

Coherent photoproduction of J/ψ at forward rapidity in Pb–Pb with the ALICE detector

Guillermo Contreras
on behalf of the **ALICE** Collaboration

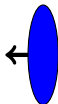
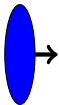
Faculty of Nuclear Sciences and Physical Engineering,
Czech Technical University in Prague, Prague, Czech Republic

December 3rd, 2012

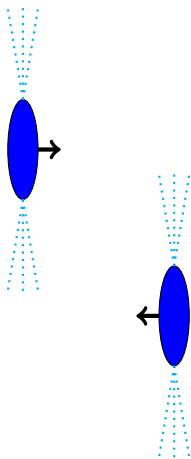
- ✓ The LHC as a γPb collider
- ✓ ALICE and J/ψ
- ✓ The measurement
(ALICE Collaboration, arXiv:1209.3715. Accepted by PLB.)
- ✓ Comparison to theory
- ✓ Of things to come
- ✓ Summary

The LHC as γ Pb collider

✓ At the LHC heavy ions are accelerated towards each other at ultra relativistic energies



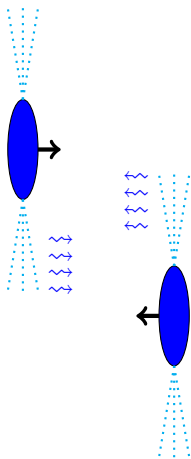
The LHC as γ Pb collider



✓ At the LHC heavy ions are accelerated towards each other at ultra relativistic energies

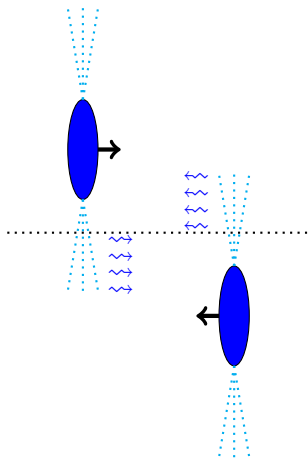
✓ Being charged particles, they are accompanied by an electromagnetic field

The LHC as γ Pb collider



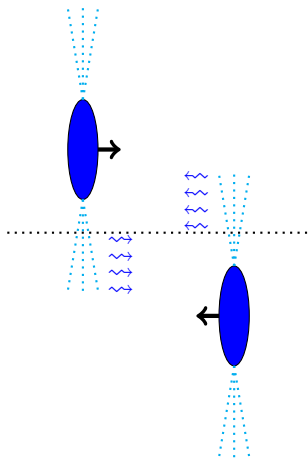
- ✓ At the LHC heavy ions are accelerated towards each other at ultra relativistic energies
- ✓ Being charged particles, they are accompanied by an electromagnetic field
- ✓ According to a proposal from Fermi, this field can be interpreted as a beam of quasi real photons

The LHC as γ Pb collider



- ✓ At the LHC heavy ions are accelerated towards each other at ultra relativistic energies
- ✓ Being charged particles, they are accompanied by an electromagnetic field
- ✓ According to a proposal from Fermi, this field can be interpreted as a beam of quasi real photons
- ✓ The strong force has a short range, so for impact parameters bigger than the sum of the radii of the incoming ions, the only possible interaction is electromagnetic.

The LHC as γ Pb collider



- ✓ At the LHC heavy ions are accelerated towards each other at ultra relativistic energies
- ✓ Being charged particles, they are accompanied by an electromagnetic field
- ✓ According to a proposal from Fermi, this field can be interpreted as a beam of quasi real photons
- ✓ The strong force has a short range, so for impact parameters bigger than the sum of the radii of the incoming ions, the only possible interaction is electromagnetic.

These processes are called **ultra peripheral collisions (UPC)**

The flux of photons

- ✓ The flux of photons is proportional to Z^2 , making heavy ions an intense source of photons
- ✓ The photons are emitted coherently by the source so the wavelength is restricted to be bigger than the nuclear radius (R_A)

The flux of photons

✓ The flux of photons is proportional to Z^2 , making heavy ions an intense source of photons

✓ The photons are emitted coherently by the source so the wavelength is restricted to be bigger than the nuclear radius (R_A)

- In the transverse direction

$$p_T \lesssim \hbar c / R_A \approx 30 \text{ MeV}$$

- The photon energy k reaches,

$$k \lesssim \hbar c \gamma_L / R_A \approx 40 \text{ GeV}$$

where γ_L is the Lorentz factor to the lab sys for the LHC energy of the 2011 Pb–Pb run

The flux of photons

✓ The flux of photons is proportional to Z^2 , making heavy ions an intense source of photons

✓ The photons are emitted coherently by the source so the wavelength is restricted to be bigger than the nuclear radius (R_A)

- In the transverse direction

$$p_T \lesssim \hbar c / R_A \approx 30 \text{ MeV}$$

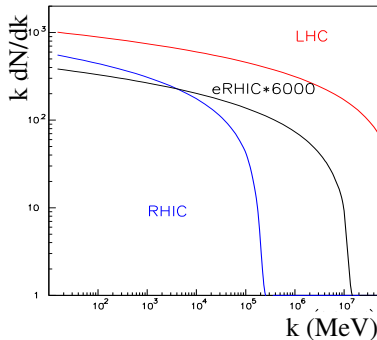
- The photon energy k reaches,

$$k \lesssim \hbar c \gamma_L / R_A \approx 40 \text{ GeV}$$

where γ_L is the Lorentz factor to the lab sys for the LHC energy of the 2011 Pb–Pb run

✓ The flux decays with photon energy. For the 2011 Pb–Pb run:

$$\sqrt{s_{\gamma A}^{\max}} \approx 450 \text{ GeV}$$



Photon flux as a function of the photon energy k in the rest frame of the target nucleus;

Fig. 4: Phys. Rep. **458** (2008) 1–171.

Coherent photo-production of J/ψ

- ✓ The photon interacts **coherently** with the opposite nucleus and produces a single vector meson
- ✓ If the photon interacts with only one nucleon, the process is called **incoherent**
- ✓ There are two contributions, one for each incoming nucleus acting as a source of the photons and the other as the target

Coherent photo-production of J/ψ

- ✓ The photon interacts **coherently** with the opposite nucleus and produces a single vector meson
- ✓ If the photon interacts with only one nucleon, the process is called **incoherent**
- ✓ There are two contributions, one for each incoming nucleus acting as a source of the photons and the other as the target
- ✓ In the case we are interested here

$$Pb\ Pb \rightarrow Pb\ (\gamma\ Pb) \rightarrow Pb\ Pb\ J/\psi$$

with the J/ψ produced at forward rapidity y

Coherent photo-production of J/ψ

✓ The photon interacts **coherently** with the opposite nucleus and produces a single vector meson

✓ If the photon interacts with only one nucleon, the process is called **incoherent**

✓ There are two contributions, one for each incoming nucleus acting as a source of the photons and the other as the target

✓ In the case we are interested here

$$Pb\ Pb \rightarrow Pb\ (\gamma\ Pb) \rightarrow Pb\ Pb\ J/\psi$$

with the J/ψ produced at forward rapidity y

✓ The kinematics of the J/ψ are

- $p_T < 30\ \text{MeV}$ ($< 250\ \text{MeV}$ for incoherent processes)
- $y = \ln(2k/M)$ where M is the mass of the J/ψ

Coherent photo-production of J/ψ

- ✓ The photon interacts **coherently** with the opposite nucleus and produces a single vector meson
- ✓ If the photon interacts with only one nucleon, the process is called **incoherent**
- ✓ There are two contributions, one for each incoming nucleus acting as a source of the photons and the other as the target
- ✓ In the case we are interested here

$$Pb Pb \rightarrow Pb (\gamma Pb) \rightarrow Pb Pb J/\psi$$

with the J/ψ produced at forward rapidity y

- ✓ The kinematics of the J/ψ are
 - $p_T < 30$ MeV (< 250 MeV for incoherent processes)
 - $y = \ln(2k/M)$ where M is the mass of the J/ψ
- ✓ Some representative values for the γA cms energy $W_{\gamma A}$:

$$y = 0 \implies W_{\gamma A} \approx 90 \text{ GeV}$$

$$y = -3 \implies W_{\gamma A} \approx \frac{400}{20} \text{ GeV}$$

Coherent photo-production of J/ψ

✓ The photon interacts **coherently** with the opposite nucleus and produces a single vector meson

✓ If the photon interacts with only one nucleon, the process is called **incoherent**

✓ There are two contributions, one for each incoming nucleus acting as a source of the photons and the other as the target

✓ In the case we are interested here

$$Pb Pb \rightarrow Pb (\gamma Pb) \rightarrow Pb Pb J/\psi$$

with the J/ψ produced at forward rapidity y

✓ The kinematics of the J/ψ are

- $p_T < 30$ MeV (< 250 MeV for incoherent processes)
- $y = \ln(2k/M)$ where M is the mass of the J/ψ

✓ Some representative values for the γA cms energy $W_{\gamma A}$:

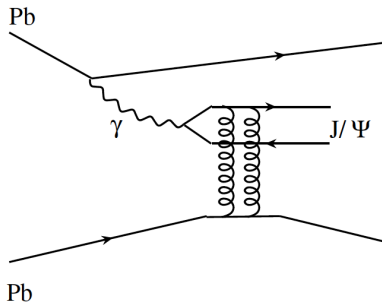
$$y = 0 \implies W_{\gamma A} \approx 90 \text{ GeV}$$

$$y = -3 \implies W_{\gamma A} \approx \frac{400}{20} \text{ GeV}$$

✓ For $y = -3$ the flux of photons at $W_{\gamma A} \approx 400$ GeV is around 4% of the flux at $W_{\gamma A} \approx 20$ GeV

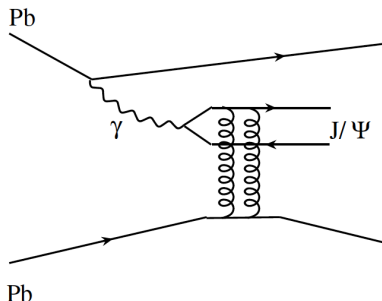
Coherent photo-production of J/ψ in QCD

In QCD models this diffractive process is mediated by a two gluon exchange:



Coherent photo-production of J/ψ in QCD

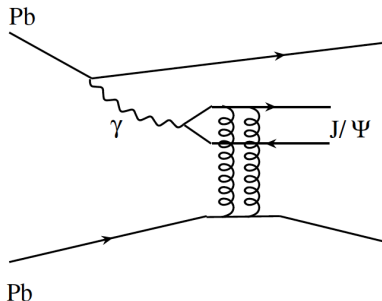
In **QCD** models this diffractive process is mediated by a two gluon exchange:



In these models, at LO, the cross section depends on the **square** of the gluon distribution in the target

Coherent photo-production of J/ψ in QCD

In QCD models this diffractive process is mediated by a two gluon exchange:



In these models, at LO, the cross section depends on the **square** of the gluon distribution in the target

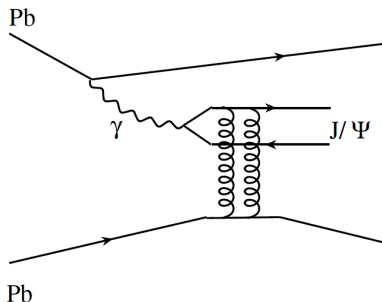
✓ In this case x -Bjorken is related to the rapidity of the J/ψ :

$$x = \frac{M}{\sqrt{s_{NN}}} e^{\pm y}$$

where the \pm refers to the two possible contributions

Coherent photo-production of J/ψ in QCD

In QCD models this diffractive process is mediated by a two gluon exchange:



In these models, at LO, the cross section depends on the **square** of the gluon distribution in the target

✓ In this case x -Bjorken is related to the rapidity of the J/ψ :

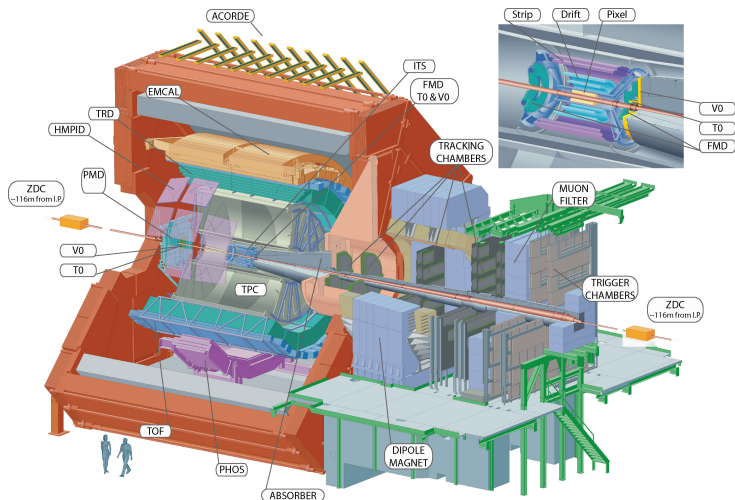
$$x = \frac{M}{\sqrt{s_{NN}}} e^{\pm y}$$

where the \pm refers to the two possible contributions

✓ At $y = 0$, values of $x \approx 10^{-3}$ are probed; while at $y = -3$ the process probes $x \approx 10^{-2}$ or $x \approx 10^{-5}$

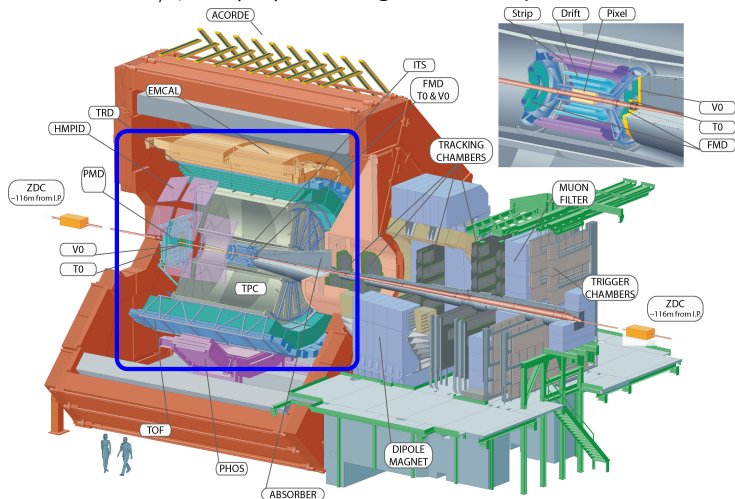
✓ In this scenario the nuclear gluon distribution and its shadowing can be probed as a function of x varying y

Excellent tracking and PID capabilities



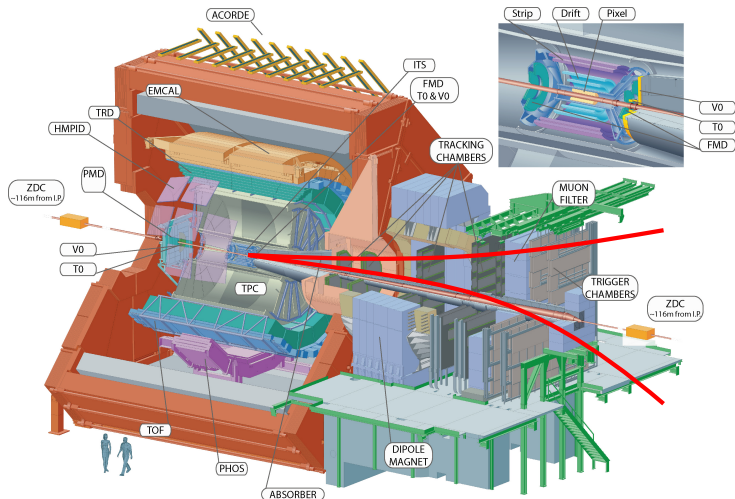
ALICE and J/ψ

At **mid rapidities** ALICE measures J/ψ in the lepton channels,
 $J/\psi \rightarrow e^+e^-$ and $J/\psi \rightarrow \mu^+\mu^-$, using its PID capabilities

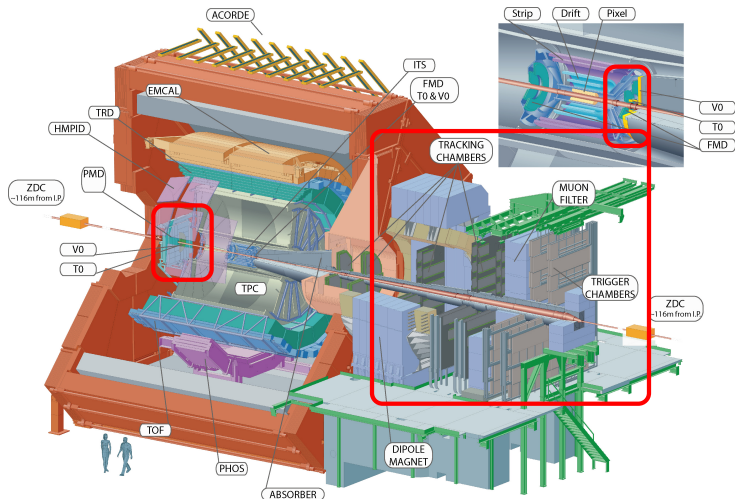


ALICE and J/ψ

At **forward rapidities** ALICE measures $J/\psi \rightarrow \mu^+ \mu^-$

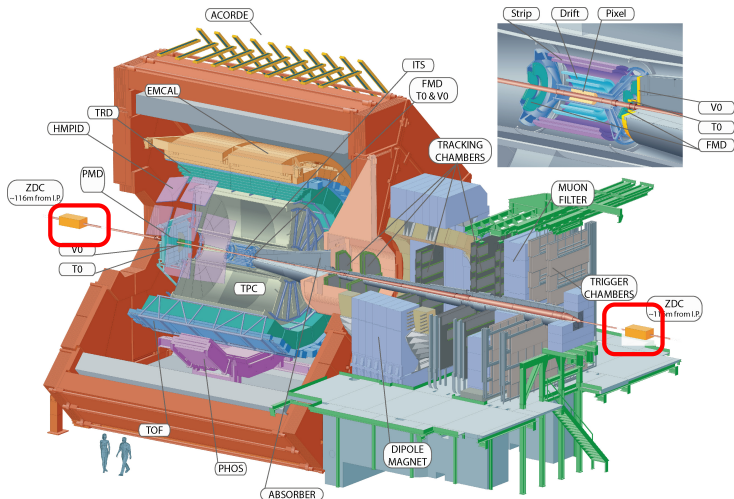


The **trigger** uses the muon spectrometer and the VZERO



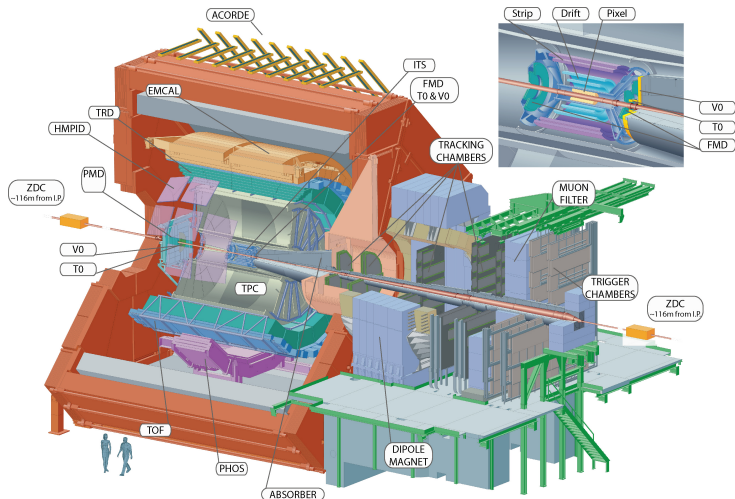
ALICE and J/ψ

The **analysis** also uses the ZDC at ± 116 m from the IP



ALICE and J/ψ

ALICE can measure the J/ψ **down to zero p_T** in both y ranges



Event selection

Using $55 \mu\text{b}^{-1}$ collected in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ during 2011

Event selection

Using $55 \mu\text{b}^{-1}$ collected in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ during 2011

✓ Transverse momentum of J/ψ less than 0.3 GeV

This implies that the transverse momentum of the muons is around 1.5 GeV each and that they are back-to-back

Event selection

Using $55 \mu\text{b}^{-1}$ collected in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ during 2011

✓ Transverse momentum of J/ψ less than 0.3 GeV

This implies that the transverse momentum of the muons is around 1.5 GeV each and that they are back-to-back

✓ Rapidity of the J/ψ in range $(-3.6, -2.6)$

Given by the overlap of the acceptances of the VZERO–C and the muon spectrometer

Event selection

Using $55 \mu\text{b}^{-1}$ collected in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ during 2011

✓ Transverse momentum of J/ψ less than 0.3 GeV

This implies that the transverse momentum of the muons is around 1.5 GeV each and that they are back-to-back

✓ Rapidity of the J/ψ in range (-3.6,-2.6)

Given by the overlap of the acceptances of the VZERO–C and the muon spectrometer

✓ Beam gas rejection

VZERO timing and (p,DCA) correlation of the muon tracks

Event selection

Using $55 \mu\text{b}^{-1}$ collected in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ during 2011

✓ Transverse momentum of J/ψ less than 0.3 GeV

This implies that the transverse momentum of the muons is around 1.5 GeV each and that they are back-to-back

✓ Rapidity of the J/ψ in range (-3.6,-2.6)

Given by the overlap of the acceptances of the VZERO–C and the muon spectrometer

✓ Beam gas rejection

VZERO timing and (p,DCA) correlation of the muon tracks

✓ Rejection of low multiplicity hadronic interactions

Little energy in the neutron ZDC; i. e., no spectators

Event selection

Using $55 \mu\text{b}^{-1}$ collected in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ during 2011

✓ Transverse momentum of J/ψ less than 0.3 GeV

This implies that the transverse momentum of the muons is around 1.5 GeV each and that they are back-to-back

✓ Rapidity of the J/ψ in range (-3.6,-2.6)

Given by the overlap of the acceptances of the VZERO–C and the muon spectrometer

✓ Beam gas rejection

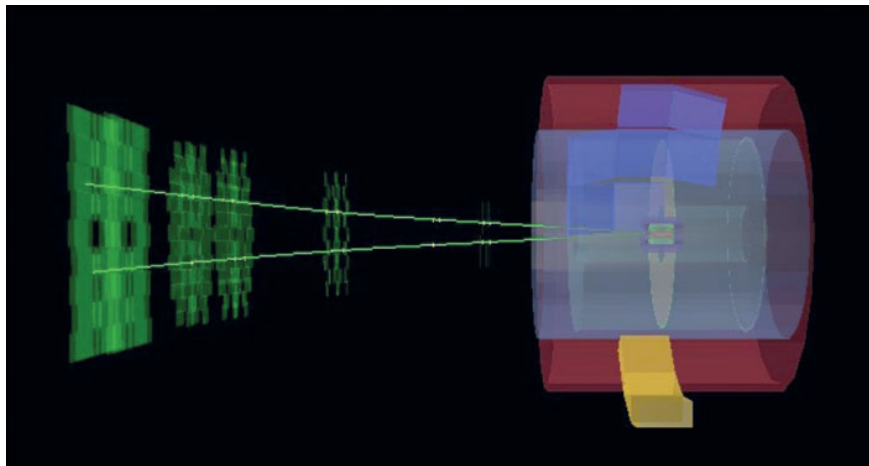
VZERO timing and (p,DCA) correlation of the muon tracks

✓ Rejection of low multiplicity hadronic interactions

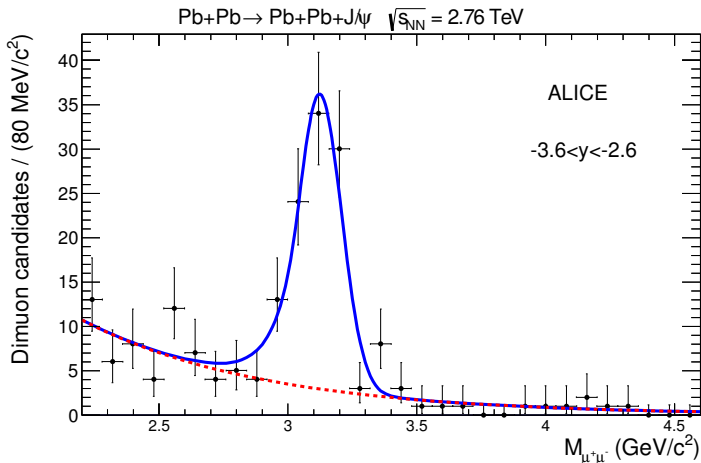
Little energy in the neutron ZDC; i. e., no spectators

117 events remained with $2.8 < M(\mu^+, \mu^-) < 3.4 \text{ GeV}$

One typical event



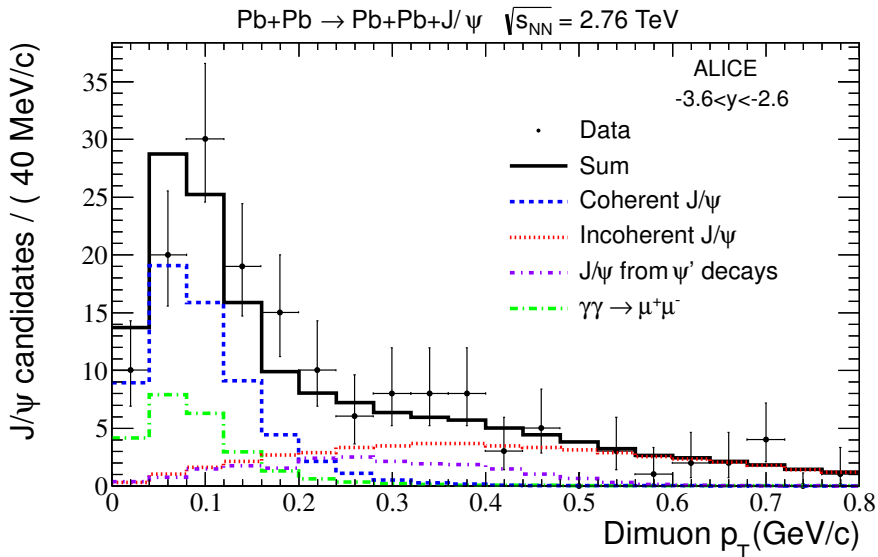
Mass distribution



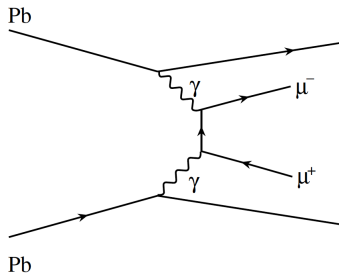
- ✓ Like-sign events: 2
- ✓ No activity somewhere else
- ✓ Signal fitted to a CB shape

- ✓ Background to an exponential
- ✓ Exponential shape compatible with expectations from $\gamma\gamma \rightarrow \mu^+\mu^-$

Contributions around the mass peak



Normalization

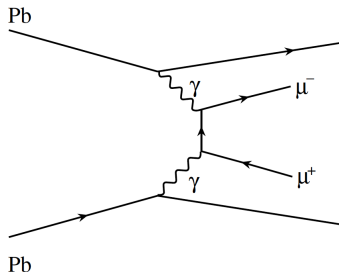


✓ Process $\gamma\gamma \rightarrow \mu^+\mu^-$ used:

$$\frac{d\sigma_{\text{coh}}}{dy} = \frac{1}{BR} \cdot \frac{N_{\text{coh}}}{N_{\gamma\gamma}} \cdot \frac{(\text{Acc x } \epsilon)_{\gamma\gamma}}{(\text{Acc x } \epsilon)_{\text{coh}}} \frac{\sigma_{\gamma\gamma}}{\Delta y}$$

✓ Standard QED process ...

Normalization



✓ Process $\gamma\gamma \rightarrow \mu^+\mu^-$ used:

$$\frac{d\sigma_{\text{coh}}}{dy} = \frac{1}{BR} \cdot \frac{N_{\text{coh}}}{N_{\gamma\gamma}} \cdot \frac{(\text{Acc x } \epsilon)_{\gamma\gamma}}{(\text{Acc x } \epsilon)_{\text{coh}}} \frac{\sigma_{\gamma\gamma}}{\Delta y}$$

✓ Standard QED process ... **but**

- ✓ Uncertainty in higher order terms due to coupling $\sim Z\sqrt{\alpha}$
- ✓ Uncertainty on minimum momentum transfer and nuclear form factor
- ✓ Previous experimental results from RHIC also have large uncertainties and can not constraint the theory

Cross section for the photoproduction of coherent J/ψ at forward rapidities in Pb–Pb with the ALICE detector

✓ Measurement at $\sqrt{s_{NN}} = 2.76$ TeV for $p_T < 0.3$ GeV and $-3.6 < y < -2.6$

Cross section for the photoproduction of coherent J/ψ at forward rapidities in Pb–Pb with the ALICE detector

- ✓ Measurement at $\sqrt{s_{NN}} = 2.76$ TeV for $p_T < 0.3$ GeV and $-3.6 < y < -2.6$
- ✓ Systematic uncertainty: **+24/-26 %**
 - Dominant contributions.
Theoretical uncertainty on $\sigma_{\gamma\gamma}$: **20 %**, signal extraction: **+9/-14 %**
 - Intermediate contributions.
Reconstruction efficiency: **6 %**, muon trigger efficiency: **5 %**, acceptance calculation **3 %**
 - Small contributions.
 $\gamma\gamma \rightarrow e^+e^-$: **2 %**, branching ratio: **1 %**

Cross section for the photoproduction of coherent J/ψ at forward rapidities in Pb–Pb with the ALICE detector

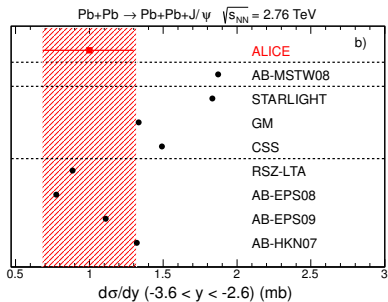
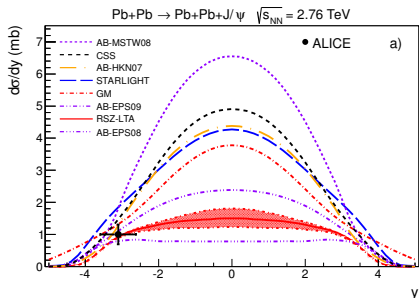
- ✓ Measurement at $\sqrt{s_{NN}} = 2.76$ TeV for $p_T < 0.3$ GeV and $-3.6 < y < -2.6$
- ✓ Systematic uncertainty: **+24/-26 %**
 - Dominant contributions.
Theoretical uncertainty on $\sigma_{\gamma\gamma}$: **20 %**, signal extraction: **+9/-14 %**
 - Intermediate contributions.
Reconstruction efficiency: **6 %**, muon trigger efficiency: **5 %**, acceptance calculation **3 %**
 - Small contributions.
 $\gamma\gamma \rightarrow e^+e^-$: **2 %**, branching ratio: **1 %**

$$\frac{d\sigma_{\text{coh}}}{dy} = 1.00 \pm 0.18 \text{ (stat)} \begin{matrix} +0.24 \\ -0.26 \end{matrix} \text{ (sys)} \text{ mb}$$

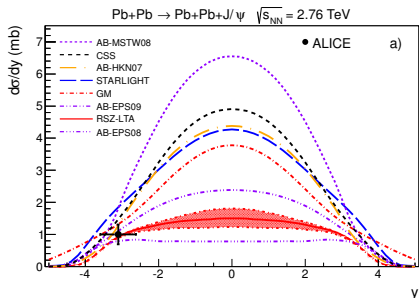
Theory predictions

- ✓ The total cross section, $\sigma(PbPb)$, is a convolution of the photon flux and the γA cross section, $\sigma(\gamma A)$
- ✓ Some differences in the treatment of the photon flux in the different calculations, but numerically the results are similar
- ✓ The main differences among the predictions come from the assumptions in the computation of $\sigma(\gamma A)$
- ✓ Active field, **5 recent predictions**:
 - S.R. Klein, J. Nystrand, Phys. Rev. C 60 (1999) 014903.
 - V.P. Goncalves, M.V.T. Machado, Phys. Rev. C 84 (2011) 011902.
 - V. Rebyakova, M. Strikman, M. Zhalov, Phys. Lett. B 710 (2012) 647.
 - A. Adeluyi, C.A. Bertulani, Phys. Rev. C 85 (2012) 044904.
 - A. Cisek, W. Schafer, A. Szczurek, Phys. Rev. C 86 (2012) 014905.

Comparison with Theory



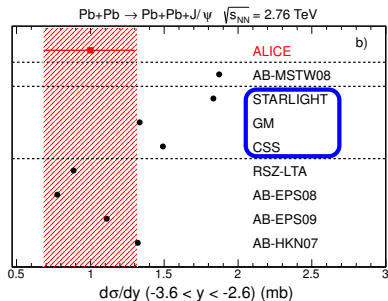
Comparison with Theory



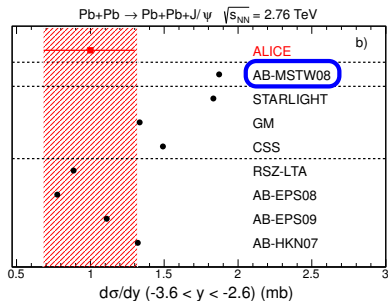
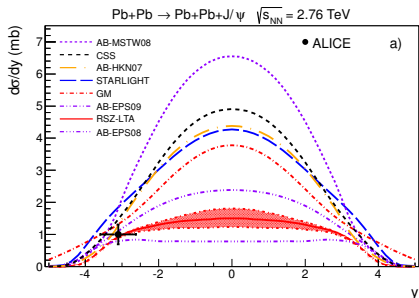
✓ STARLIGHT, from Klein and Nystrand, uses a GVDM coupled to a Glauber approach to link the γA to the γp cross section, where the later is obtained from a parameterization of HERA data

✓ GM is based on the color dipole model, where the scattering amplitude depends on the nuclear profile and the dipole nucleon cross section, which is taken from the IIM model which incorporates saturation

✓ The CSS model uses a Glauber approach and the color dipole nucleon amplitude based on the unintegrated gluon distribution of the proton



Comparison with Theory



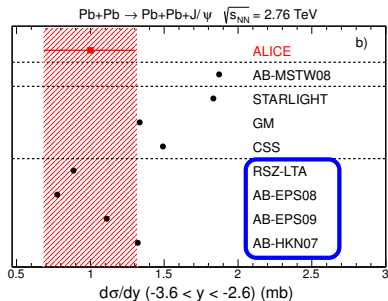
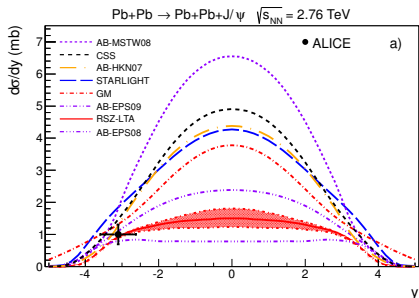
✓ STARLIGHT, from Klein and Nystrand, uses a GVDM coupled to a Glauber approach to link the γA to the γp cross section, where the later is obtained from a parameterization of HERA data

✓ GM is based on the color dipole model, where the scattering amplitude depends on the nuclear profile and the dipole nucleon cross section, which is taken from the IIM model which incorporates saturation

✓ The CSS model uses a Glauber approach and the color dipole nucleon amplitude based on the unintegrated gluon distribution of the proton

✓ The AB models use the LO pQCD amplitude scaled to correct for some other effects. For the gluon distribution, AB-MSTW08 assumes no nuclear effects

Comparison with Theory



✓ STARLIGHT, from Klein and Nystrand, uses a GVDM coupled to a Glauber approach to link the γA to the γp cross section, where the later is obtained from a parameterization of HERA data

✓ GM is based on the color dipole model, where the scattering amplitude depends on the nuclear profile and the dipole nucleon cross section, which is taken from the IIM model which incorporates saturation

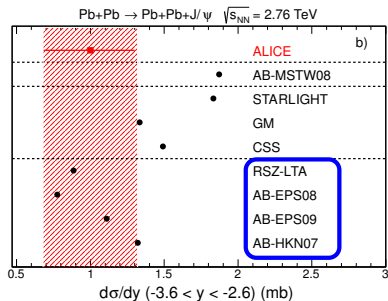
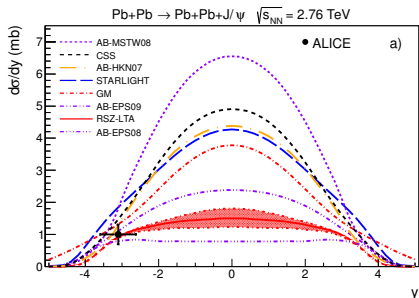
✓ The CSS model uses a Glauber approach and the color dipole nucleon amplitude based on the unintegrated gluon distribution of the proton

✓ The AB models use the LO pQCD amplitude scaled to correct for some other effects. For the gluon distribution, AB-MSTW08 assumes no nuclear effects

✓ The other AB models incorporate nuclear effects according to the EPS08, EPS09 or HKN07 prescriptions

✓ RSZ-LTA is based on the LO pQCD amplitude for two gluon exchange where the nuclear gluon density incorporates shadowing computed in the leading twist approximation

Comparison with Theory



✓ STARLIGHT, from Klein and Nystrand, uses a GVDM coupled to a Glauber approach to link the γA to the γp cross section, where the later is obtained from a parameterization of HERA data

✓ GM is based on the color dipole model, where the scattering amplitude depends on the nuclear profile and the dipole nucleon cross section, which is taken from the IIM model which incorporates saturation

✓ The CSS model uses a Glauber approach and the color dipole nucleon amplitude based on the unintegrated gluon distribution of the proton

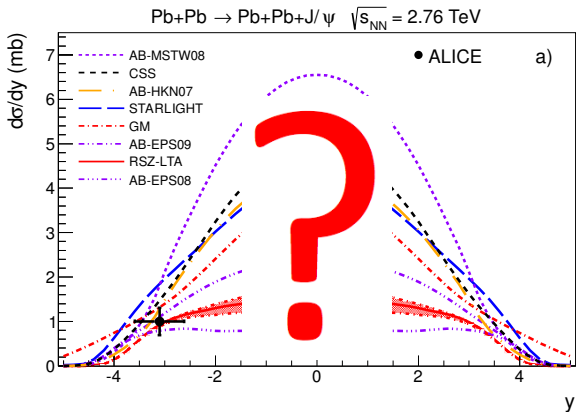
✓ The AB models use the LO pQCD amplitude scaled to correct for some other effects. For the gluon distribution, AB-MSTW08 assumes no nuclear effects

✓ The other AB models incorporate nuclear effects according to the EPS08, EPS09 or HKN07 prescriptions

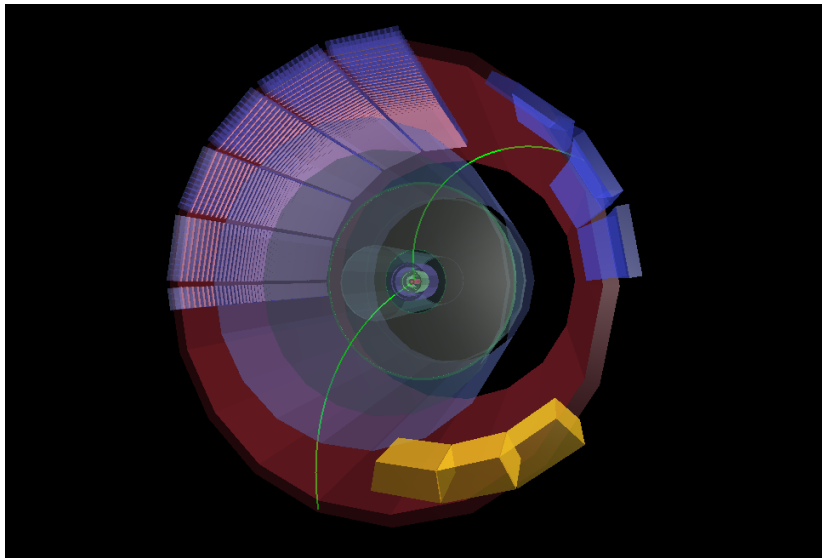
✓ RSZ-LTA is based on the LO pQCD amplitude for two gluon exchange where the nuclear gluon density incorporates shadowing computed in the leading twist approximation

Models which incorporate shadowing – with the scale fixed near the mass of the J/ ψ – are closer to data

Of things to come: The natural question

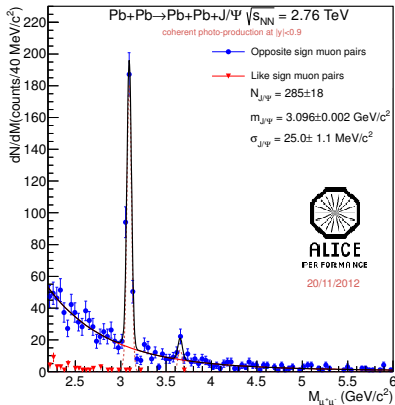


ALICE has the answer

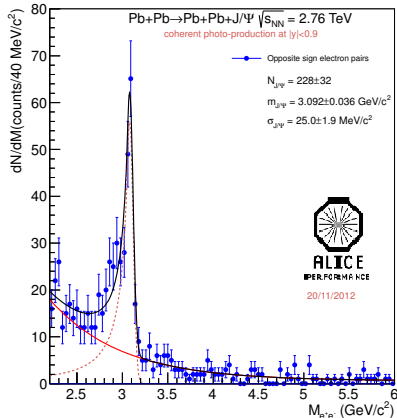


The answer is almost available

✓ Using the PID capabilities of ALICE in the central region, it is possible to measure separately $J/\psi \rightarrow \mu^+ \mu^-$ and $J/\psi \rightarrow e^+ e^-$



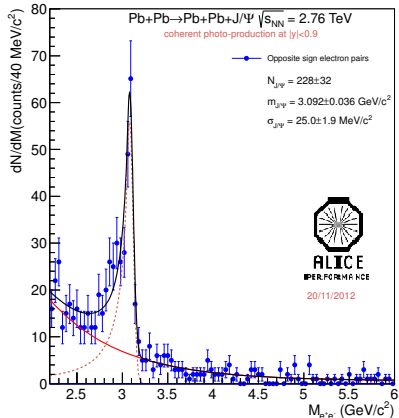
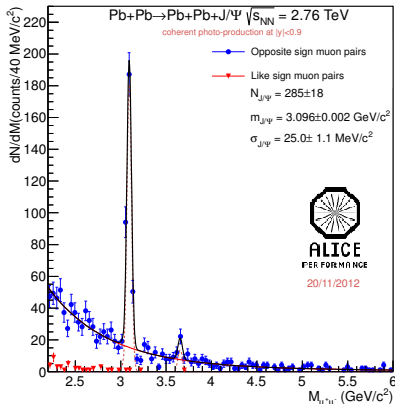
ALI-CONF-45234



ALI-CONF-45238

The answer is almost available

✓ Using the PID capabilities of ALICE in the central region, it is possible to measure separately $J/\psi \rightarrow \mu^+ \mu^-$ and $J/\psi \rightarrow e^+ e^-$



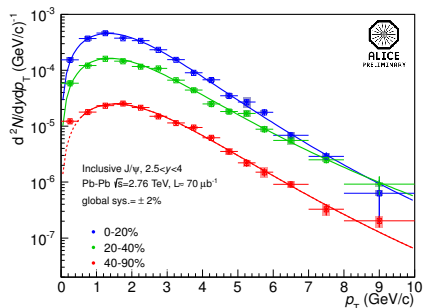
Cross section at $y = 0$ will be ready soon!

Of things to come: UPC in PC?

- ✓ In principle one could have the same mechanism for peripheral interactions, but overlapped with the accompanying hadronic interaction of the nuclei
- ✓ No way to disentangle the coherent signal ...

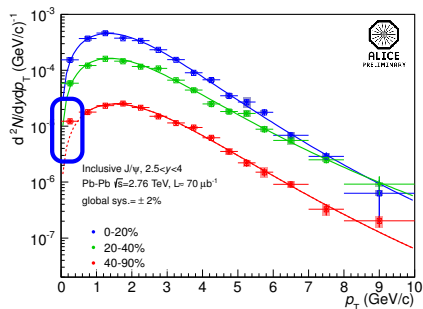
Of things to come: UPC in PC?

- ✓ In principle one could have the same mechanism for peripheral interactions, but overlapped with the accompanying hadronic interaction of the nuclei
- ✓ No way to disentangle the coherent signal ...



Of things to come: UPC in PC?

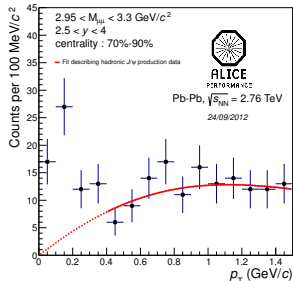
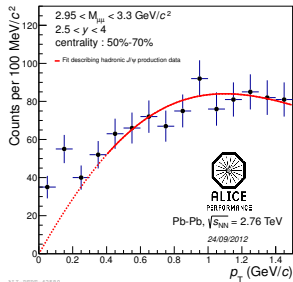
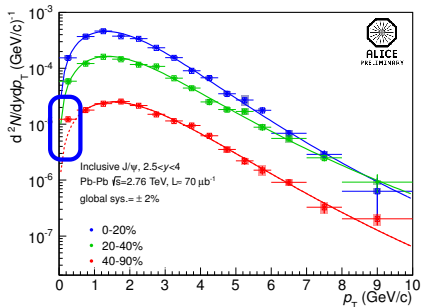
- ✓ In principle one could have the same mechanism for peripheral interactions, but overlapped with the accompanying hadronic interaction of the nuclei
- ✓ No way to disentangle the coherent signal ... **right?**



ALI-PREL-36165

Of things to come: UPC in PC?

- ✓ In principle one could have the same mechanism for peripheral interactions, but overlapped with the accompanying hadronic interaction of the nuclei
- ✓ No way to disentangle the coherent signal ... **right?**



Of things to come: pA collisions at the LHC

- ✓ In January 2013 – just one month and a half to go – the LHC will provide both pA and Ap collisions
- ✓ As the system is asymmetric it would be possible to disentangle the source and the target at forward rapidities
- ✓ The most probable source will be the lead ion
- ✓ With the planned luminosities it could be possible to access the gluon in the proton down to $x \approx 10^{-5}$
- ✓ More difficult to access the low x nuclear gluon
- ✓ Anyway, exciting times just around the corner!

- ✓ Coherent production is also possible for other vector mesons
- ✓ Do not miss Kyrre Skjerdal's talk on Wednesday at 4:00 pm

Photoproduction of ρ_0 in Ultra-Peripheral Nuclear Collisions at ALICE

Summary

- ✓ The LHC can also be seen as a γA collider
- ✓ ALICE has measured the exclusive coherent photo-production of J/ψ at forward rapidities, $-3.6 < y < -2.6$:

$$\frac{d\sigma_{\text{coh}}}{dy} = 1.00 \pm 0.18 \text{ (stat)} \begin{matrix} +0.24 \\ -0.26 \end{matrix} \text{ (sys)} \text{ mb}$$

- ✓ Five recent theory predictions have been compared to the measurement
- ✓ Closer to the data are predictions incorporating nuclear gluon shadowing
- ✓ The measurement at $y = 0$ will be available soon
- ✓ As well as measurements of coherent ρ photoproduction
- ✓ And the pA LHC data is about to arrive
- ✓ **Stay tuned!**