

Heavy quark collisional energy loss

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– KRUGER 2012 –

PRESCRIPTIONS

PHARMACIST ON DUTY

PETER HIGGS

PETER HIGGS

BUY 1 GET 1

Heavy ion collisions



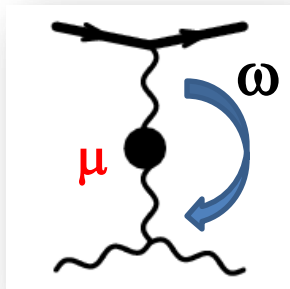
Quark-gluon plasma signals

- **soft probes** ○ ...
 - ...

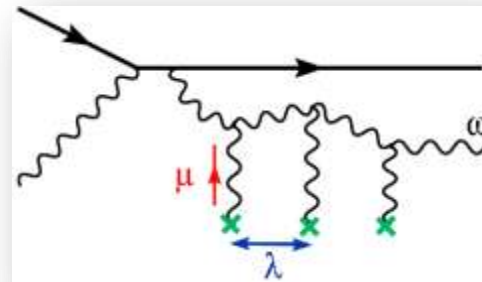
- **hard probes** ○ jet quenching ... **partonic energy loss**
 - ...

2 mechanisms ←

binary collisions



induced radiation



Energy Loss of Energetic Partons in Quark - Gluon Plasma: Possible Extinction of High $p(t)$ Jets in Hadron - Hadron Collisions.

J.D. Bjorken (Fermilab). Aug 1982. 20 pp.

FERMILAB-PUB-82-059-THY, FERMILAB-PUB-82-059-T

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[KEK scanned document](#); [Fermilab Library Server \(fulltext available\)](#)

[Detailed record](#) - [Cited by 39 records](#)

High energy quarks and gluons propagating through quark-gluon plasma suffer differential energy loss via elastic scattering from quanta in the plasma. This mechanism is very similar in structure to ionization loss of charged particles in ordinary matter. The dE/dx is roughly proportional to the square of the plasma temperature. For hadron-hadron collisions with high associated multiplicity and with transverse energy dE_T/dy in excess of 10 GeV per unit rapidity, it is possible that quark-gluon plasma is produced in the collision. If so, a produced secondary high- p_T quark or gluon might lose tens of GeV of its initial transverse momentum while plowing through quark-gluon plasma produced in its local environment. High energy hadron jet experiments should be analysed as function of associated multiplicity to search for this effect. An interesting signature may be events in which the hard collision occurs near the edge of the overlap region, with one jet escaping without absorption and the other fully absorbed.

similar to
ionization
loss

jet
quenching

correlations

$dE/dx \dots (t)$

Bjorken 1982

$$\frac{dE}{dx} \sim \alpha^2 T^2 \ln \frac{ET}{m_D^2}$$

Bethe-Bloch ... Bjorken

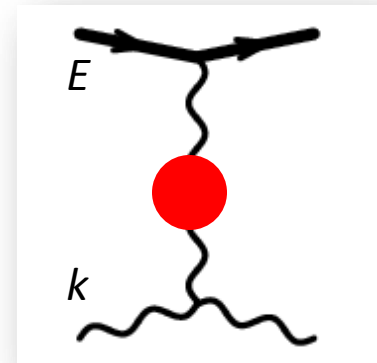
$$\frac{dE_i}{dx} \simeq \int_{k^3} n_i(k) \int dt \frac{d\sigma_i}{dt} \underbrace{\text{FluxFactor} \times \omega}_{\simeq -t/2k}$$

targets $i = \{\text{gluons, light flavors}\}$

dominating at large E : $\frac{d\sigma}{dt} \propto \frac{\alpha^2}{t^2}$

$$\frac{dE}{dx} \propto \int \frac{d^3k}{(2\pi)^3} \frac{n(k)}{2k} \int_{-ET}^{-\mu^2} dt \frac{\alpha^2}{t^2} (-t)$$

$$\propto T^2 \alpha^2 \ln \frac{ET}{\mu^2}$$



screening

uncertainties ...

dE/dx ... (t)

Braaten & Thoma 1991

$$\frac{dE}{dx} \sim \alpha^2 T^2 \left(\ln \frac{ET}{m_D^2} + \tilde{c} \right)$$

Bjorken 1982

$$\frac{dE}{dx} \sim \alpha^2 T^2 \ln \frac{ET}{m_D^2}$$

dE/dx ... (t)

AP 2006

$$\frac{dE}{dx} \sim \alpha(m_D^2)\alpha(ET) T^2 \ln \frac{ET}{m_D^2}$$

Braaten & Thoma 1991

$$\frac{dE}{dx} \sim \alpha^2 T^2 \left(\ln \frac{ET}{m_D^2} + \tilde{c} \right)$$

Bjorken 1982

$$\frac{dE}{dx} \sim \alpha^2 T^2 \ln \frac{ET}{m_D^2}$$

Whatever they do – ... RUN ...?



Whatever they do - RUN

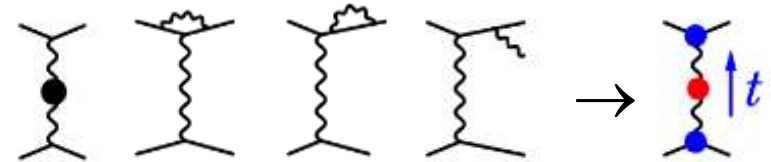
quantum corrections to Born X-section

thermal contributions

vacuum contributions

screening

running coupling



running coupling $\alpha(t) = \frac{4\pi/\beta_0}{\ln(-t/\Lambda^2)}$ **crucial** for QCD energy loss

$$\frac{dE}{dx} \propto \int \frac{d^3k}{(2\pi)^3} \frac{n(k)}{2k} \int_{-s}^{-m_D^2} dt \left(\frac{4\pi/\beta_0}{\ln(-t/\Lambda^2)} \right)^2 \frac{1}{t^2} (-t)$$

$$\propto T^2 \alpha(m_D^2) \alpha(ET) \ln \frac{ET}{m_D^2} \quad \text{AP 2006}$$

- ❖ **predictive** leading log result
- ❖ NB: $dE/dx \rightarrow$ **constant** for large E

dE/dx ... (t)

Peigné & AP 2008

$$\frac{dE}{dx} \sim \alpha(m_D^2)\alpha(ET)T^2 \ln \frac{ET}{m_D^2}$$
$$+ \alpha(M^2)\alpha(ET)T^2 \frac{2}{9} \ln \frac{ET}{M^2}$$
$$+ \alpha(ET)\alpha(m_D M) T^2 \tilde{c}$$

AP 2006

$$\frac{dE}{dx} \sim \alpha(m_D^2)\alpha(ET) T^2 \ln \frac{ET}{m_D^2}$$

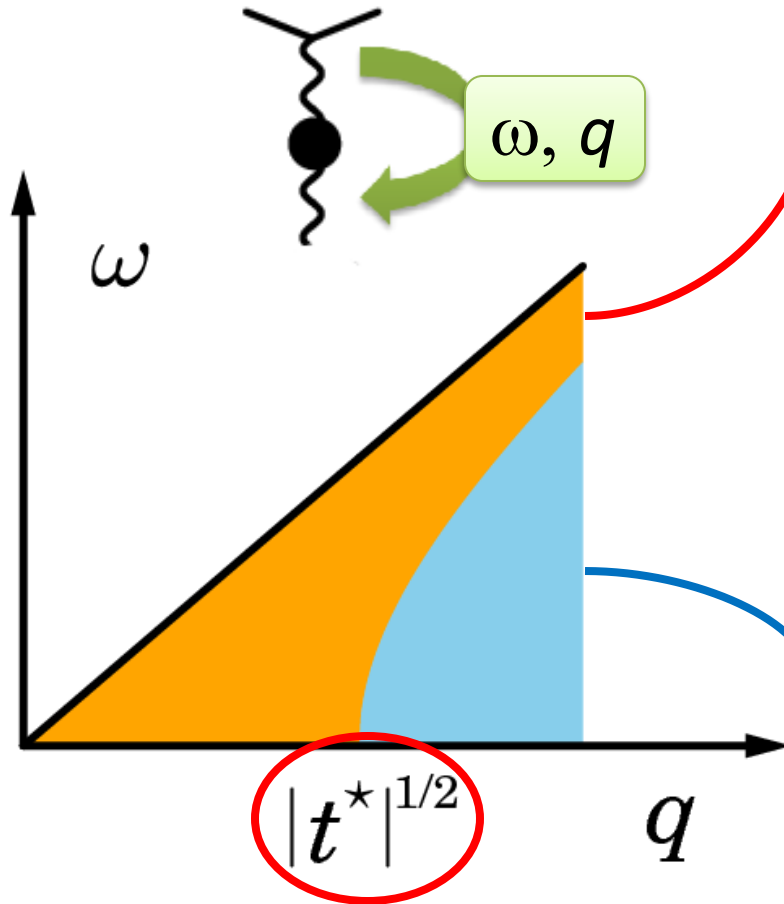
Braaten & Thoma 1991

$$\frac{dE}{dx} \sim \alpha^2 T^2 \left(\ln \frac{ET}{m_D^2} + \tilde{c} \right)$$

Bjorken 1982

$$\frac{dE}{dx} \sim \alpha^2 T^2 \ln \frac{ET}{m_D^2}$$

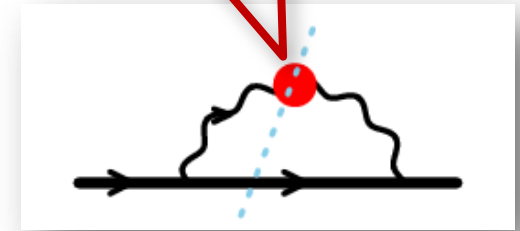
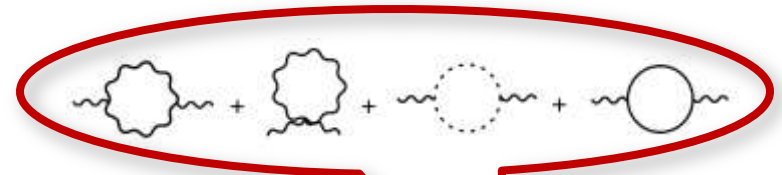
PP approach



separation scale,
drops out in final result

soft contribution

hard thermal loop
perturbation theory



hard contribution

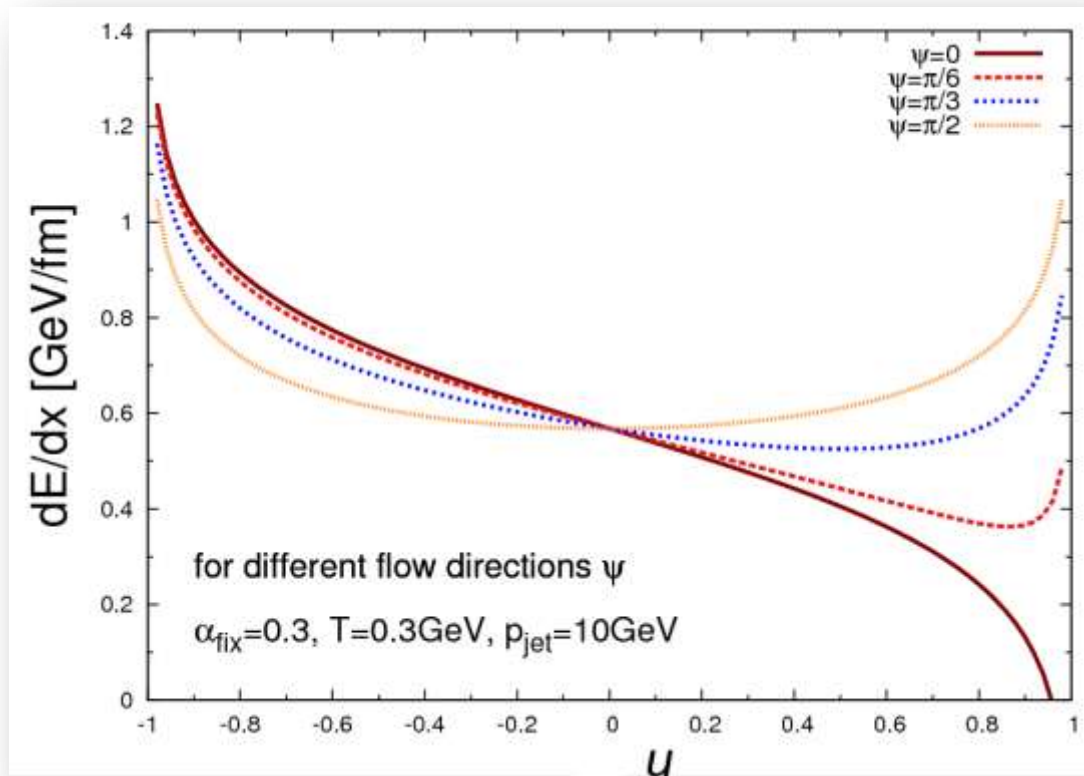
with Born X-sections

dE/dx with flow

$$\frac{dE}{dx} \sim \int \frac{d^3k}{(2\pi)^3} \frac{n(K_\mu u^\mu)}{2k} \int dt \frac{d\sigma}{dt}(-t)$$

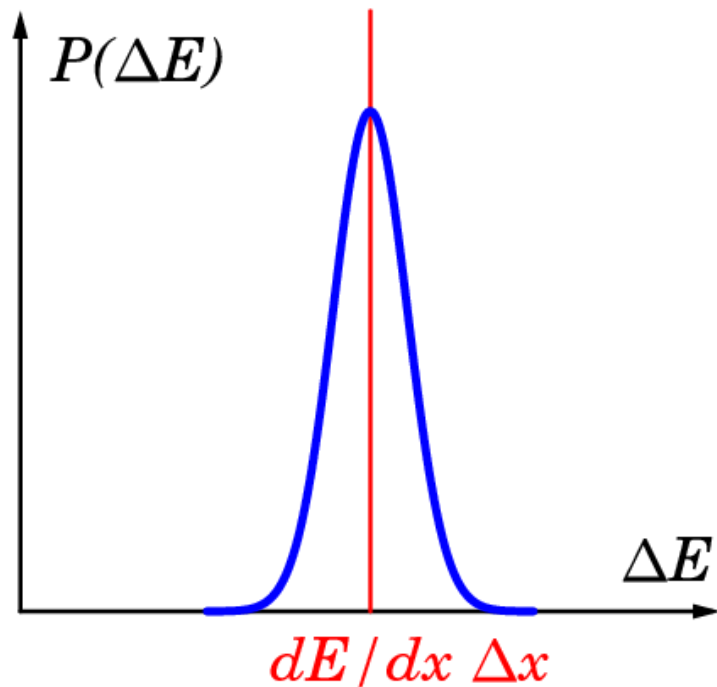
$$\sim T^2 \alpha(P_\mu u^\mu T) \left[\alpha(m_D^2) \ln \frac{P_\mu u^\mu T}{m_D^2} + \frac{2}{9} \alpha(M^2) \ln \frac{P_\mu u^\mu T}{M^2} \right]$$

Meistrenko,
Uphoff, Greiner
& AP 2012



Probabilistic description

dE/dx is 'stopping power' ... there is also 'range straggling'
→ need **probabilistic approach**



➤ how to NOT get the PDF

- deterministic loss
- Gauss smearing
- via Fokker-Planck eqn.

Collision rate

✓ quantum statistics

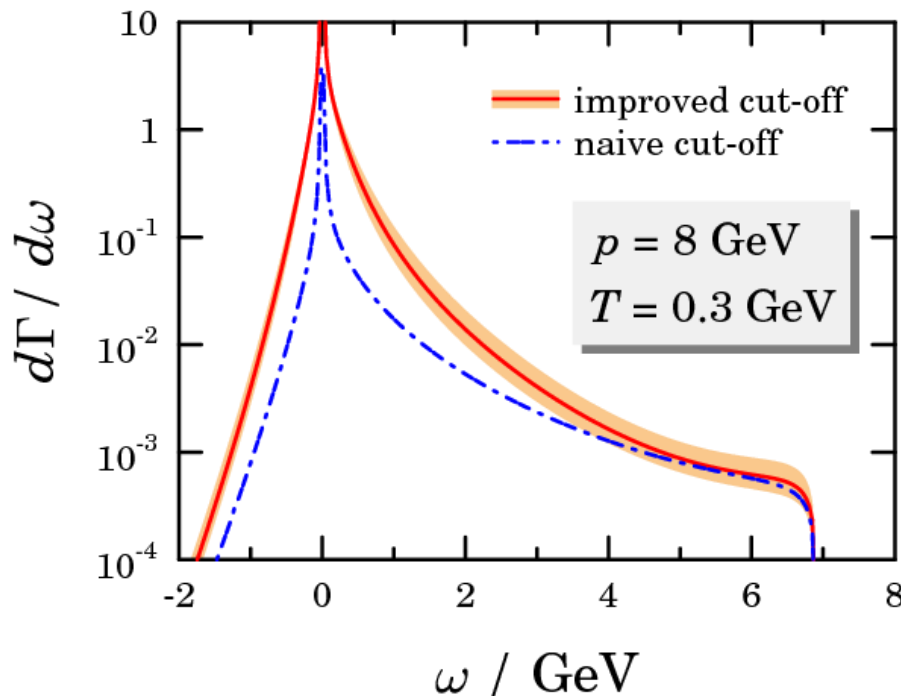
✓ kinematics

$$\frac{dE}{dx} = \frac{v^{-1}}{2E} \int_k \frac{n_k}{2k} \int_{k'} \frac{1 \pm n_{k'}}{2k'} \int_{p'} \frac{(2\pi)^4}{2E'} \delta^{(4)}(P+K-P'-K') \langle |\mathcal{M}|^2 \rangle_\omega$$

$$= \int d\omega \frac{d\Gamma}{d\omega} \omega$$

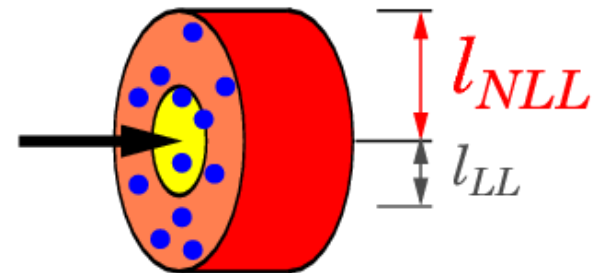
✓ running coupling
✓ effective mass cutoff

$$\mu^2(t) = \kappa 4\pi \left(1 + \frac{n_f}{6}\right) \alpha(t) T^2$$



$$\kappa \approx 0.2$$

→ enhanced interaction rate

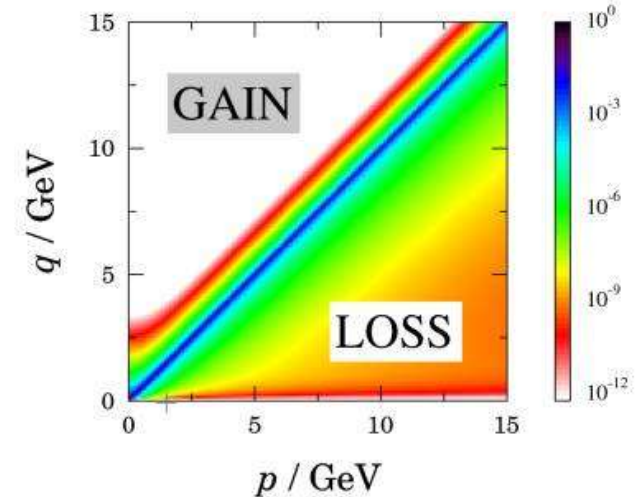


AP 2008

Markov (= Boltzmann) evolution

Consider **ensemble** of test particles.

Evolution of spectrum is **1st order Markov process**.

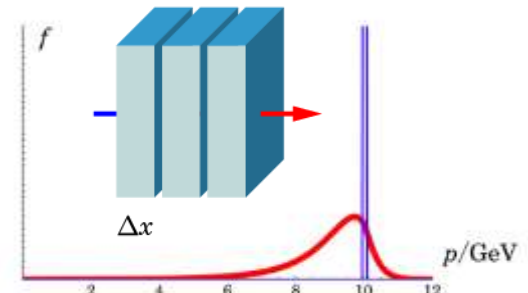
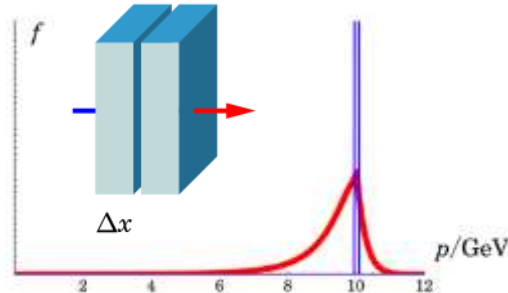
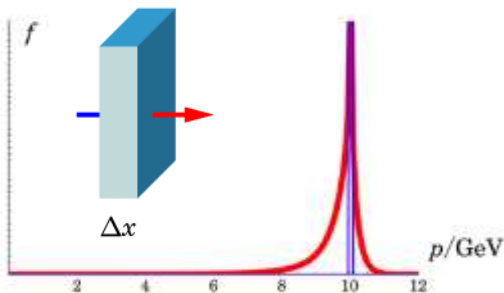


- in small interval δt :

$$f(q) \rightarrow f(q, \delta t) = f(q) - \delta t \Gamma(q) f(q) + \delta t \int dp \mathcal{P}(q, p) f(p)$$

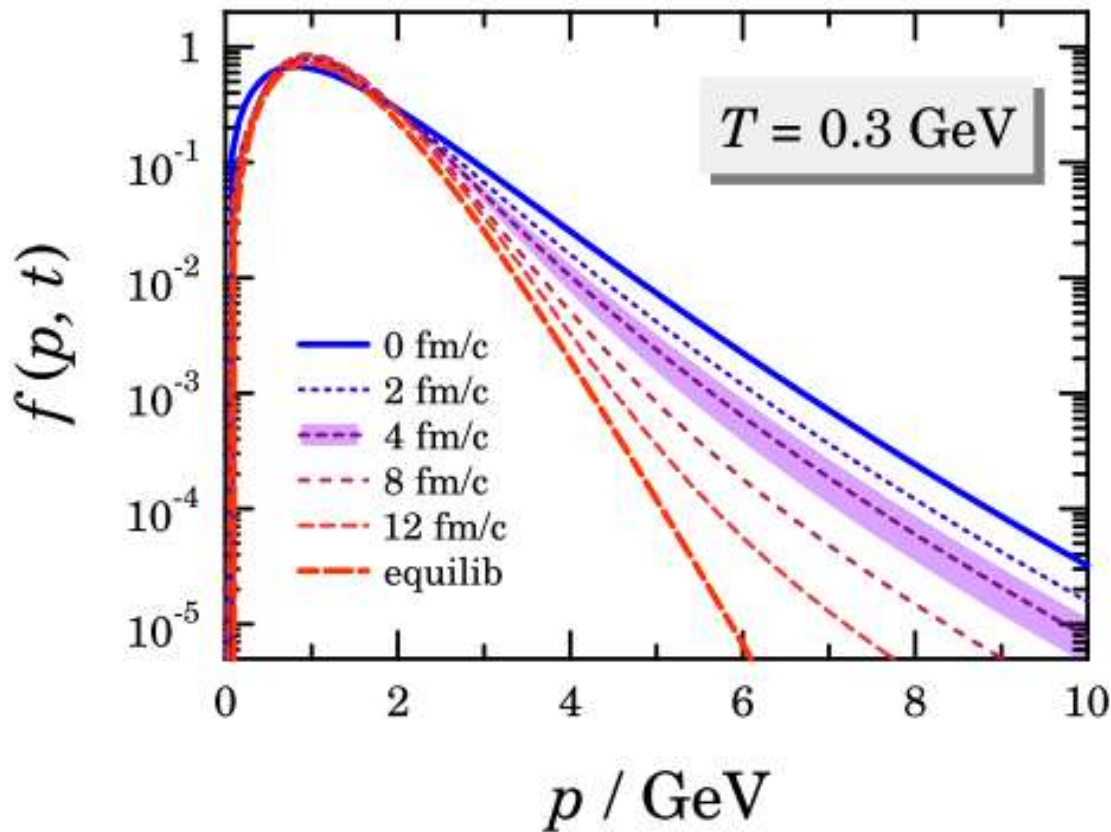
$$\equiv \int dp \mathcal{T}(q, p; \delta t) f(p) \quad \text{transition matrix}$$

- discretize (bin) momenta: $f_q(t) = (\mathcal{T}_{qp})^n f_p$ evolution = matrix power



Charm (partonic) equilibration

Initial charm spectrum in heavy ion collisions: $f(p, 0) = \left. \frac{dN}{dp} \right|_{ini} \sim p^{-\nu}$

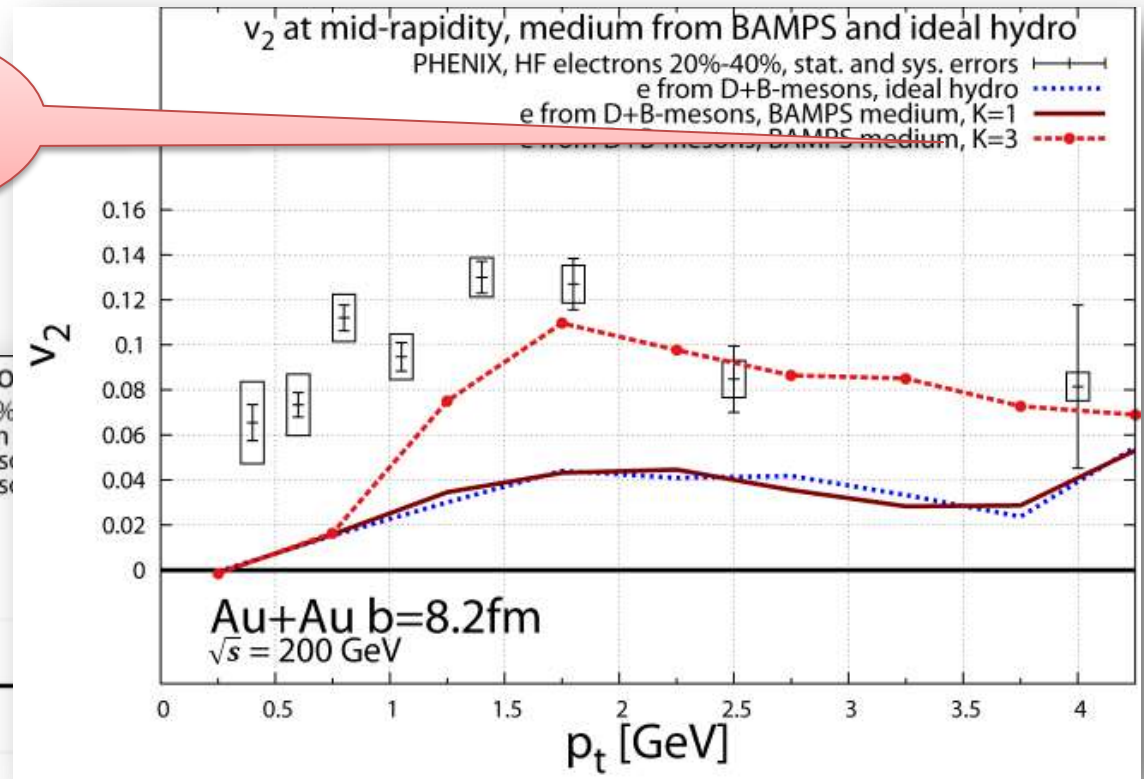
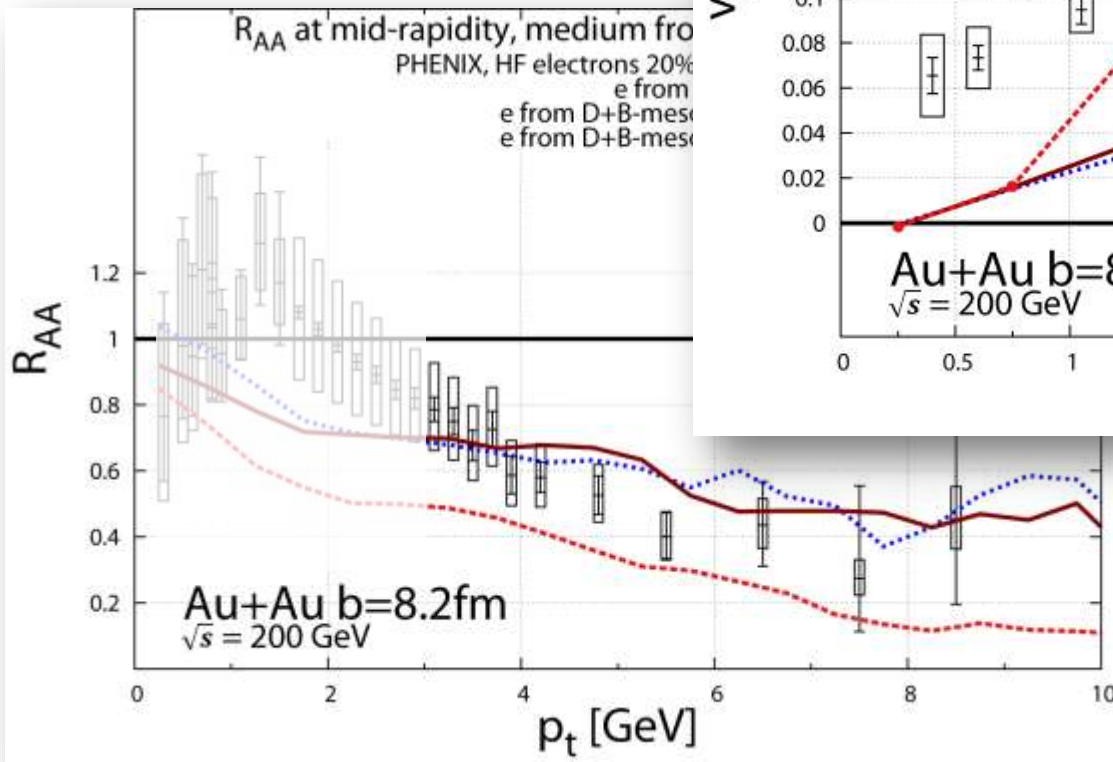


κ^{-1} -faster approach
to therm. equilibrium

AP 2008

Heavy flavor R_{AA} & v_2

Xsection
scaled by
 $K = 3$



Meistrenko,
Uphoff, Greiner
& AP 2012

Resumé

- ❖ need to be more precise than `prescriptions`/questionable approximations
- ❖ collisions need to be considered, besides gluon radiation, to understand heavy charm quenching and anisotropy
- ❖ stay tuned 😊

