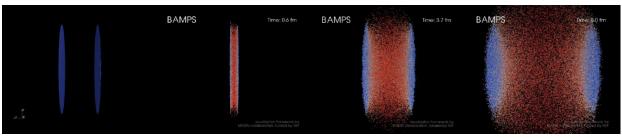


# Heavy vs. light quark energy loss in ultra-relativistic heavy-ion collisions

## Jan Uphoff

with O. Fochler, Z. Xu and C. Greiner

Based on Phys. Lett. B 717, 430 (2012) and Phys. Rev. D88 (2013)



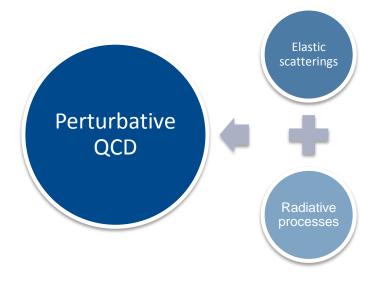




Hard Probes, Stellenbosch 5 November 2013

# Heavy quark energy loss mechanism





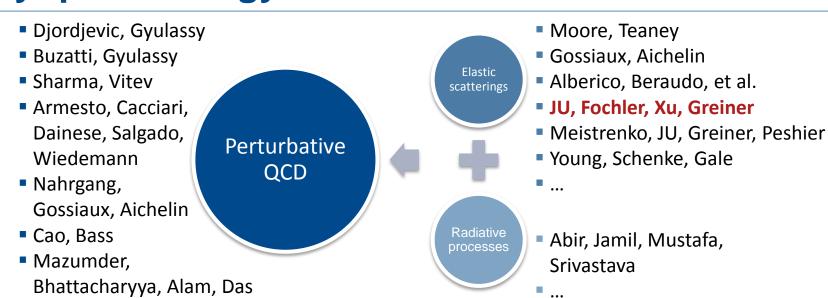




## Heavy quark energy loss mechanism

JU, Fochler, Xu, Greiner







#### **BAMPS**



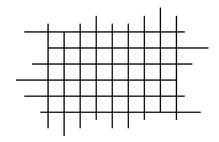
#### BAMPS: Boltzmann Approach to MultiParton Scatterings

- 3+1 dimensional, fully dynamic parton transport model
- solves the Boltzmann equations for on-shell partons with pQCD interactions

$$\left(\frac{\partial}{\partial t} + \frac{\mathbf{p}_i}{E_i} \frac{\partial}{\partial \mathbf{r}}\right) f_i(\mathbf{r}, \mathbf{p}_i, t) = \mathcal{C}_i^{2 \to 2} + \mathcal{C}_i^{2 \leftrightarrow 3} + \dots$$
Z. Xu & C. Greiner, Phys. Rev. C71 (2005) Phys. Rev. C76 (2007)

Phys. Rev. C76 (2007)

Divide collision zone into cells



**Using stochastic method** 

$$P_{2\to 2} = v_{\rm rel} \frac{\sigma_{2\to 2}}{N_{\rm test}} \frac{\Delta t}{\Delta^3 x}$$

# Interactions in BAMPS with $N_{flavor} = 3+2$



## **Light flavors**

# g g o g g $g g o q ar{q}$ q g o g g and $q ar{q} o q' ar{q}'$ q g o q g and $q g o q' ar{q}'$ q g o q g and $\bar{q} g o \bar{q} g$ $q \bar{q} o q \bar{q}$ q q o q q and $\bar{q} q o q \bar{q}$ q q o q q and q q o q q'

$$g g \leftrightarrow g g g$$
 $q g \leftrightarrow q g g$  and  $\bar{q} g \leftrightarrow \bar{q} g g$ 
 $q \bar{q} \leftrightarrow q \bar{q} g$ 
 $q q \leftrightarrow q q g$  and  $\bar{q} \bar{q} \leftrightarrow \bar{q} \bar{q} g$ 
 $q q' \leftrightarrow q q' g$  and  $q \bar{q}' \leftrightarrow q \bar{q}' g$ 

# Heavy flavors

$$g + g \rightarrow Q + Q$$

$$Q + \bar{Q} \rightarrow g + g$$

$$q + \bar{q} \rightarrow Q + \bar{Q}$$

$$Q + \bar{Q} \rightarrow q + \bar{q}$$

$$Q + Q \rightarrow q + Q$$

$$q + Q \rightarrow q + Q$$

$$q + Q \rightarrow q + Q$$

$$g + J/\psi \rightarrow c + \bar{c}$$

$$c + \bar{c} \rightarrow g + J/\psi$$

## inelastic

binary

$$g + Q \rightarrow g + Q + g$$
  
 $q + Q \rightarrow q + Q + g$ 

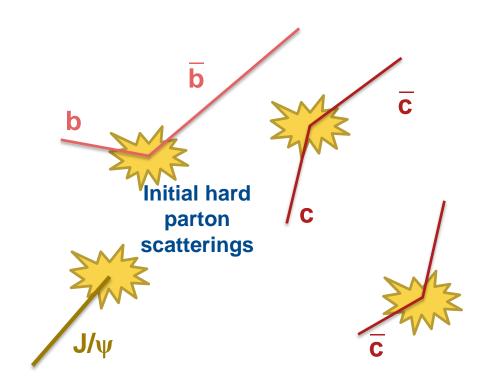




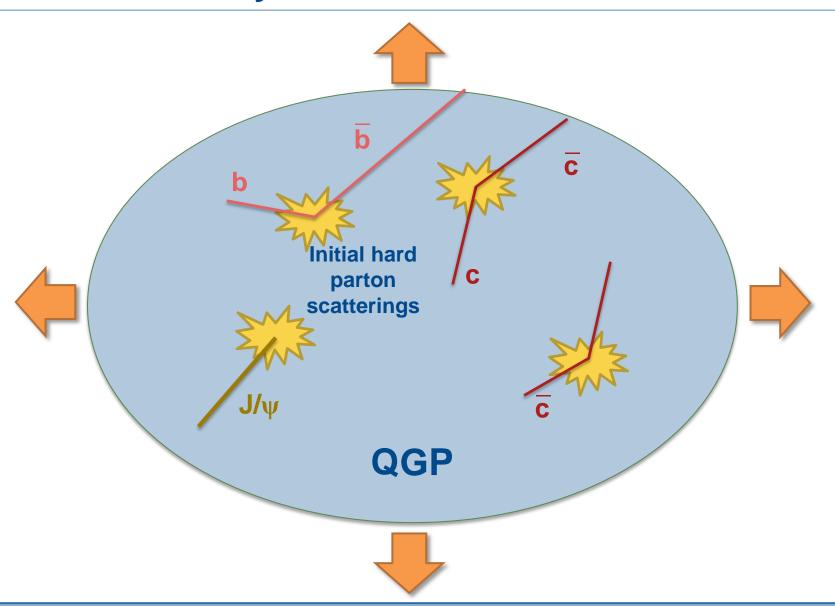




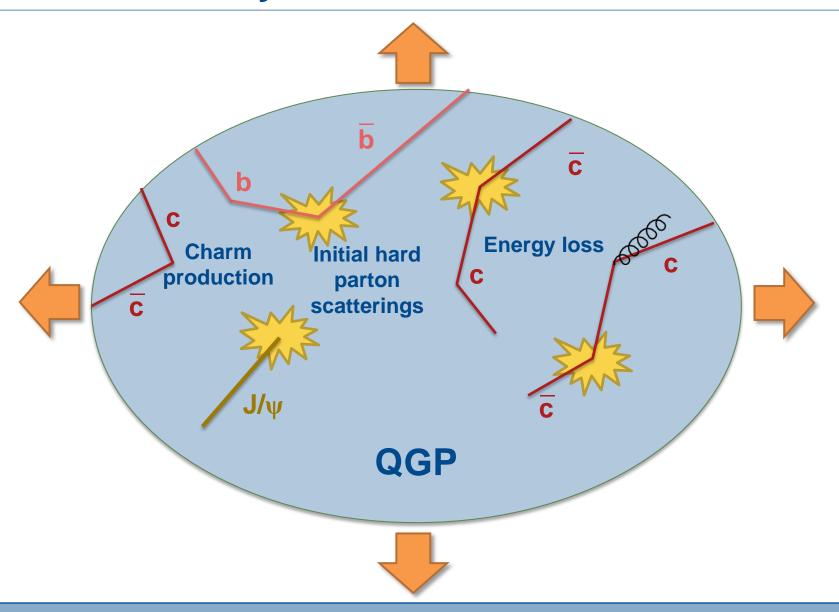




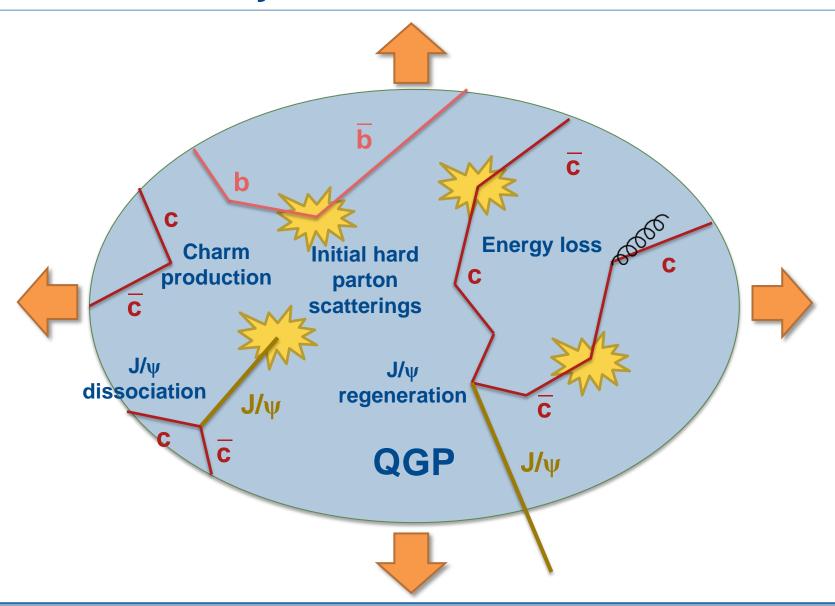




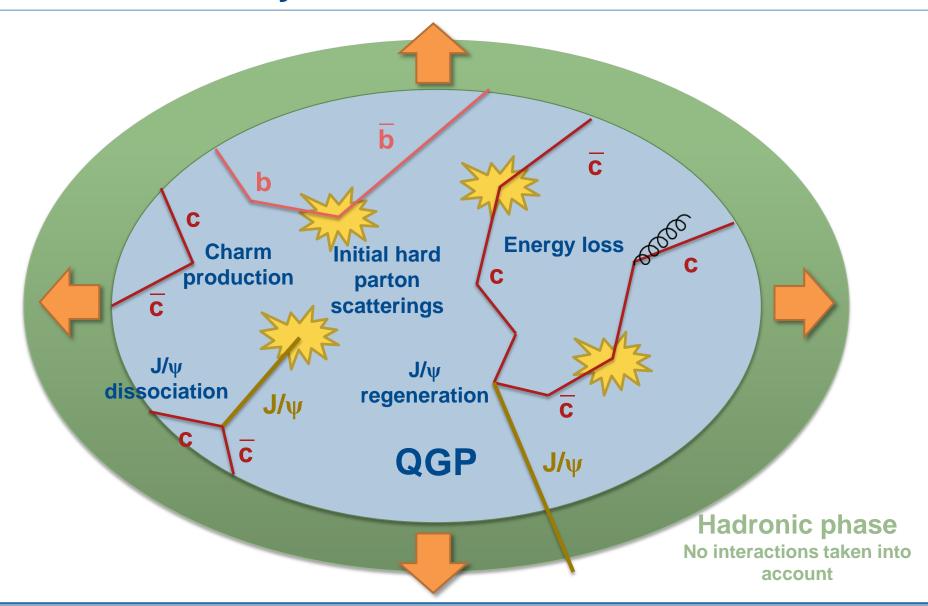




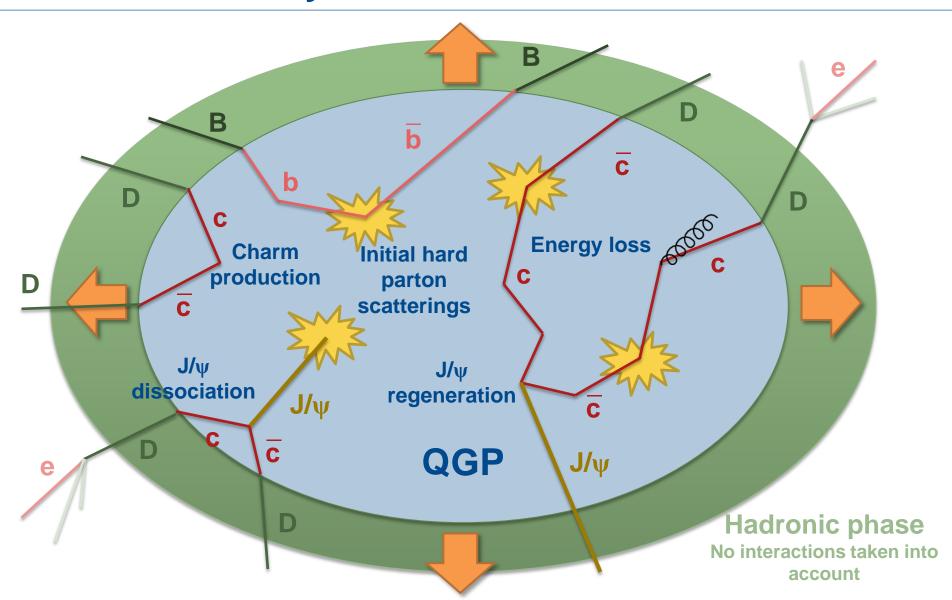












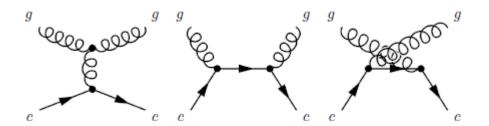
# **Heavy quark scattering**



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#### **Leading order perturbative QCD:**

$$g + Q \rightarrow g + Q$$
  
 $q + Q \rightarrow q + Q$ 



# Improved Debye screening by comparing to HTL

Bethke et al. 2006

$$\alpha_s(Q)$$

0.4

Deep Inelastic Scattering

• ete Annihilation

Hadron Collisions

Heavy Quarkonia

0.3

 $\alpha_s(Q)$ 
 $\alpha_s(M_z) = 0.1189 \pm 0.0010$ 

10

 $\alpha_s(M_z) = 0.1189 \pm 0.0010$ 

$$\frac{1}{t} \to \frac{1}{t - \kappa \, m_D^2}$$

$$\kappa = \frac{1}{2e} \approx 0.2$$

A. Peshier, Nucl.Phys. A888 (2012)

P.B. Gossiaux, J. Aichelin, Phys.Rev.C78 (2008)

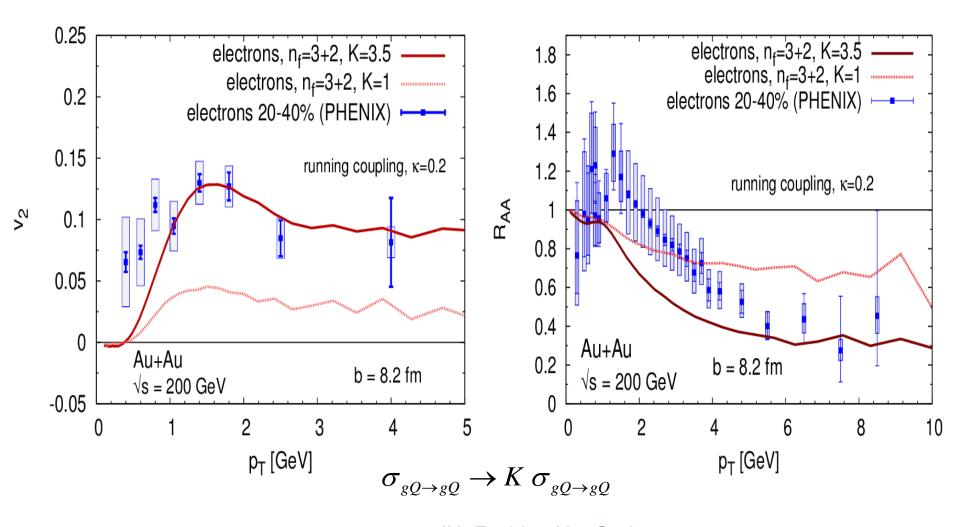
#### **Running coupling**

Details: JU, Fochler, Xu, Greiner Phys. Rev. C 84 (2011)

# Heavy quark v<sub>2</sub> and R<sub>AA</sub> at RHIC





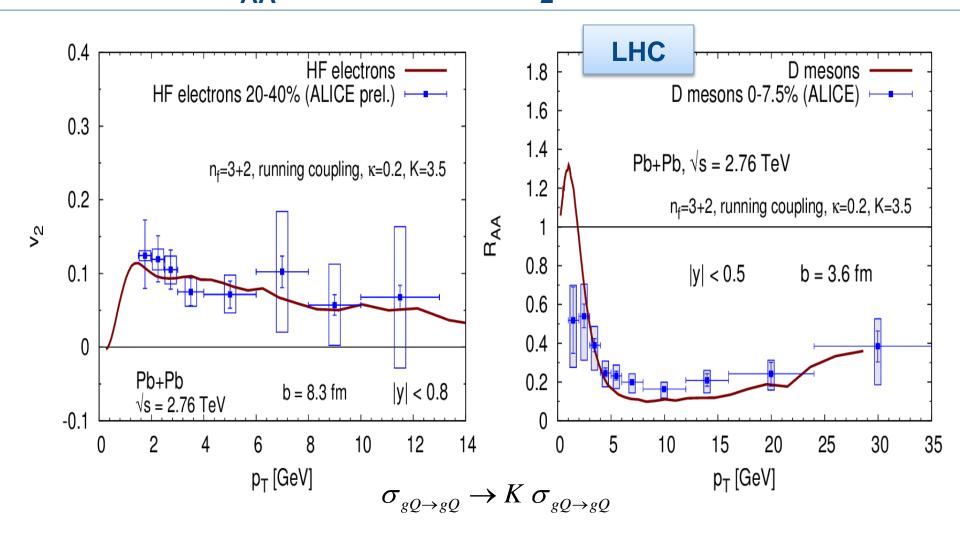


only elastic heavy quark processes

JU, Fochler, Xu, Greiner Phys. Lett. B 717 (2012) PHENIX data, Phys.Rev. C84 (2011)

# D meson $R_{AA}$ and electron $v_2$ at LHC





only elastic heavy quark processes

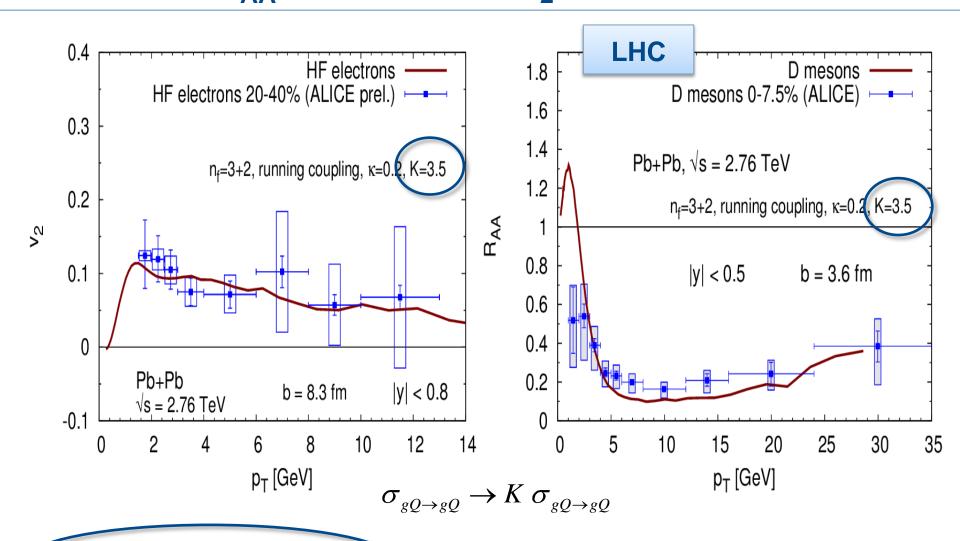
JU, Fochler, Xu, Greiner Phys. Lett. B 717 (2012)

ALICE data, QM12

14

# D meson $R_{AA}$ and electron $v_2$ at LHC





only elastic heavy quark processes

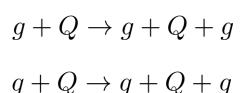
JU, Fochler, Xu, Greiner Phys. Lett. B 717 (2012)

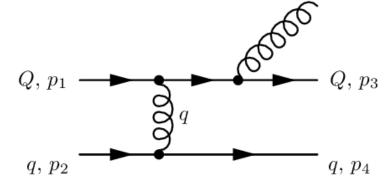
ALICE data, QM12

# Radiative processes: Improved Gunion-Bertsch matrix element



Can radiative processes account for K~3.5?





#### Improved Gunion-Bertsch matrix element generalized to heavy quarks:

$$\left| \overline{\mathcal{M}}_{qQ \to qQg} \right|^2 = 12g^2 (1 - \bar{x})^2 \left| \overline{\mathcal{M}}_0^{qQ} \right|^2 \left[ \frac{\mathbf{k}_{\perp}}{k_{\perp}^2 + x^2 M^2} + \frac{\mathbf{q}_{\perp} - \mathbf{k}_{\perp}}{(\mathbf{q}_{\perp} - \mathbf{k}_{\perp})^2 + x^2 M^2} \right]^2$$

Fochler, JU, Xu, Greiner, Phys. Rev. D88 (2013)

In accordance to scalar QCD result at mid- and forward rapidity from Gossiaux, Aichelin, Gousset, Guiho, J.Phys.G37 (2010)

# Radiative pQCD processes



#### **Exact matrix element**

Kunszt, Pietarinen, Reya, Phys.Rev. D21 (1980)

$$\left|\overline{\mathcal{M}}\right|^2 = -16\sum_{i,j=1}^5 C_{ij} \frac{N_{ij}}{D_{ij}}$$

$$C = \frac{1}{3} \begin{bmatrix} 8 & 1 & 9 & -2 & -7 \\ 8 & 9 & -7 & -2 \\ & 18 & -9 & -9 \\ & & 8 & 1 \\ & & & 8 \end{bmatrix}$$

 $N_{11} = x_{31} \left( -x_{42} x_{53} - x_{43} x_{52} + 2 m_Q^2 x_{54} \right) + 2 m_Q^2 \left( x_{41} x_{52} + x_{42} x_{51} + x_{42} x_{53} + x_{43} x_{52} + 2 m_Q^2 x_{54} \right),$  $N_{12} = x_{12} \left[ x_{41} (2x_{52} + x_{53}) + x_{42} (2x_{51} + x_{53}) + x_{43} (x_{51} + x_{52}) + 4m_Q^2 x_{54} \right]$ 

- $+ x_{23}[x_{41}(-2x_{51} + x_{52}) + x_{42}x_{51} + 2m_Q^2x_{54}]$
- $+ x_{31}[x_{41}x_{52} + x_{42}(x_{51} 2x_{52}) + 2m_Q^2x_{54}] 4m_Q^2x_{43}x_{53}$

 $N_{13} = x_{12} \left[ -2x_{21}x_{54} + x_{41}(4x_{52} + 3x_{53}) + x_{42}(4x_{51} + 3x_{53}) + x_{43}(3x_{51} + 3x_{52} + 2x_{53}) + 8m_0^2 x_{54} \right]$ 

- $+x_{23}[x_{41}(-6x_{51}+x_{52}-x_{53})+(x_{42}-x_{43})x_{51}+4m_0^2x_{54}]$
- $+x_{31}[x_{41}x_{52}+x_{42}(x_{51}-2x_{52}-3x_{53})-3x_{43}x_{52}]$
- $+2m_0^2[x_{41}(2x_{51}+4x_{52}+3x_{53})+x_{42}(4x_{51}-2x_{52}+5x_{53})+x_{43}(3x_{51}+5x_{52})+8m_0^2x_{54}],$

 $N_{14} = x_{51}(-x_{12}x_{43} + x_{23}x_{41} + x_{42}x_{43}) + x_{31}[2x_{41}x_{52} + x_{42}(x_{51} - x_{54}) + 2x_{43}x_{52} + 2m_0^2x_{54}]$ 

 $+x_{41}[2(x_{41}+x_{43})x_{52}+x_{42}(2x_{51}+x_{53})+2m_0^2(x_{53}+2x_{54})]+2m_0^2x_{43}(x_{53}+x_{54}-x_{51}-x_{52})$ 

 $N_{15} = N_{14}(4 - 5)$ ,  $N_{22} = N_{11}(1 - 2)$ ,  $N_{23} = N_{13}(1 - 2)$ ,  $N_{24} = N_{14}(1 - 2)$ ,  $N_{25} = N_{24}(4 - 5)$ 

 $N_{33} = x_{12} \left[ 2x_{54} (x_{12} - x_{23} - x_{31}) + x_{41} (-2x_{51} + 6x_{52} + 5x_{53}) + x_{42} (-2x_{52} + 6x_{51} + 5x_{53}) \right]$ 

- $+x_{43}(5x_{51}+5x_{52}+4x_{53})+28m_0^2x_{54}$
- $+x_{25}[-2x_{31}x_{54}+x_{41}(-6x_{51}+x_{52}-3x_{53})+(x_{42}-3x_{43})x_{51}]$
- $+ x_{31}[x_{41}x_{52} + x_{42}(x_{51} 6x_{52} 3x_{53}) 3x_{43}x_{52}]$

 $+2m_{Q}^{2}[x_{41}(-4x_{51}+4x_{52}+7x_{53})+x_{42}(4x_{51}-4x_{52}+7x_{53})+x_{43}(7x_{51}+7x_{52}+2x_{53})+24m_{Q}^{2}x_{54}],$ 

 $N_{34} = x_{12}[x_{54}(x_{23} + x_{31} + x_{42} + x_{41}) - 2x_{43}(x_{51} + x_{52})]$ 

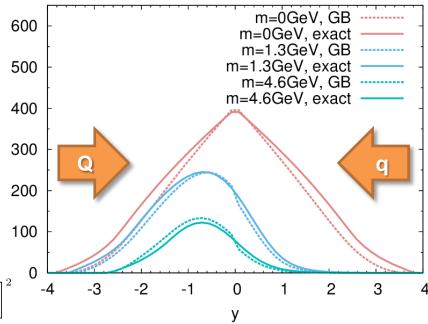
- $+x_{25}[2x_{31}x_{54}+x_{41}(x_{52}-x_{53}-x_{54})+3x_{51}(x_{42}+x_{43})+8m_Q^2x_{54}]$
- $+ x_{51}[x_{42}(x_{51} x_{53} x_{54}) + 3x_{52}(x_{41} + x_{43}) + 8m_Q^2x_{54}]$
- $+ x_{41}[3(x_{41} + x_{43})x_{52} + x_{42}(x_{51} + x_{52} + 2x_{53}) + 2m_Q^2(x_{53} + 5x_{54})]$
- $+x_{42}[3(x_{42}+x_{43})x_{51}+2m_Q^2(x_{53}+5x_{54})]+2m_Q^2x_{43}(2x_{53}+2x_{54}-5x_{51}-5x_{52})$

 $N_{44} = x_{43}(-x_{23}x_{51} - x_{31}x_{52} - 2m_Q^2x_{53}),$ 

 $N_{45} = x_{54} \big[ x_{23} (x_{41} + x_{51}) + x_{31} (x_{42} + x_{52}) + 4 m_{\rm Q}^{\ 2} (x_{53} + x_{54} + x_{43}) \big]$ 

 $+\,x_{41}[x_{52}(x_{43}+2x_{54})+x_{53}(x_{52}-2x_{42})\,]+x_{51}[-\,2x_{43}x_{52}+x_{42}(x_{53}+2x_{54}+x_{43})\,]\,,$ 

 $N_{35} = N_{34}(4 \rightarrow 5)$ ,  $N_{55} = N_{44}(4 \rightarrow 5)$ .



#### **Gunion Bertsch (GB) approximation**

$$\left| \overline{\mathcal{M}}_{qQ \to qQg} \right|^2 = 12g^2 (1 - \bar{x})^2 \left| \overline{\mathcal{M}}_0^{qQ} \right|^2 \left[ \frac{\mathbf{k}_{\perp}}{k_{\perp}^2 + x^2 M^2} + \frac{\mathbf{q}_{\perp} - \mathbf{k}_{\perp}}{(\mathbf{q}_{\perp} - \mathbf{k}_{\perp})^2 + x^2 M^2} \right]^2$$

Fochler, JU, Xu, Greiner, Phys. Rev. D88 (2013)

dσ/dy (a.u.)

#### **Dead cone effect**



$$\left| \overline{\mathcal{M}}_{qQ \to qQg} \right|^2 = \left| \overline{\mathcal{M}}_{qq \to qqg} \right|^2 \mathcal{D}$$

#### **Heavy quark suppression factor**

$$\mathcal{D} = \frac{1}{\left(1 + \frac{M^2}{\theta^2 E^2}\right)^2} = \frac{1}{\left(1 + \frac{\theta_D^2}{\theta^2}\right)^2}$$

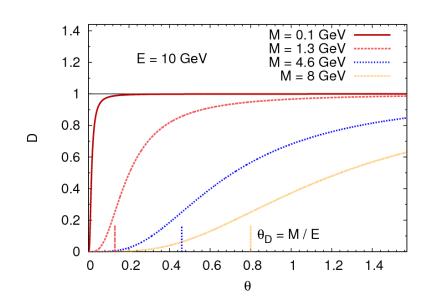
Dokshitzer, Kharzeev, Phys.Lett. B519 (2001)

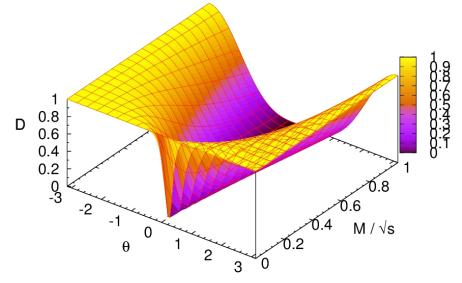
$$\theta_D = \frac{M}{E}$$

# More accurate: valid for all order of mass M and also for large angles

$$\mathcal{D} = \frac{1}{1 + \frac{M^2}{s \tan^2(\frac{\theta}{2})}}$$

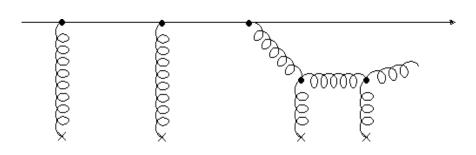
Abir, Greiner, Martinez, Mustafa, JU, Phys.Rev. D85 (2012)





### LPM cut-off





#### Mean free path

$$\lambda > X \tau$$

#### **Formation time**

 $2 \to 3$  process only allowed if mean free path of jet larger than formation time of radiated gluon

$$X = 0$$

No LPM

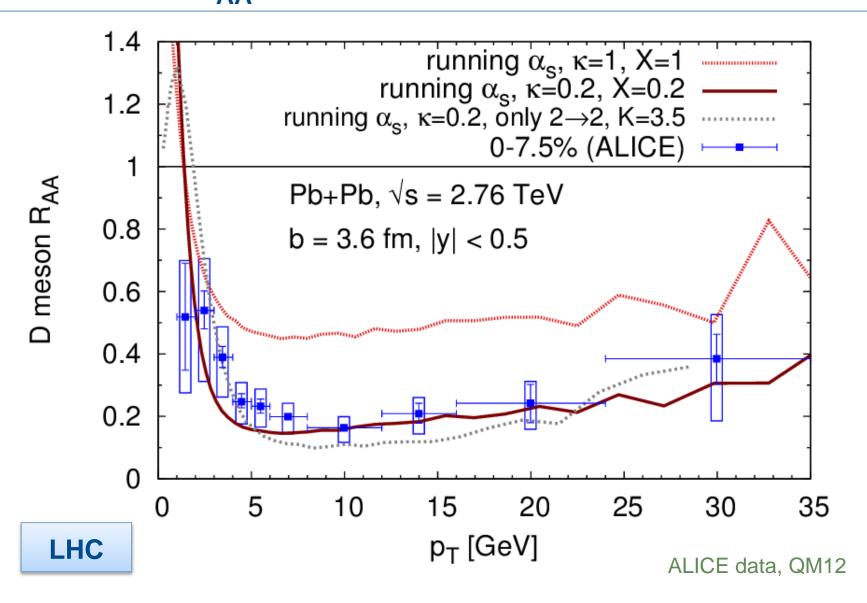
$$X = 1$$

Only completely independent scatterings (forbids too many interactions)

Allows effectively some interference effects

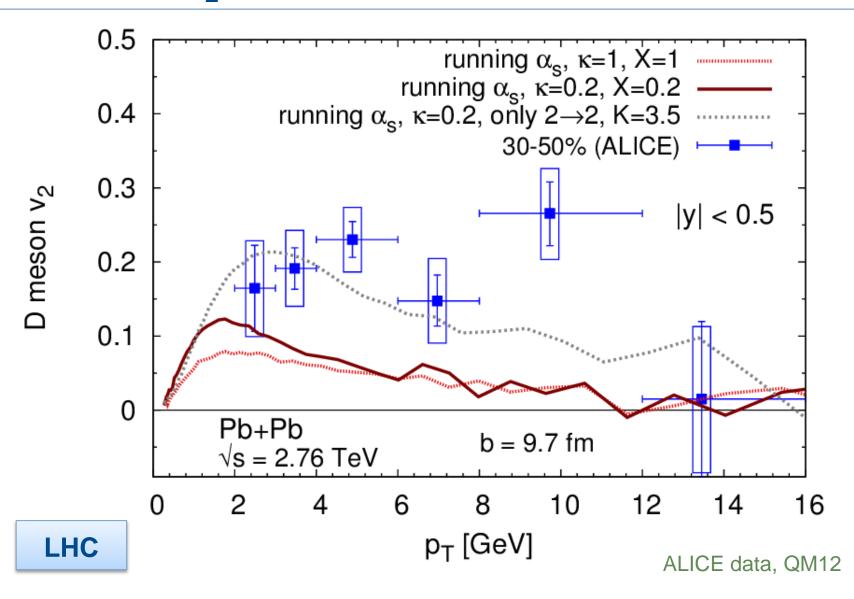
# With improved screening: D meson $R_{AA}$ at LHC





# D meson v<sub>2</sub> at LHC



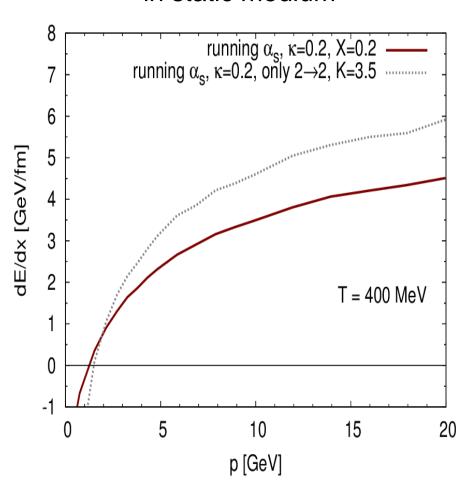


# **Energy loss and transport cross section**



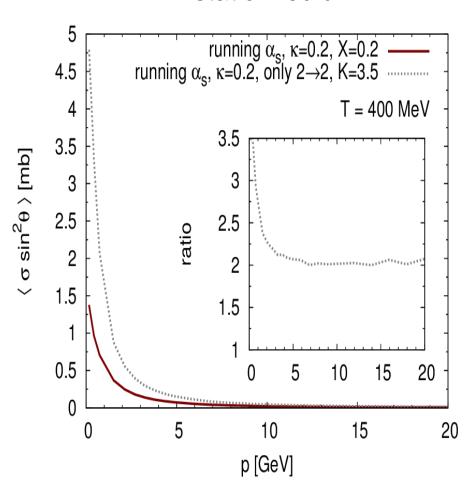
#### **Energy loss**

in static medium



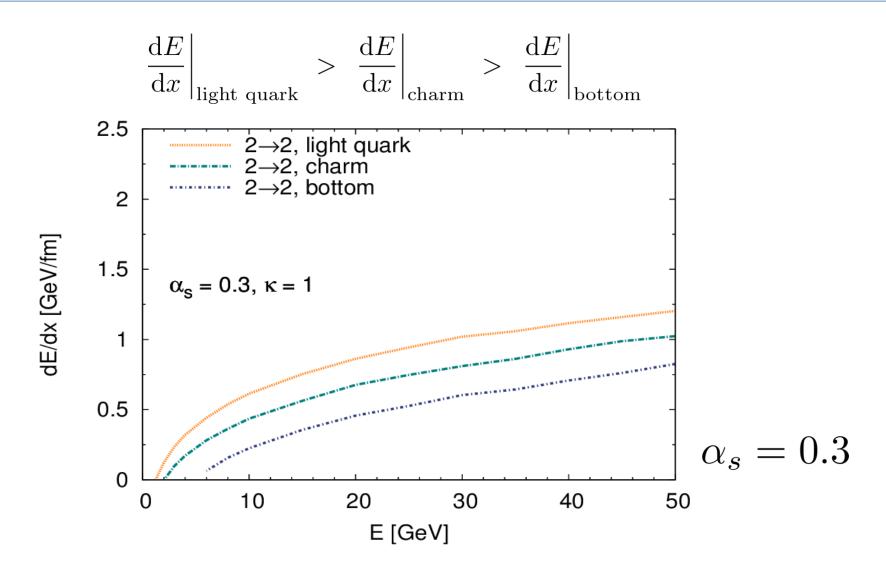
#### **Transport cross section**

in static medium



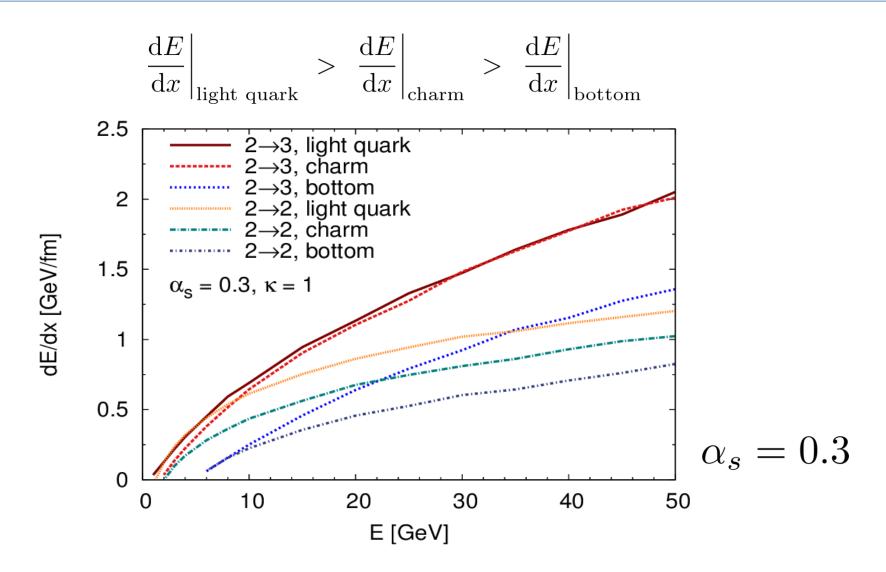
# Radiative energy loss





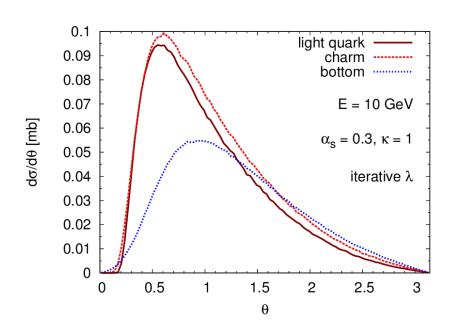
# Radiative energy loss

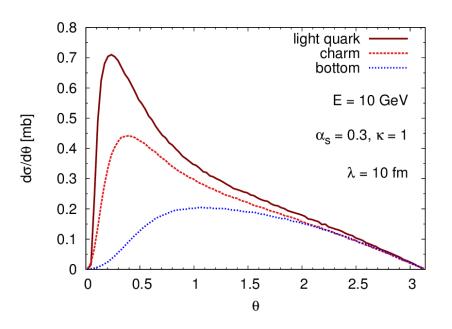




# **Angle distribution in lab frame**







With LPM

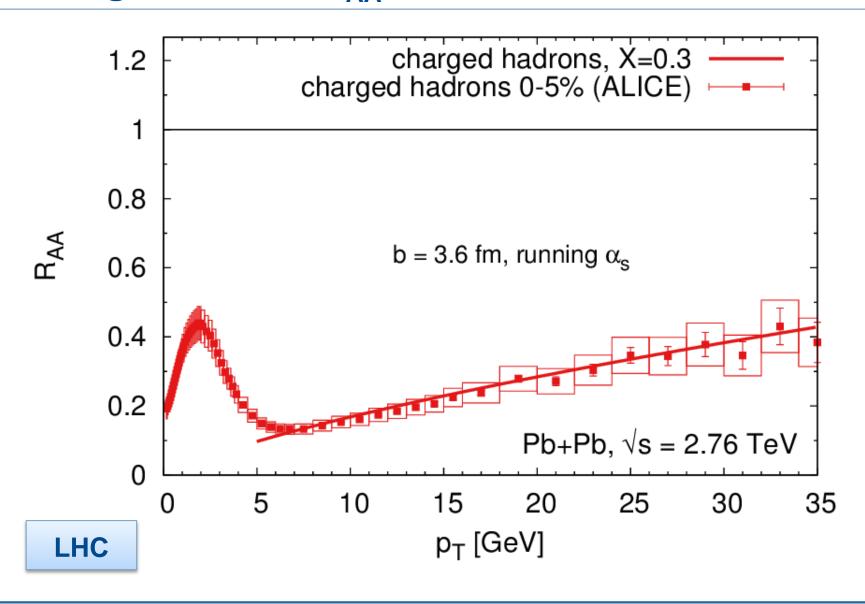
Without LPM

Dead cone due to mass is overlayed by second dead cone from LPM cut-off

Dead cone due to mass is visible

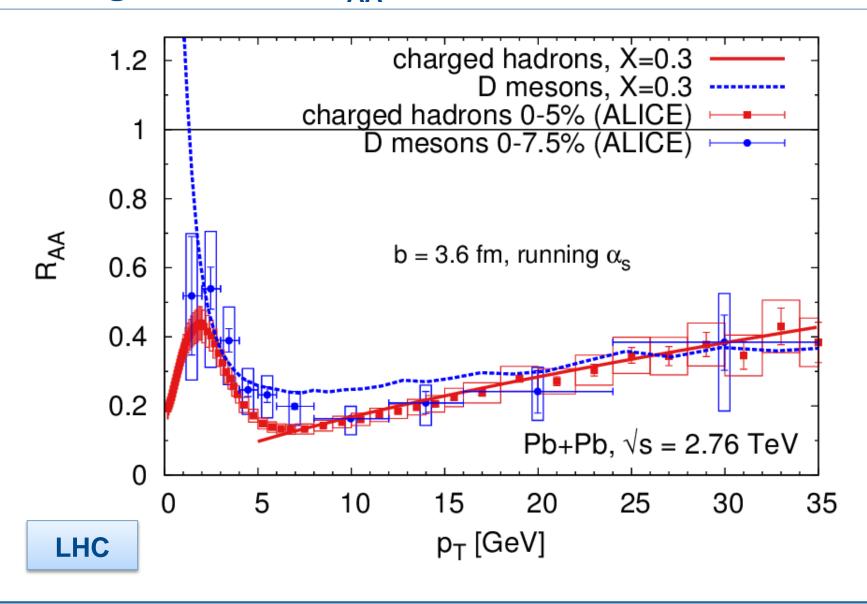
# Heavy flavor and charged hadron R<sub>AA</sub> at LHC





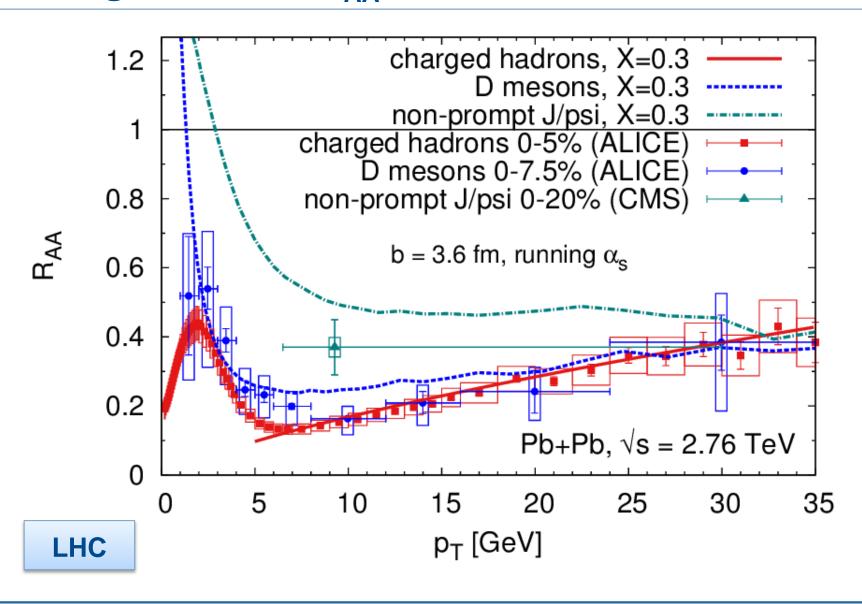
# Heavy flavor and charged hadron R<sub>AA</sub> at LHC





# Heavy flavor and charged hadron R<sub>AA</sub> at LHC





#### **Conclusions & outlook**



#### Full space-time evolution of QGP with charm and bottom quarks

#### Only binary collisions:

With running coupling and improved Debye screening, v<sub>2</sub> and R<sub>AA</sub> agreement only with K=3.5

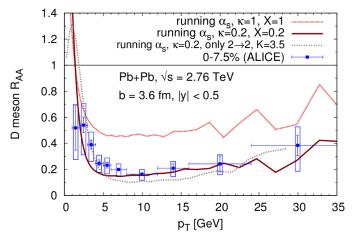
#### Radiative and binary collisions:

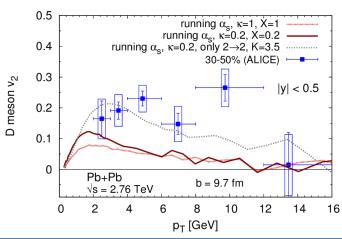
- Not enough energy loss
   → improvement of LPM
- R<sub>AA</sub> and v<sub>2</sub> simultaneously seems difficult
- R<sub>AA</sub> of light and heavy hadrons can be described

Further details in Phys. Lett. B 717, 430 (2012) and Phys. Rev. D88 (2013)

#### **Future tasks:**

Improvement of LPM effect



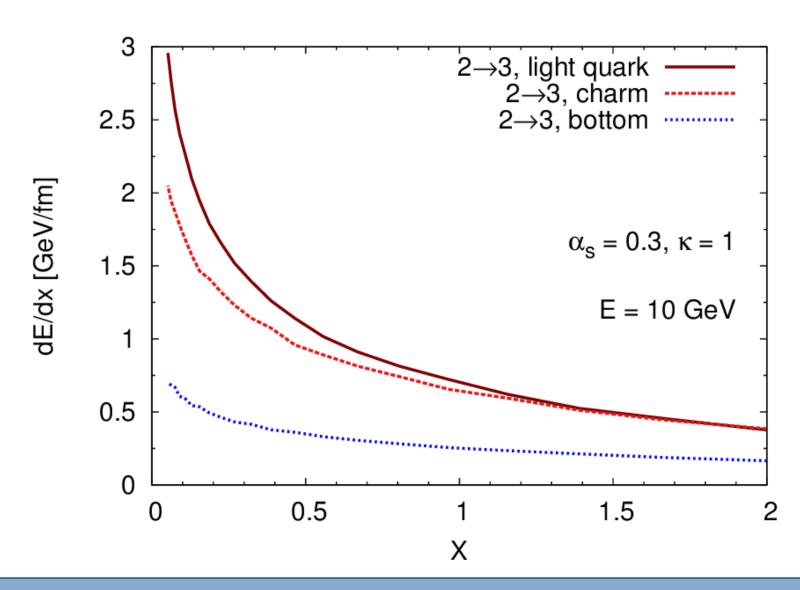




# Thank you for your attention.

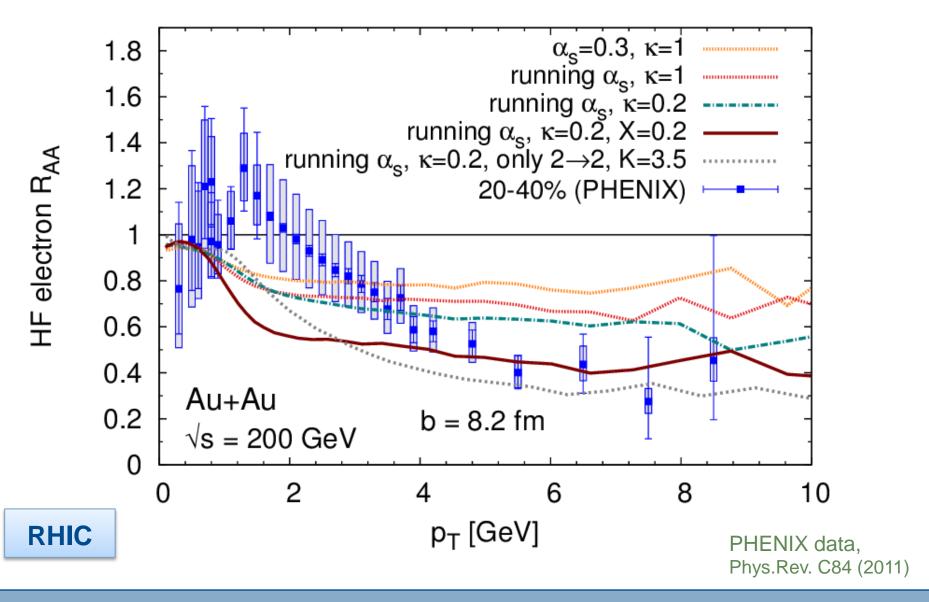
# LPM: X dependence





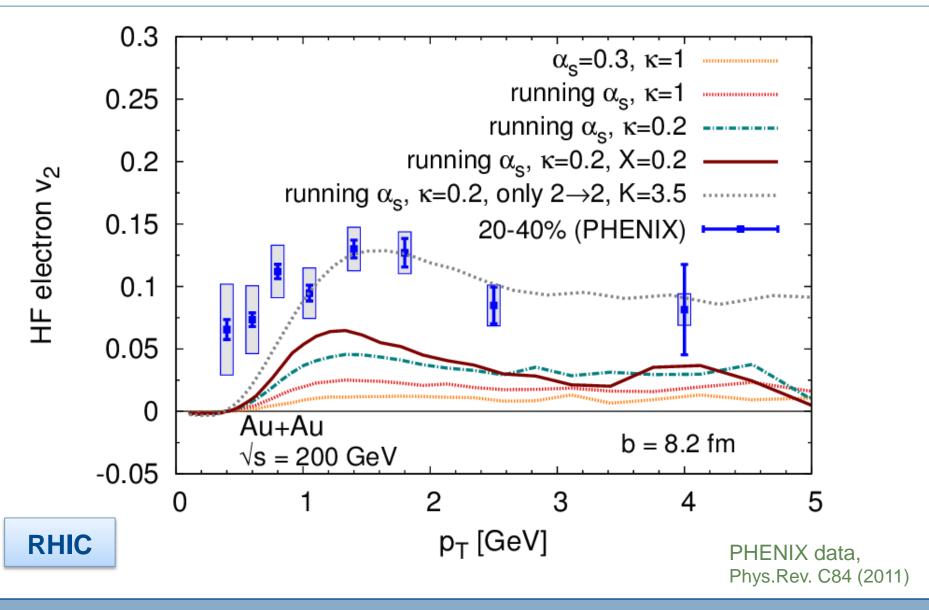
# **Heavy quark R<sub>AA</sub> at RHIC**





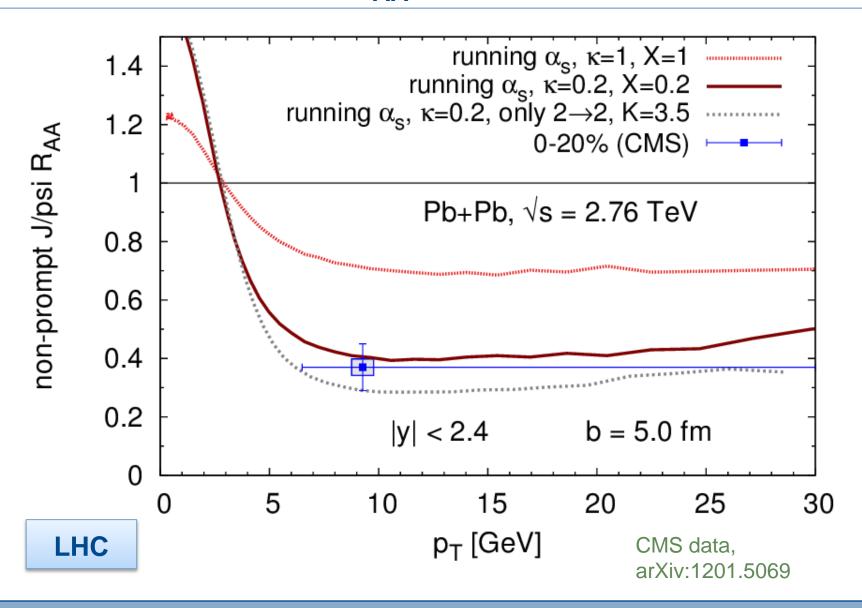
# Heavy quark v<sub>2</sub> at RHIC





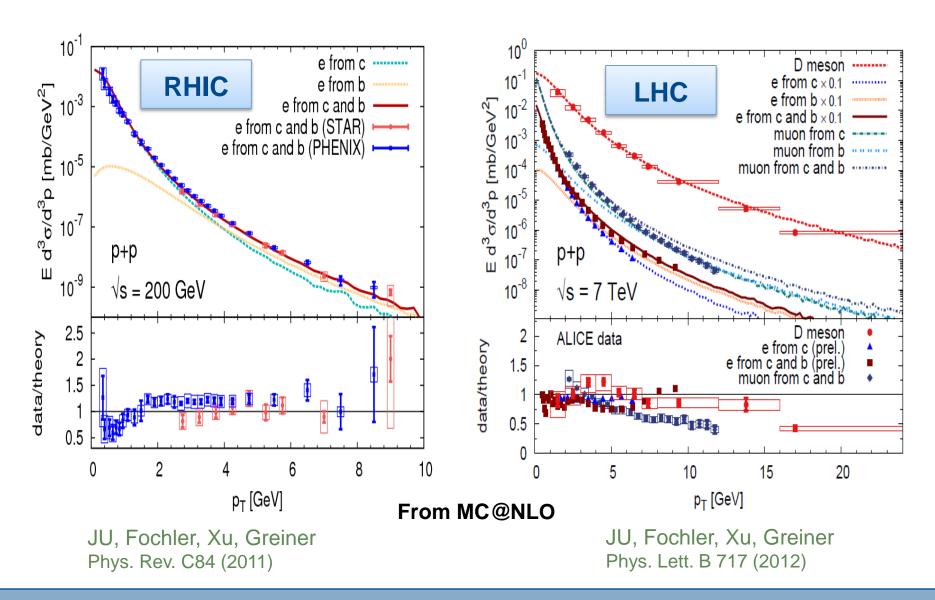
# Non-prompt J/psi R<sub>AA</sub> at LHC





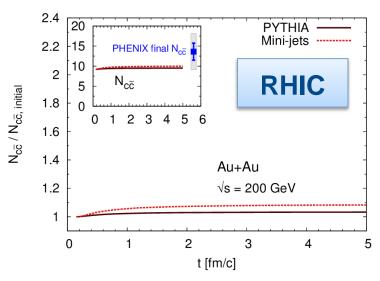
# **Initial heavy flavor spectrum**

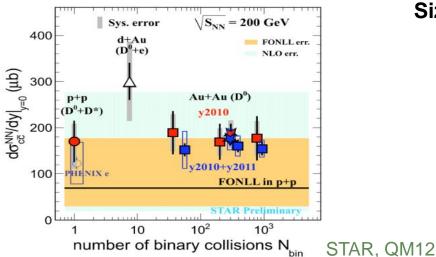




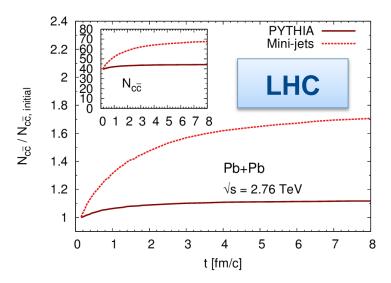
#### **Total charm production**







JU, Fochler, Xu, Greiner, Phys. Rev. C 82 (2010)

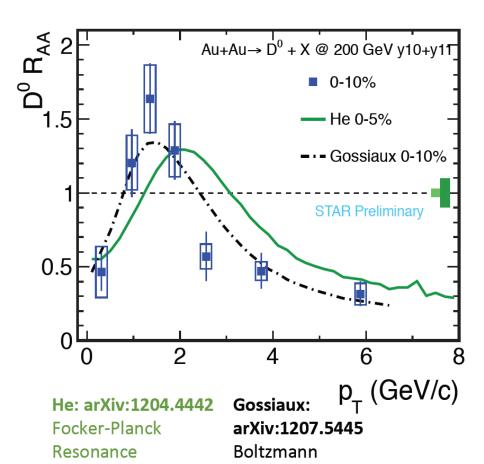


Sizeable charm production in QGP at LHC



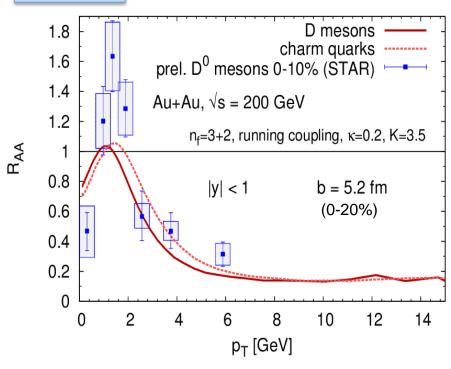
# D meson R<sub>AA</sub> from STAR





pQCD with running coupling

#### RHIC



JU, Fochler, Xu, Greiner

STAR data, QM 2012

recombination

# **Fragmentation and Decay**



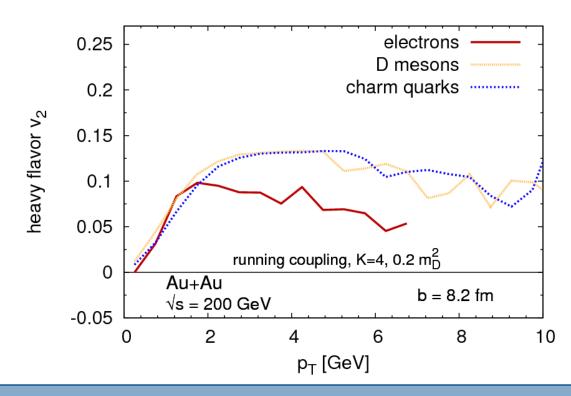
Peterson fragmentation

Peterson et al., Phys. Rev. D27 (1983)

$$D_{H/Q}(z) = \frac{N}{z \left(1 - \frac{1}{z} - \frac{\epsilon_Q}{1 - z}\right)^2} \qquad z = \frac{|\vec{p}_H|}{|\vec{p}_Q|} \qquad \epsilon_c = 0.05$$

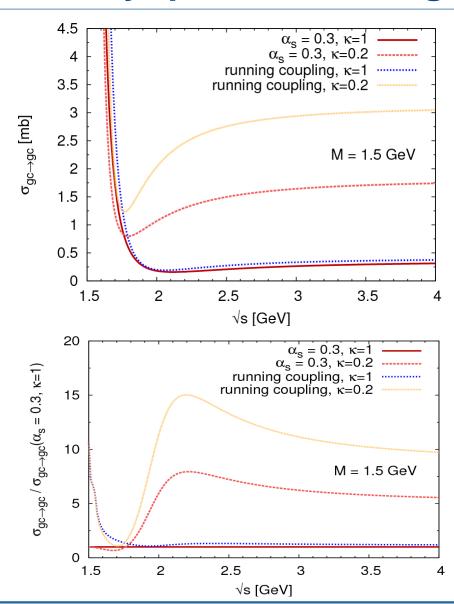
 Decay to electrons with PYTHIA

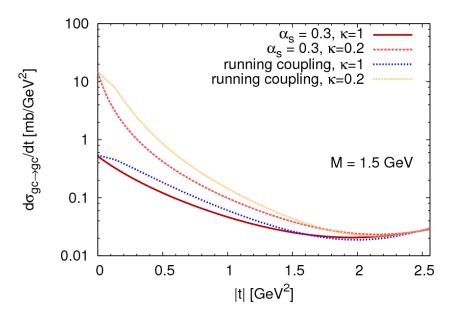
Impact of hadronization and decay small



## Heavy quark scattering cross section

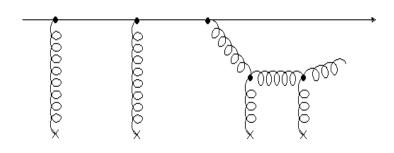






#### LPM effect vs. dead cone effect

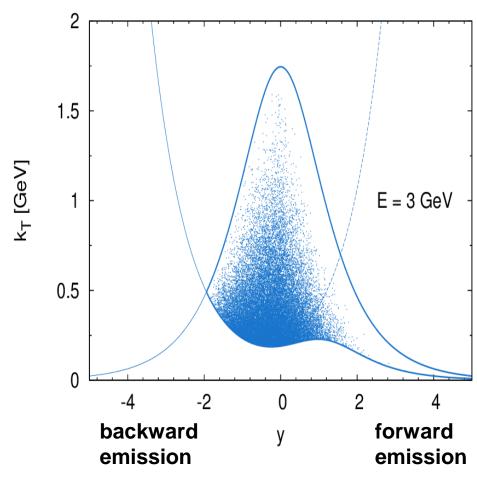




$$\lambda > \tau$$

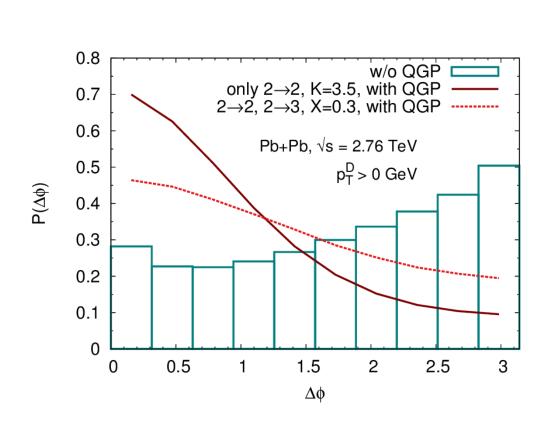
 $2 \rightarrow 3$  process only allowed if mean free path of jet larger than formation time of radiated gluon

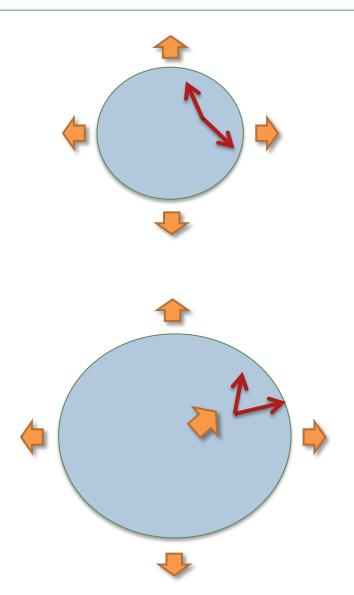




## D meson angle correlations

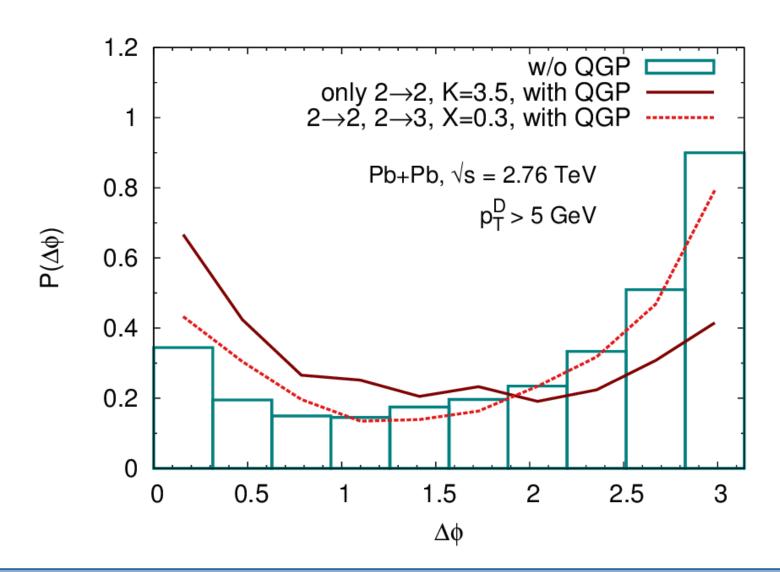






## D meson angle correlations





# **Momentum imbalance A**<sub>D</sub>



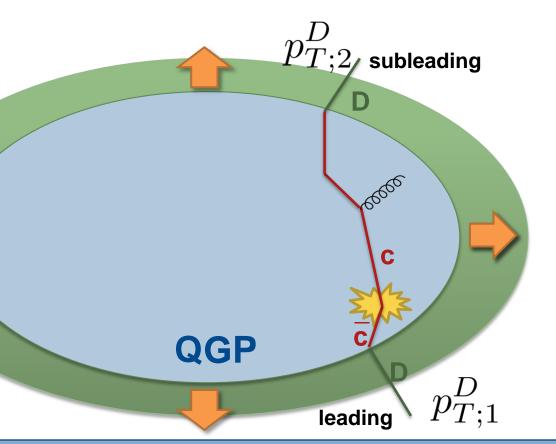
#### D meson momentum imbalance

$$A_D = \frac{p_{T;1}^D - p_{T;2}^D}{p_{T;1}^D + p_{T;2}^D}$$

No jet reconstruction necessary

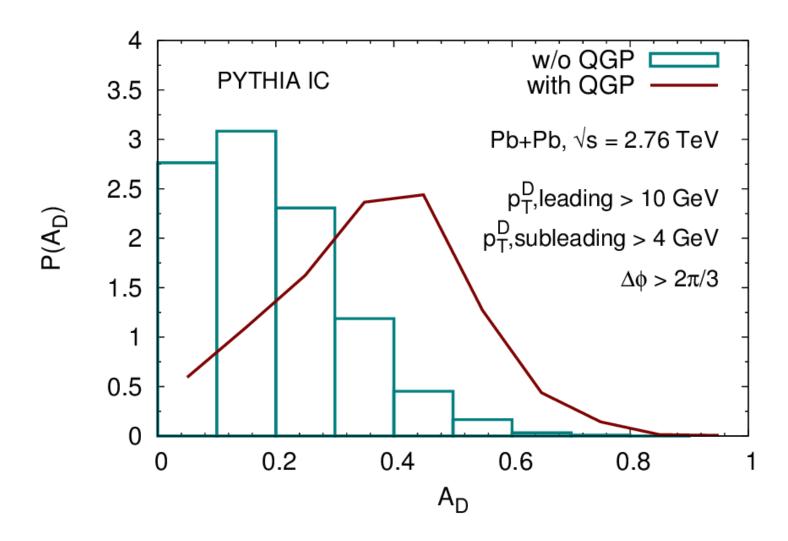
#### **Analogous to**

$$A_J = \frac{p_{T;1}^J - p_{T;2}^J}{p_{T;1}^J + p_{T;2}^J}$$



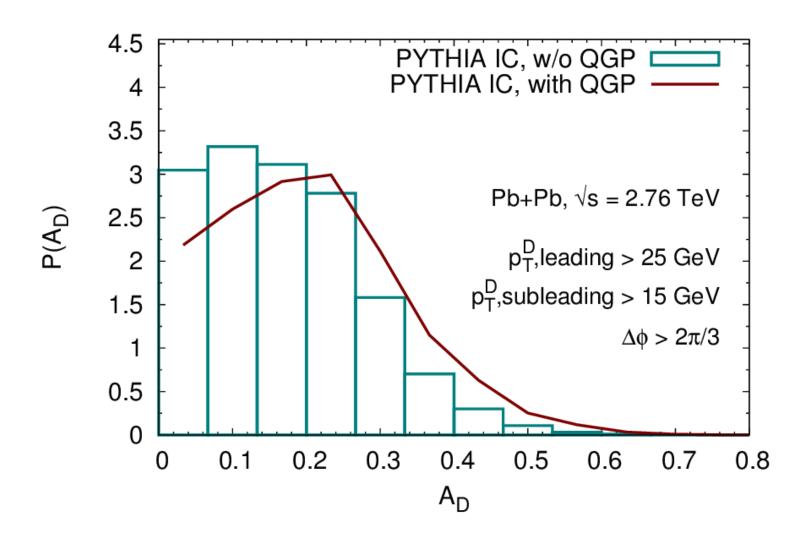
# Momentum imbalance A<sub>D</sub> for low triggers





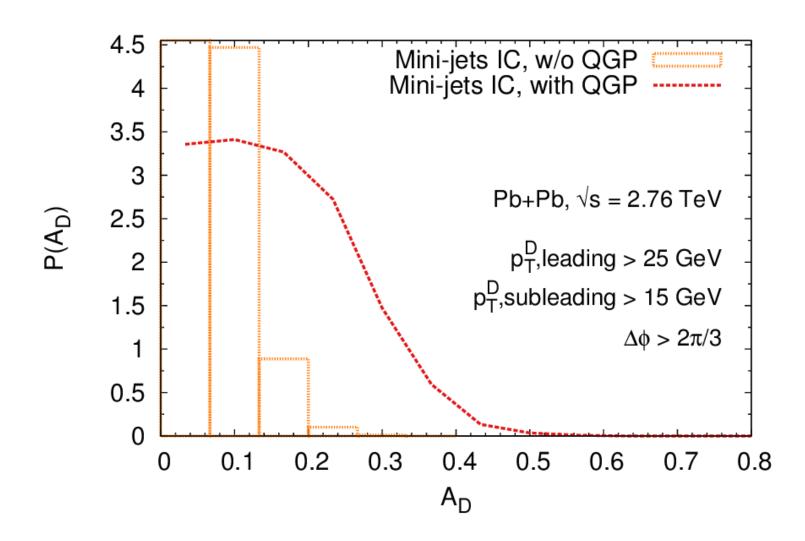
# **A**<sub>D</sub> for high triggers – PYTHIA





# **A**<sub>D</sub> for high triggers – Mini-jets





## Length imbalance



