

Applications of JIMWLK Evolution to Exclusive J/Ψ Production



Andrecia Ramnath

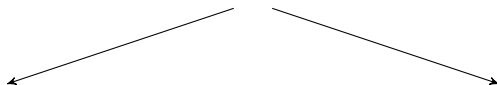
Supervisors: Prof. H. Weigert & Dr. A. Hamilton

Hard Probes Conference, STIAS (South Africa), 5 November 2013

Outline

Goal:

Determine **cross-section** for exclusive $J/\Psi \rightarrow \mu^+ \mu^-$ in nuclear collisions



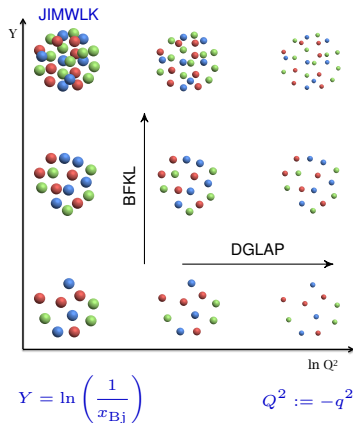
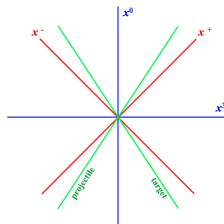
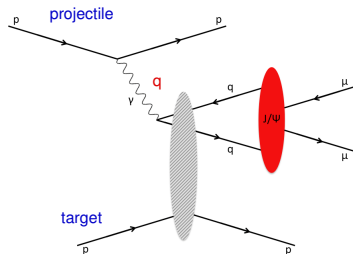
Theory:

- **Colour Glass Condensate** (CGC)
- Appropriate effective theory - **JIMWLK** evolution
- Truncate infinite hierarchy - **Gaussian Truncation** (GT)

Experiment:

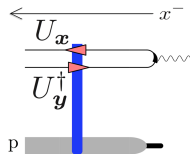
- **Vertexing method** to identify exclusive events in collision data
- Test method on **Starlight Monte Carlo**

Saturation Effects



Colour Interaction

Colour diagrams



Quark lines $\rightarrow$ Wilson lines

$$U_{\mathbf{x}} := \text{P exp} \left\{ -ig \int dx^- \mathbf{A}^+(0, x^-, \mathbf{x}) \right\}.$$

Generalise to generic U -correlator $\langle \dots \rangle_Y$

$\sim$ cross-section

JIMWLK Equation

(Jalilian-Marian - Iancu - McLerran - Weigert - Leonidov - Kovner)

Nonlinear, functional renormalization group equation

$$\frac{d}{dY} \langle \dots \rangle_Y = -H_{\text{JIMWLK}} \langle \dots \rangle_Y$$

$$H_{\text{JIMWLK}} := -\frac{\alpha_s}{2\pi^2} \int d^2z \mathcal{K}_{\mathbf{x}\mathbf{z}\mathbf{y}} i\nabla_{\mathbf{x}}^a \left[(1 - U_{\mathbf{x}}^\dagger U_{\mathbf{z}})(1 - U_{\mathbf{z}}^\dagger U_{\mathbf{y}}) \right]^{ab} i\nabla_{\mathbf{y}}^b$$

$$\mathcal{K}_{\mathbf{x}\mathbf{z}\mathbf{y}} = \frac{(\mathbf{x} - \mathbf{z}) \cdot (\mathbf{z} - \mathbf{y})}{(\mathbf{x} - \mathbf{z})^2 (\mathbf{z} - \mathbf{y})^2}$$

$$i\nabla_{\mathbf{x}}^a := -[U_{\mathbf{x}} t^a]^{ij} \frac{\delta}{\delta U_{\mathbf{x}}^{ij}} \longrightarrow$$



Evolution

E.g. Typical correlator

$$\langle \dots \rangle_Y = \left\langle \frac{\text{tr}(U_{\mathbf{x}} U_{\mathbf{y}}^\dagger)}{N_c} \right\rangle_Y = \frac{1}{o} \quad \text{Diagram: A vertical blue line with a single horizontal white loop at the top, representing a correlator with one insertion.$$

For J/Ψ we need correlators like

$$\langle \dots \rangle_Y = \left\langle \frac{\text{tr}(U_{\mathbf{x}} U_{\mathbf{y}}^\dagger) \text{tr}(U_{\mathbf{x}'} U_{\mathbf{y}'}^\dagger)}{N_c^2} \right\rangle_Y = \frac{1}{o^2} \quad \text{Diagram: A vertical blue line with two horizontal white loops at the top, representing a correlator with two insertions.$$

Evolution

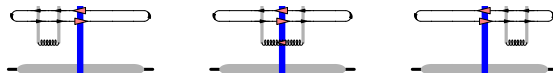
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LO



NLO



etc.

JIMWLK equation



Balitsky Hierarchy

...truncate!

Gaussian Truncation

Approximation:

$$\langle \dots \rangle_Y = \exp \left\{ -\frac{1}{2} \int^Y dY' \int d^2u d^2v G_{Y',uv} i\nabla_u^a i\nabla_v^a \right\} \dots$$

or conveniently,

$$\frac{d}{dY} \langle \dots \rangle_Y = -\frac{1}{2} \left\langle \int d^2u d^2v G_{Y,uv} i\nabla_u^a i\nabla_v^a \dots \right\rangle_Y$$

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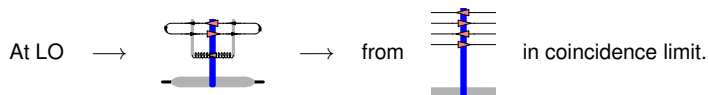
E.g. $\frac{d}{dY} \left(\frac{1}{\circ} \text{diagram} \right) = \frac{1}{\circ} \text{diagram}$

Gaussian Truncation

Use GT to find expressions for $n + 1$ -order terms in evolution equation at n -order.

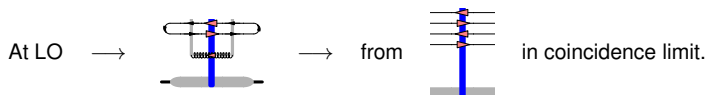
- **Group theory** constraints automatically satisfied exactly (check coincidence limits).
- Unlike BK equation, **no large N_c** limit necessary.
- **Black disc** limit vanishes, i.e. $\text{tr}(U_{\mathbf{x}} U_{\mathbf{y}}^\dagger) \rightarrow 0$ when $|\mathbf{x} - \mathbf{y}| \rightarrow \infty$

Parametrisation Equations



Two dipoles $\implies$ multiplet basis $\frac{1}{\omega} \begin{pmatrix} \circ \\ \circ \end{pmatrix}, \frac{1}{\omega^{\frac{1}{2}}} \begin{pmatrix} \circ \\ \circ \end{pmatrix}$

Parametrisation Equations



Two dipoles $\Rightarrow$ multiplet basis $\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix}$

$$\frac{d}{dY} \begin{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix} \\ \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix} \end{pmatrix} (Y) = \begin{pmatrix} a_Y & c_Y \\ c_Y & b_Y \end{pmatrix} \begin{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix} \\ \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix} \end{pmatrix}$$

Marquet & Weigert, arXiv:1003.0813 (2010)

Beyond GT

Previously,

$$\frac{d}{dY} \langle \dots \rangle_Y = -\frac{1}{2} \left\langle \int d^2u d^2v G_{Y,u\bar{v}} i\nabla_u^a i\nabla_{\bar{v}}^a \dots \right\rangle_Y$$

Now include on the RHS:

$$\int d^2u d^2v d^2w \left(f^{abc} G_{uvw}^f + d^{abc} G_{uvw}^d \right) i\nabla_u^a i\nabla_v^b i\nabla_w^c$$

where G_{uvw}^f antisymmetric and G_{uvw}^d symmetric.

(There is a piece from the antisymmetric term that has a double derivative:

$$\frac{iN_c}{4} \int d^2u d^2v \left(-G_{uvu}^f - G_{vuv}^f + G_{uvv}^f + G_{vuu}^f + G_{uuv}^f + G_{vvu}^f \right) \{i\nabla_u^a, i\nabla_v^a\}$$

Beyond the GT

Final result for the evolution equation for a quark dipole

$$\frac{d}{dY} \left\langle \left(G'_{\mathbf{x}\mathbf{y}} + 3i G_{\mathbf{x}\mathbf{y}}^{\mathcal{O}} \right) (Y) \right\rangle = \frac{\alpha_s}{\pi^2} \tilde{\mathcal{K}}_{\mathbf{x}\mathbf{z}\mathbf{y}} \left\langle 1 - e^{\frac{N_c}{2}} \left\{ \left(G'_{\mathbf{x}\mathbf{y}} - G'_{\mathbf{x}\mathbf{z}} - G'_{\mathbf{z}\mathbf{y}} \right) + 3i \left(G_{\mathbf{x}\mathbf{y}}^{\mathcal{O}} - G_{\mathbf{x}\mathbf{z}}^{\mathcal{O}} - G_{\mathbf{z}\mathbf{y}}^{\mathcal{O}} \right) \right\} (Y) \right\rangle_Y$$

$$\text{where } G_{\mathbf{x}\mathbf{y}}^{\mathcal{O}} := \frac{C_d}{2} \left(G_{\mathbf{y}\mathbf{x}\mathbf{x}}^d - G_{\mathbf{y}\mathbf{y}\mathbf{x}}^d \right)$$

Beyond the GT

Final result for the evolution equation for a quark dipole

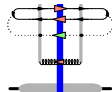
$$\frac{d}{dY} \left\langle \left(G'_{xy} + 3i G_{xy}^O \right) (Y) \right\rangle = \frac{\alpha_s}{\pi^2} \tilde{\mathcal{K}}_{xyz} \left\langle 1 - e^{\frac{N_c}{2} \left\{ \left(G'_{xy} - G'_{xz} - G'_{zy} \right) + 3i \left(G_{xy}^O - G_{xz}^O - G_{zy}^O \right) \right\} (Y)} \right\rangle_Y$$

$$\text{where } G_{xy}^O := \frac{C_d}{2} \left(G_{yx}^d - G_{yyx}^d \right)$$

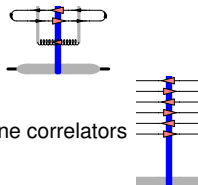
Further investigation

To understand G_{xy}^O , need to find evolution equation for

⇒ parametrise



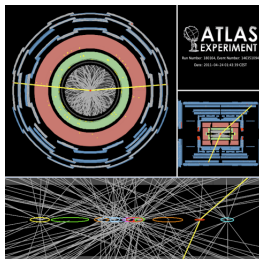
⇒ go to 6 Wilson line correlators



Finally, result for $\left\langle \dots \right\rangle_Y$ goes into cross-section.

Experiment

Determine cross-section from collision data.

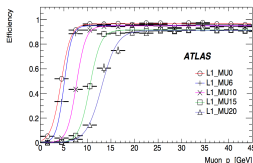
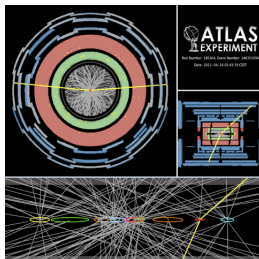


Selection Criteria:

- Each event has 2 muons within detector acceptance
- The muons come from a single vertex
- No other tracks come from the vertex
- The muon pair has a low p_T

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Challenges:

- Low p_T muon efficiency is poor
- No guarantee the protons did not dissociate
- Pile-up when luminosity is high $\Rightarrow$ go to p-Pb

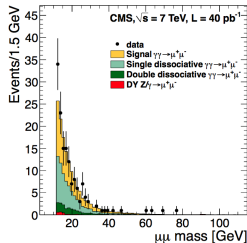
CMS and LHCb Exclusive Studies

CMS [arXiv:1111.5536 (2012)]

Exclusive $\gamma\gamma \rightarrow \mu^+ \mu^-$

148 events; approx. half are exclusive

$$\sigma(pp \rightarrow p\mu^+\mu^-p) = 3.38 + 0.58(\text{stat.}) \pm 0.16(\text{syst.}) \pm 0.14(\text{lumi.}) \text{ pb}$$



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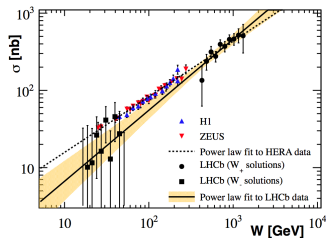
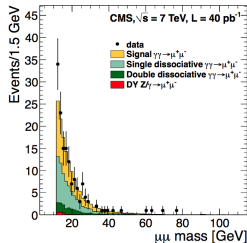
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LHCb [arXiv:1301.7084 (2013)]

$$\mathcal{L} = 36 \text{ pb}^{-1}$$

$$\sigma_{pp \rightarrow J/\Psi(\rightarrow \mu^+ \mu^-)}(2.0 < \eta_\mu < 4.5) = 307 \pm 21(\text{stat.}) \pm 36(\text{syst.}) \text{ pb}$$



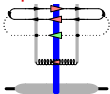
Conclusions

- The **Gaussian truncation** is a useful tool to truncate the infinite hierarchy of equations, the Balitsky hierarchy (equivalent to JIMWLK equation).
- **Parametrisation equations** help to find simple expressions for correlators.
- Several **experimental challenges** in detecting exclusive events → ongoing work with ATLAS.

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Outlook

- Determine **6-dipole correlator** parametrisation equations in order to find evolution equation for .
- Go beyond three-derivative operators.

End