}/A}{\partial \rho} \nabla \rho \longrightarrow$$

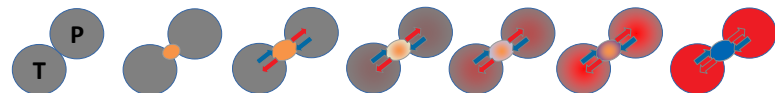
Isospin Diffusion:

- driven by the isospin gradient
- tends to equilibrate neutron–proton imbalance.



Isospin Drift:

- driven by the density gradient.
- induces a net neutron flux towards low density regions



Physics case

$^{58}\text{Ni}+^{58}\text{Ni}$ @ 32, 52 and 74 AMeV

$^{58}\text{Ni}+^{58}\text{Ni}$ at 32, 52 and 74 AMeV: evolution of the dynamics with incident energy and centrality in semispherical collisions.

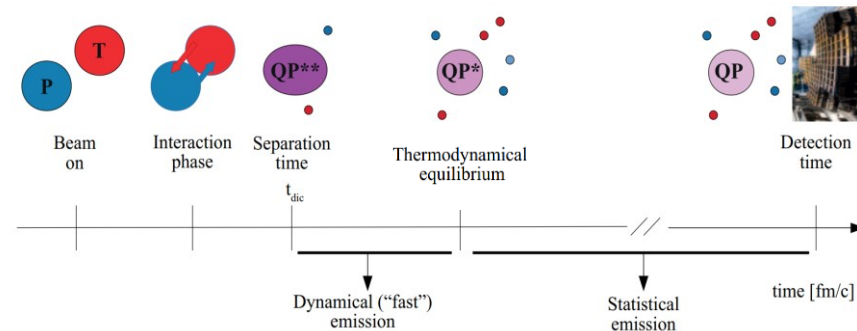
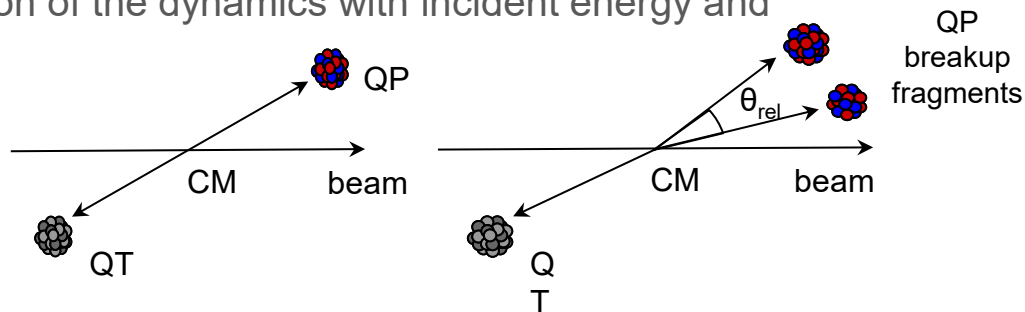
- QP phase space: selection of the QP evaporative (**QPr**) and the QP breakup (**QPb**) channels

- Centrality estimator:

reduced momentum $p_{red} = \left(\frac{p_z^{QP}}{p_{beam}} \right)$

- From the cold detected fragments, we aim at infer information on the dynamical phase:

- **Primary source reconstruction:** size and excitation energy of the primary QP* and the neck source
- **Isospin drift:** isospin investigation of the evaporative and midvelocity emissions



Experimental setup

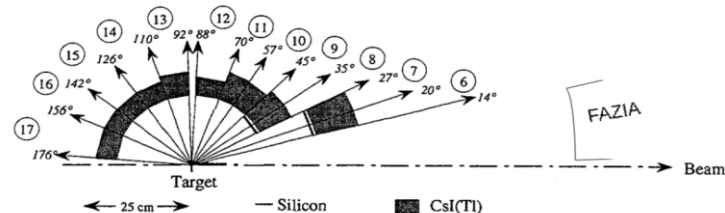
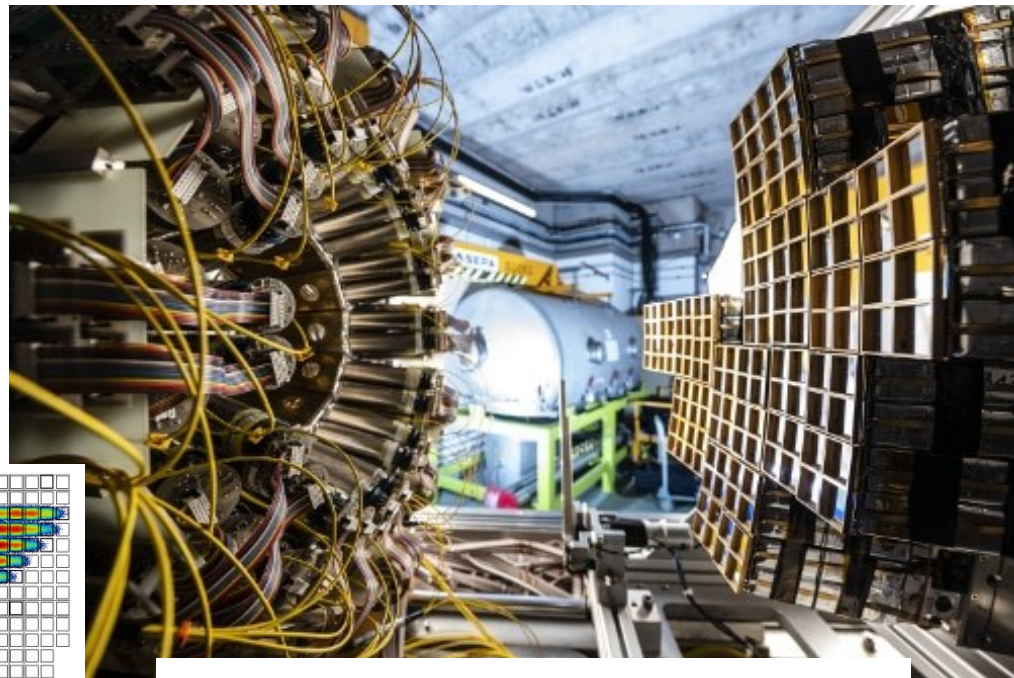
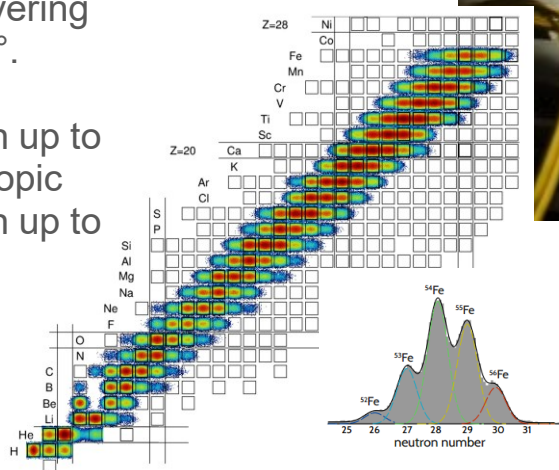
The INDRA-FAZIA apparatus

INDRA

- 12 rings covering from 14° to 176°
- Charge discrimination up to uranium, mass discrimination up to $Z \sim 10$.

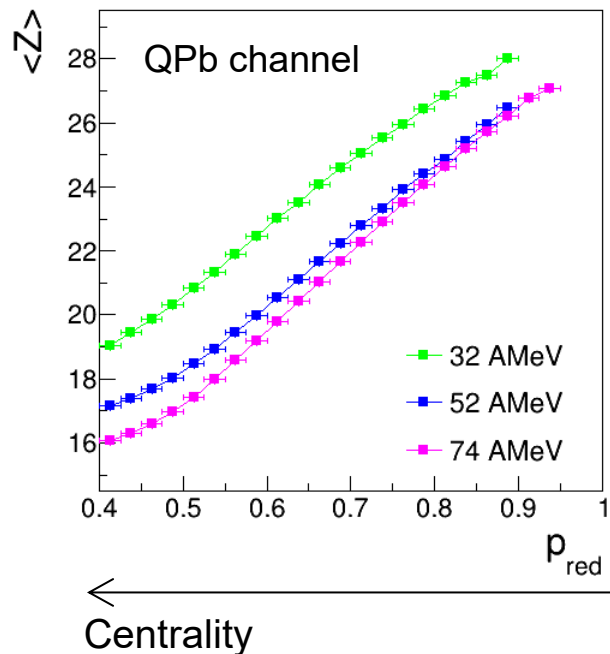
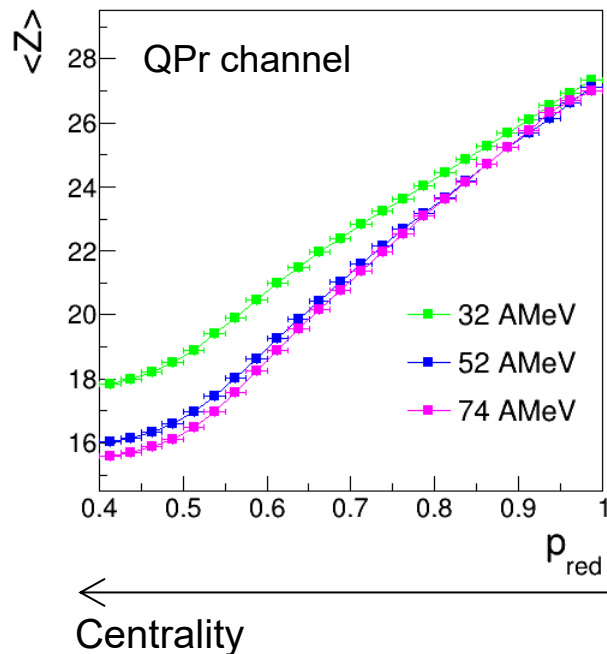
FAZIA

- 12 blocks covering from 2° to 14° .
- Charge discrimination up to uranium, isotopic discrimination up to $Z \sim 25$



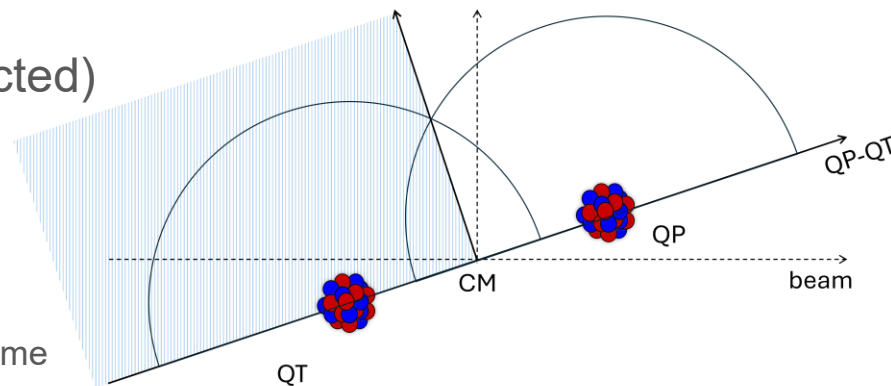
Secondary fragments

- **Quasi-projectile** (remnant or reconstructed)

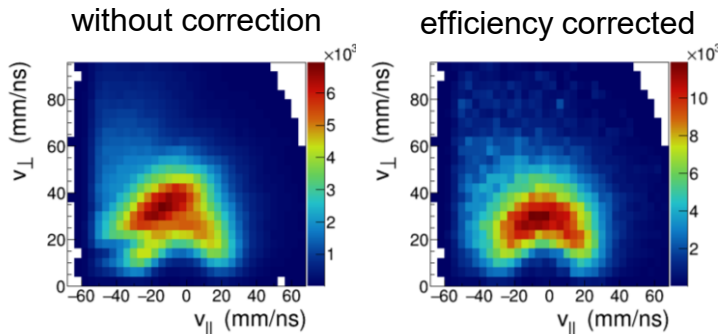


Secondary fragments

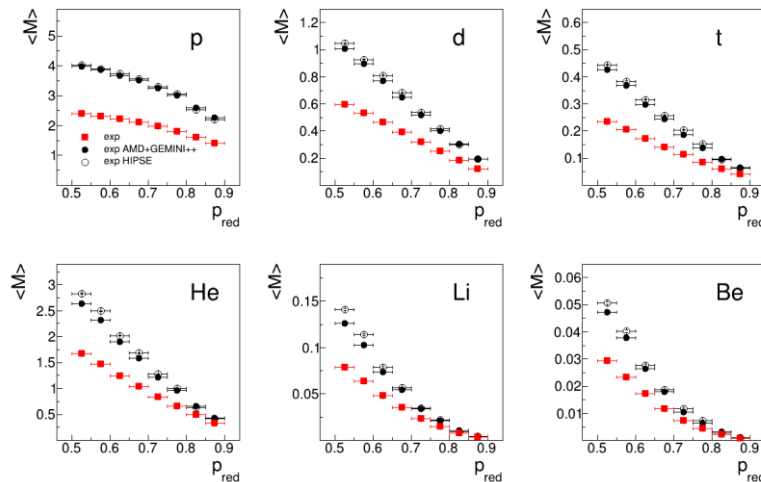
- Quasi-projectile (remnant or reconstructed)
- Light charged fragments ($Z < 5$)**
 - QP phase space:** selection of fragments (identified in mass) emitted in the forward hemisphere of the centre of mass reference frame (not-dashed region)
 - Efficiency correction:** obtained using AMD+GEMINI++ simulations



protons @ 74 AMeV
($p_{red} = 0.6-0.7$)

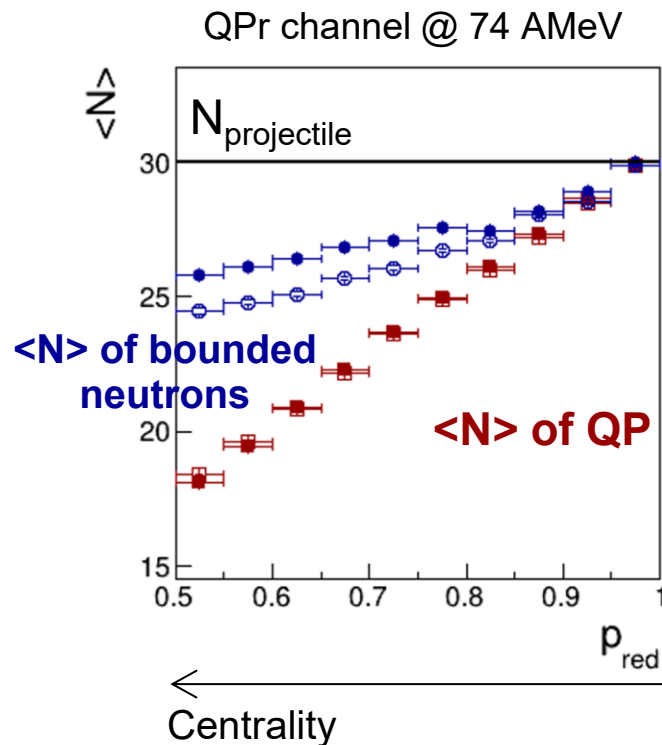


$^{58}\text{Ni} + ^{58}\text{Ni}$ @ 74 AMeV



Secondary fragments

- Quasi-projectile (remnant or reconstructed)
- Light charged fragments ($Z < 5$)
- **Free neutrons**
 - Not detected
 - The average number of free neutrons can be estimated as the difference between the N_{proj} and the $\langle N \rangle$ estimated considering the bounded neutrons.



Disentanglement between evaporative and midvelocity emissions

S. Piantelli, PRC 74, 034609 (2006)
A. Mangiarotti, PRL 93, 232701 (2004)

• Yields

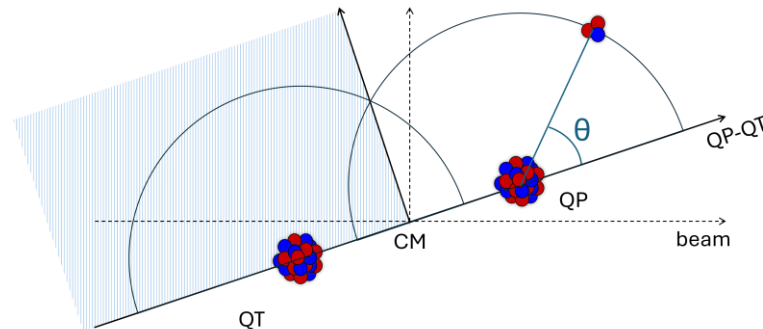
- **Light charged fragments:**
 - angular distributions of θ , i.e. the polar angle between the charged particle in the QP frame and the QP-QT separation axis
→ superposition of two emission components

- **statistical evaporation:**
approximately a $\sin(\theta)$ distribution
- large excess at backward angles which is ascribed to **midvelocity emission**.

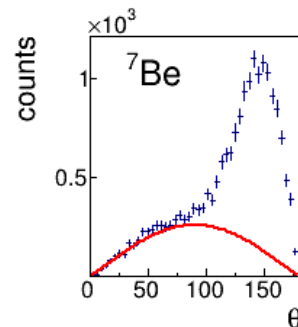
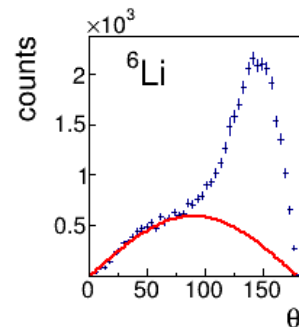
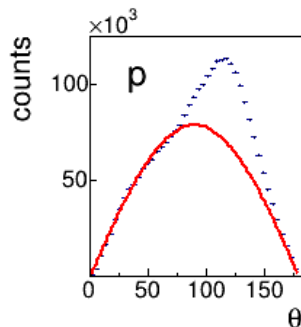
- **Free neutrons:**
 - hypothesis must be done

• Energies:

- Nuclear calorimetry technique



$^{58}\text{Ni}+^{58}\text{Ni}$ @ 74 AMeV



Primary source characteristics

QP* source

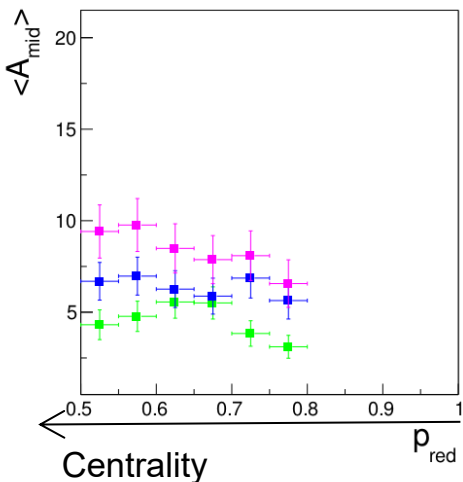
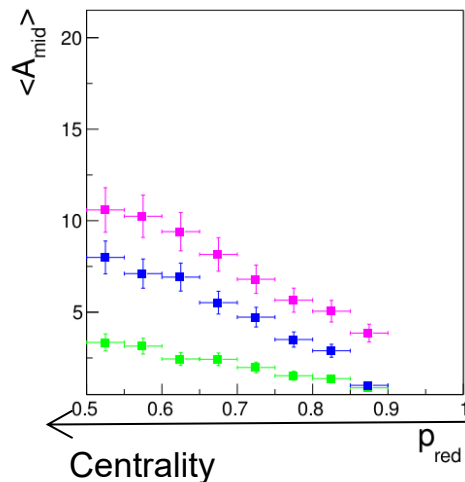
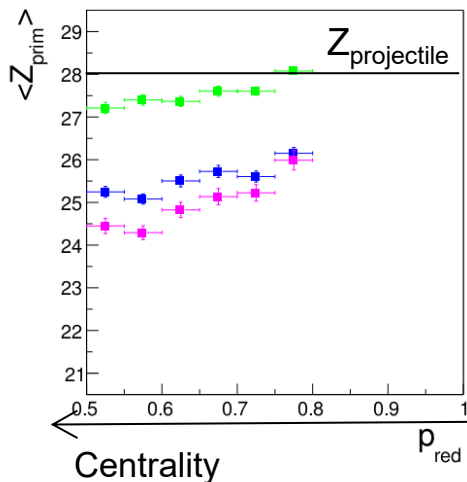
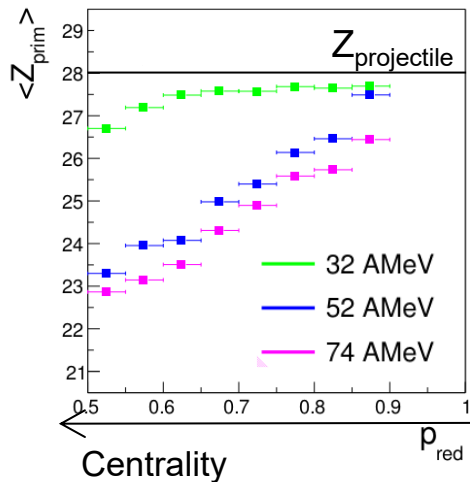
neck source

QPr channel

QPb channel

QPr channel

QPb channel

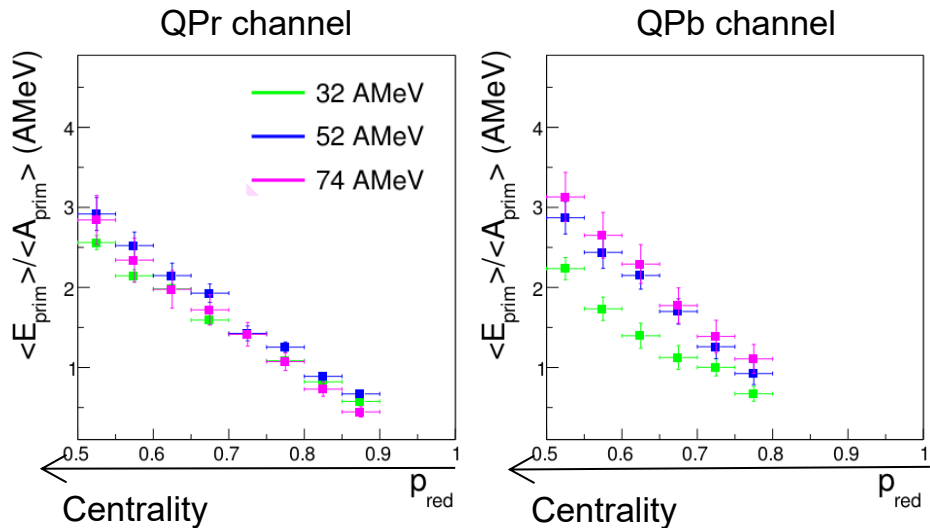


- the size of the QP* depends on both centrality and incident energy: the QP* remains close to the projectile at 32 AMeV, but becomes much lighter in central collisions at 52 and 74 AMeV

- the midvelocity contribution increases with centrality and clearly grows with beam energy

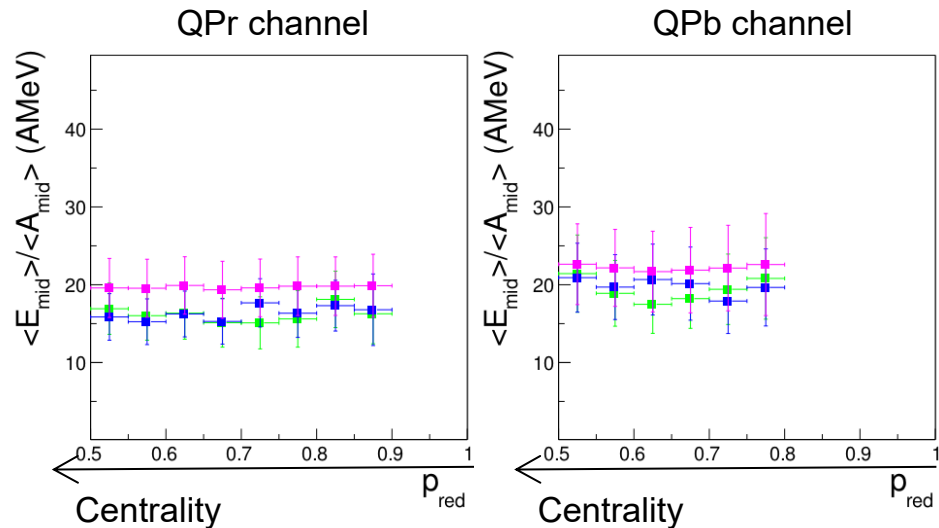
Primary source characteristics

QP* source



- the excitation energy per nucleon of the QP*, as expected for an equilibrated source, remains below 3 AMeV and increases with centrality

neck source

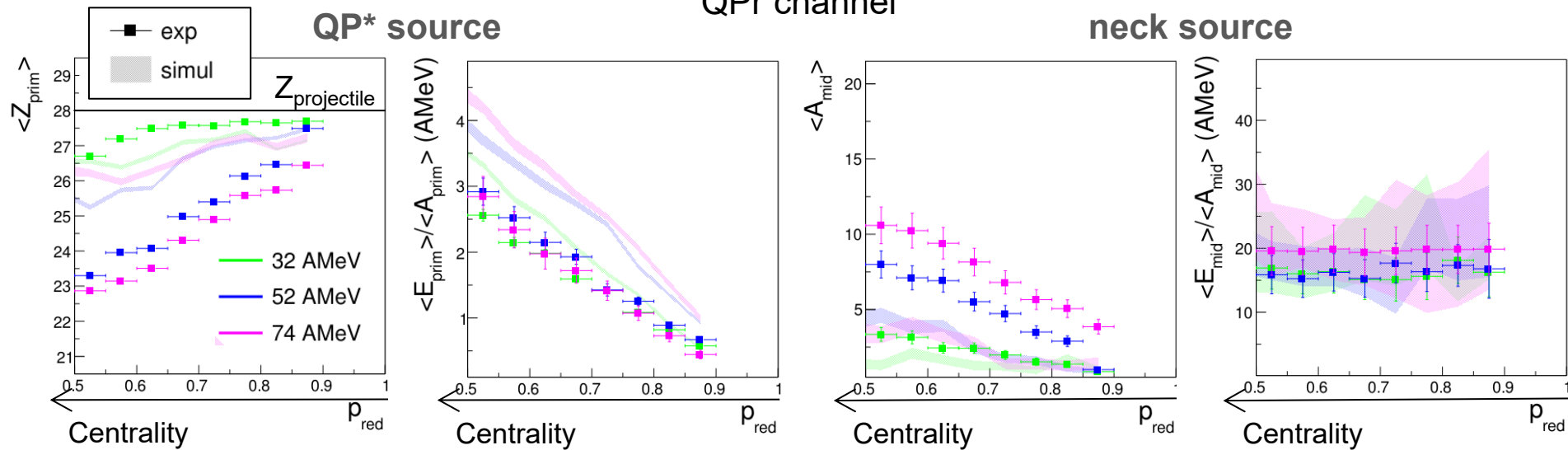


- the energy density of the neck source is significantly higher (~ 15 - 20 AMeV) and shows no dependence on either centrality or beam energy.

Primary source characteristics

Comparison with the AMD+GEMINI++ simulations

QPr channel



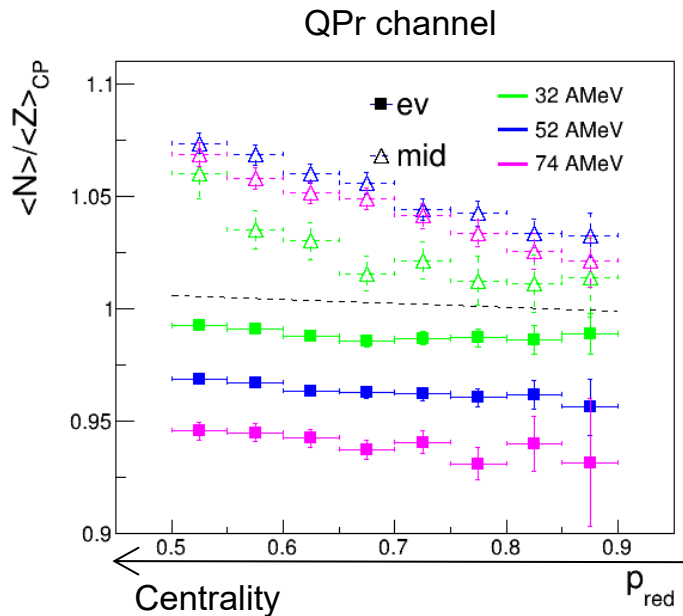
- overall satisfactory reproduction of the secondary fragment characteristics
- differences in the properties of the primary sources: underestimation of the midvelocity contribution
- Within AMD, the properties of the primary QP* source can be independently evaluated as a function of time \rightarrow the proposed reconstruction procedure allows to approach the characteristics of the fragments at the end of the dynamical stage (150-300 fm/c).

Isospin drift

isospin ratio for complex particles: $\langle N \rangle / \langle Z \rangle_{CP} = \frac{\sum_i \sum_v N_v^i}{\sum_i \sum_v Z_v^i},$

E. Galichet, PRC 79, 064614 (2009)

- where the v index numbers the different complex particles in the ith event, and the outer sum runs over all the events in the selected pred bin.
- d, t, ^3He , ^4He , ^6He , ^6Li , ^7Li , ^8Li , ^9Li , ^7Be , ^9Be , and ^{10}Be are taken into account as complex particles
- free protons and neutrons are excluded



- the **neutron enrichment of the neck** is clearly observed
- midvelocity emission: the N/Z is essentially independent on the beam energy.
- evaporative component: an interesting dependence on the beam energy emerges.
 - different isospin content of the QP* source which reflects on the evaporated particles
 - as the energy increases, the growing midvelocity contribution leaves the QP* slightly more neutron poor
- Since the neutron enrichment of the midvelocity component weakly depends on the beam energy, this effect is mainly driven by the increasing abundance of midvelocity yields.

Conclusions

- Experimental access to the dynamical phase (150-300 fm/c) in both the QPr and QPb channels
- Similar results in the two channels \rightarrow following evolution of the QP* largely independent of its size and excitation energy.
- Primary source reconstruction
 - The excitation energy of the primary QP* is below 3 AMeV
 - The midvelocity contribution increases with the incident energy
 - The energy density of the neck source is 15-20 AMeV \rightarrow dynamical origin from a hot overlapping region
 - The AMD+GEMINI++ simulations underestimate the midvelocity contribution
- Isospin Drift
 - The neutron enrichment of the neck is clearly observed
 - The midvelocity emission plays a key role in determining the isospin characteristics of the primary QP*, mainly driven by the increasing amount of neck emissions.

Thank you for your attention



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