

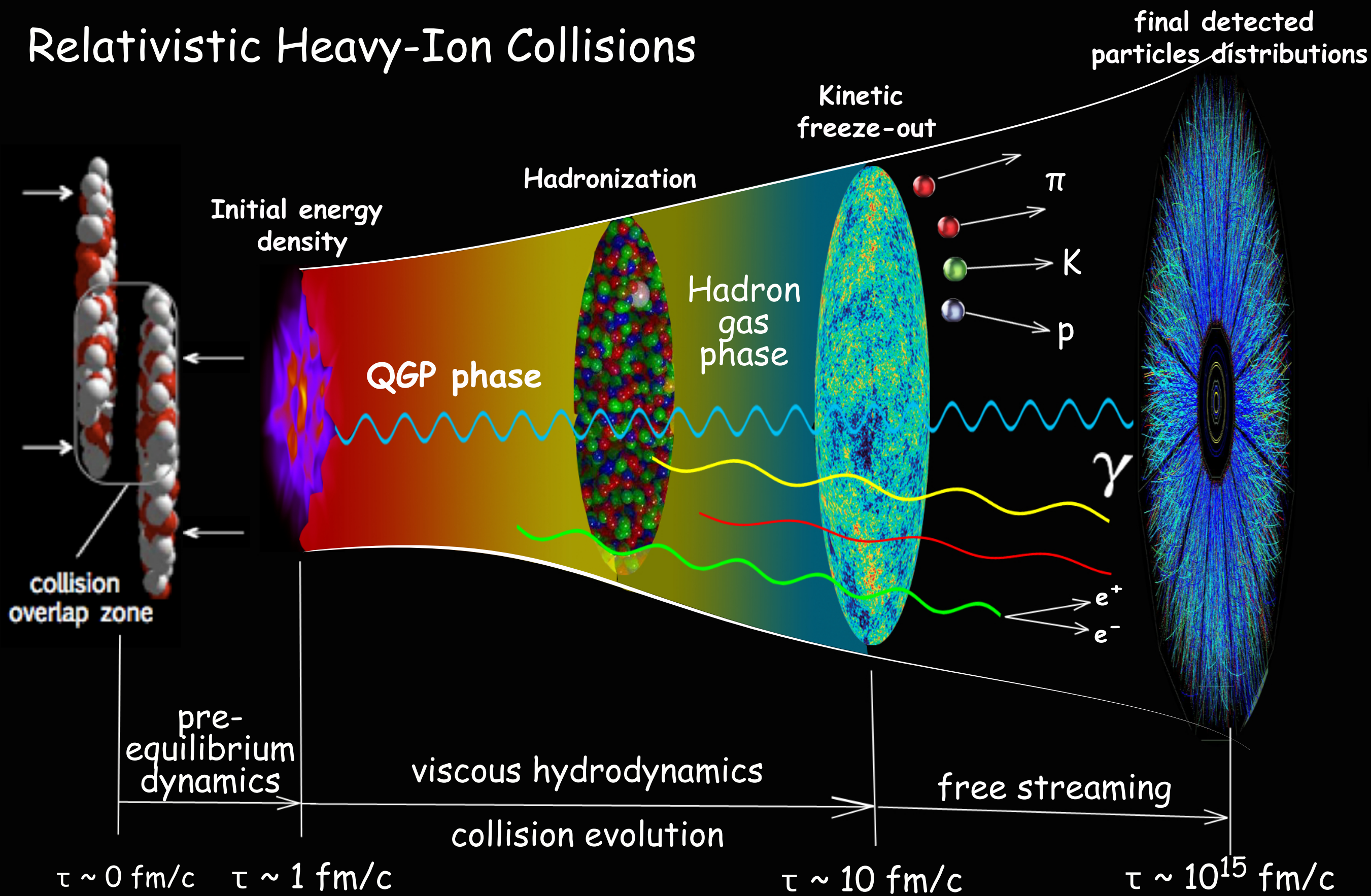
Anisotropic flow of thermal photons as a quark-gluon plasma viscometer

Chun Shen
The Ohio State University

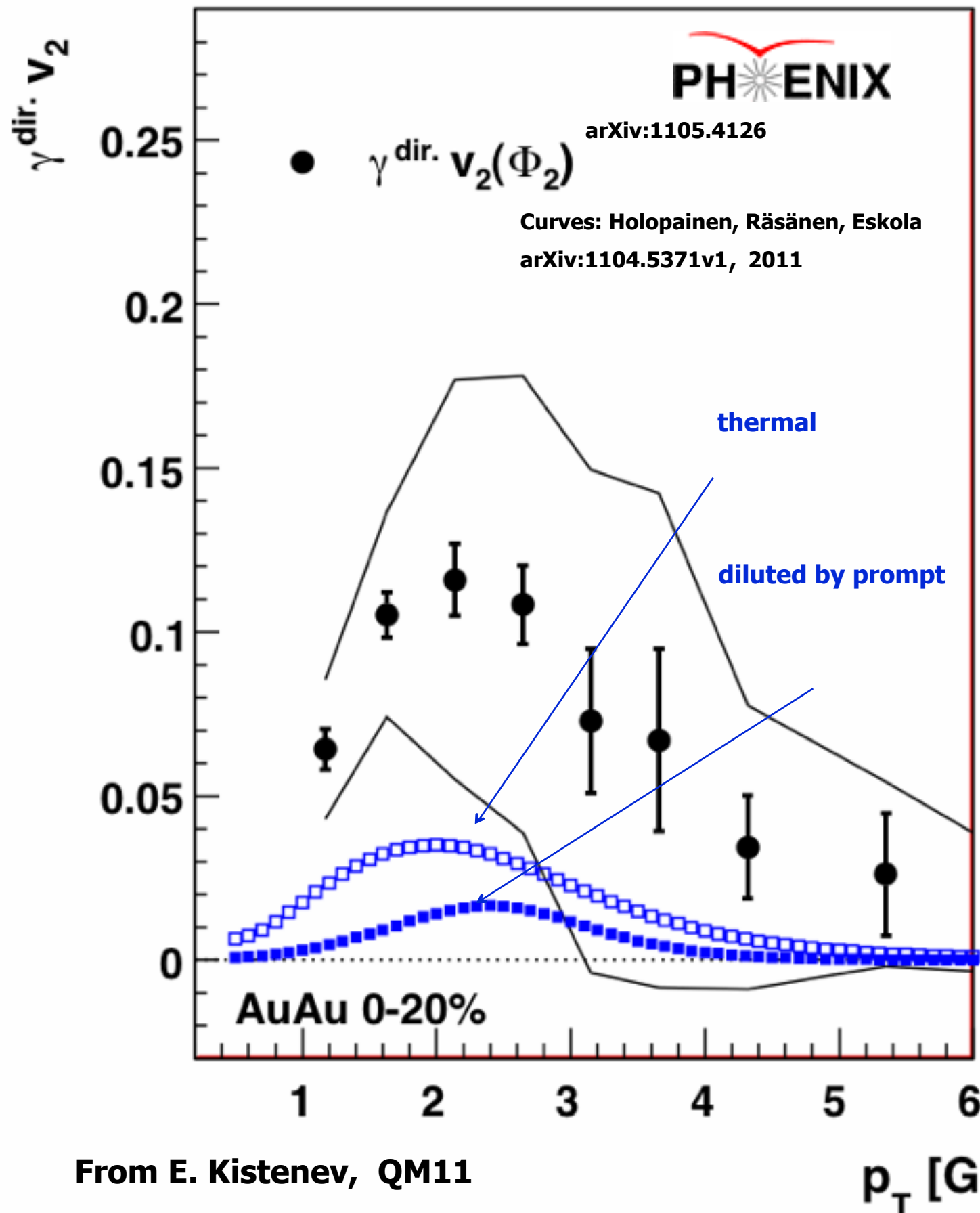
In collaboration with Jean-Francois Paquet,
Ulrich Heinz, and Charles Gale

Little Bang

Relativistic Heavy-Ion Collisions



Challenge from Experiment



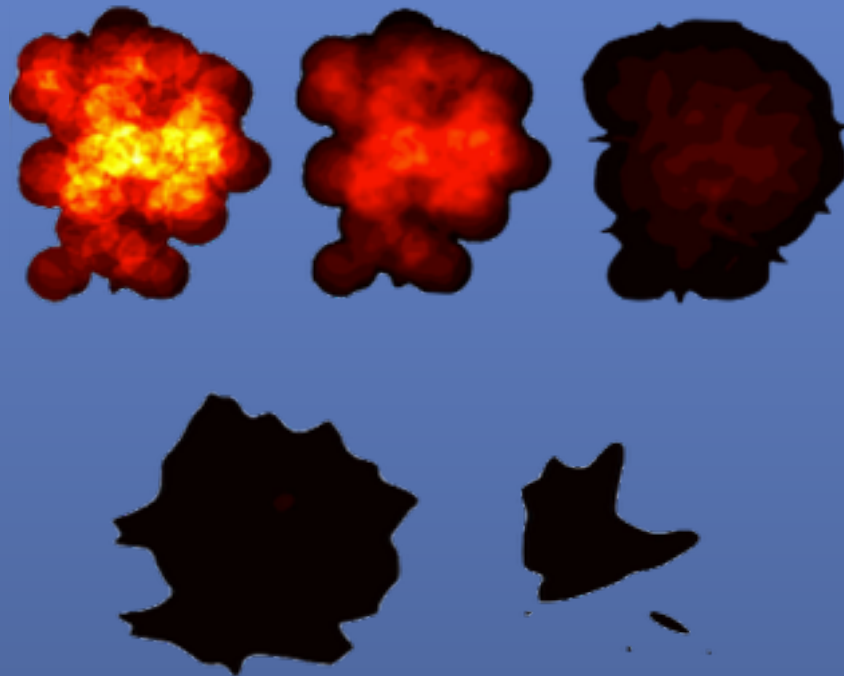
- PHENIX measurements show **large** direct photon v_2 at $p_T < 4$ GeV
- The state-of-the-art calculation underestimates the data by a factor of **5!**

Work flow of Theoretical Calculation

Initial Condition
Generators
(MC-KLN, MC-Glauber)

Thermal Photon
Emission Rates

Hydrodynamic
Simulations
(VISH2+1)



HydroInfo
Package

$$e, s, p, T,$$

$$u^\mu, \pi^{\mu\nu}$$

Thermal Photon
Interface

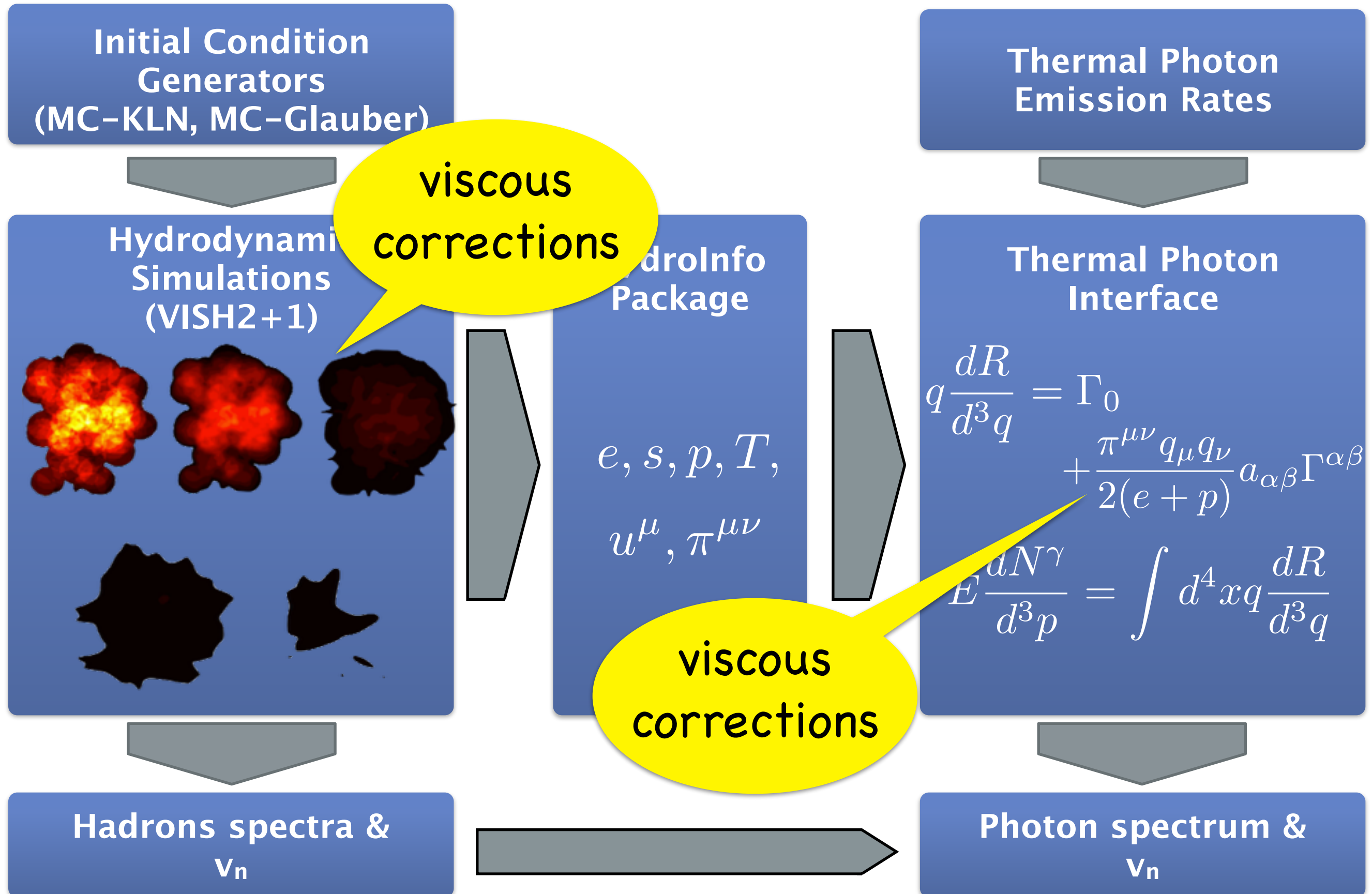
$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} q_\mu q_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

$$E \frac{dN^\gamma}{d^3p} = \int d^4x q \frac{dR}{d^3q}$$

Hadrons spectra &
 V_n

Photon spectrum &
 V_n

Work flow of Theoretical Calculation



General Formalism

Thermal photon emission rates can be calculated by

$$E_q \frac{dR}{d^3q} = \int \frac{d^3p_1}{2E_1(2\pi)^3} \frac{d^3p_2}{2E_2(2\pi)^3} \frac{d^3p_3}{2E_3(2\pi)^3} \frac{1}{2(2\pi)^3} |\mathcal{M}|^2 \\ \times f_1(p_1^\mu) f_2(p_2^\mu) (1 \pm f_3(p_3^\mu)) (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - q)$$

With

$$f(p^\mu) = f_0(E) + f_0(E)(1 \pm f_0(E)) \frac{\pi^{\mu\nu} \hat{p}_\mu \hat{p}_\nu}{2(e + p)} \chi\left(\frac{p}{T}\right)$$

We can expand photon emission rates around the thermal equilibrium:

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta},$$

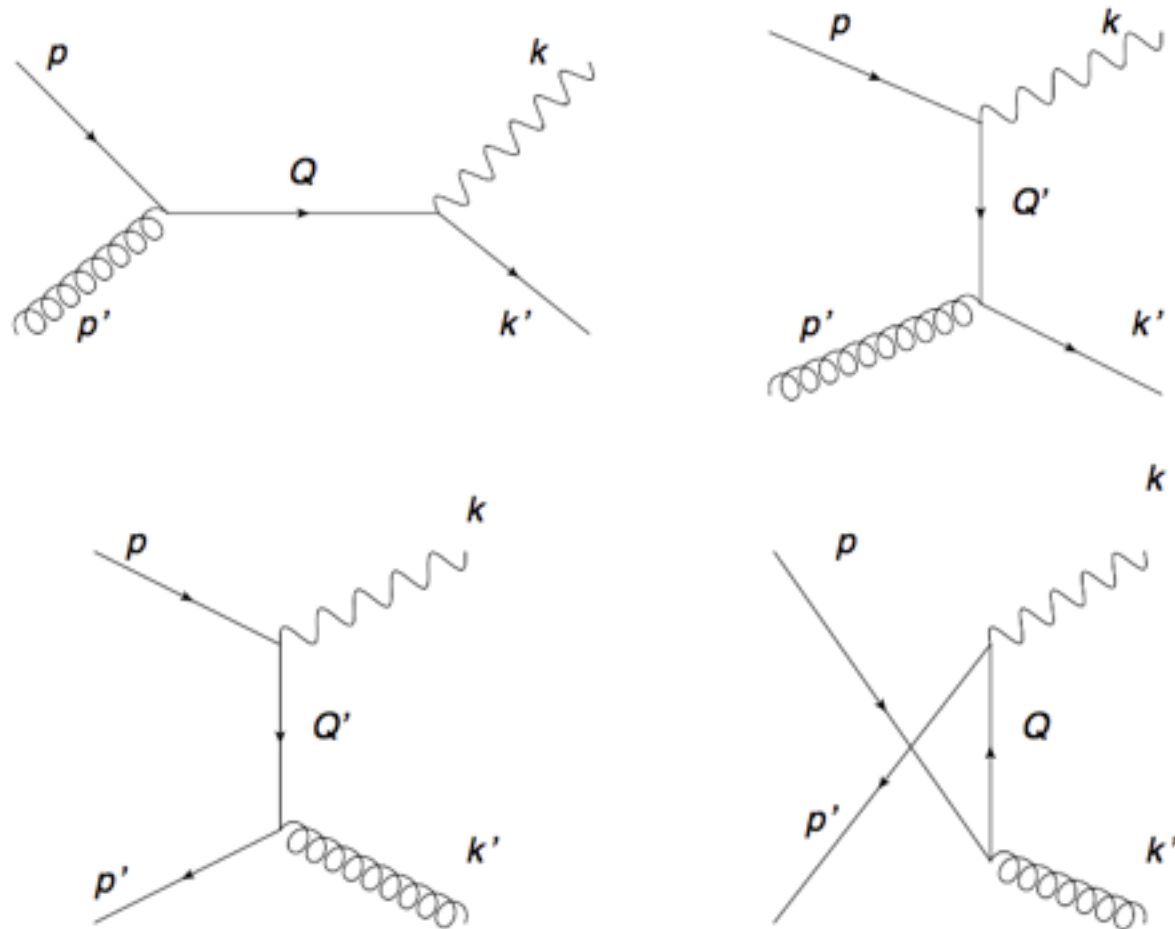
$$a_{\mu\nu} = \frac{3}{2(u \cdot \hat{q})^4} \hat{q}_\mu \hat{q}_\nu + \frac{1}{(u \cdot \hat{q})^2} u_\mu u_\nu + \frac{1}{2(u \cdot \hat{q})^2} g_{\mu\nu} - \frac{3}{2(u \cdot \hat{q})^3} (\hat{q}_\mu u_\nu + \hat{q}_\nu u_\mu).$$

General Formalism

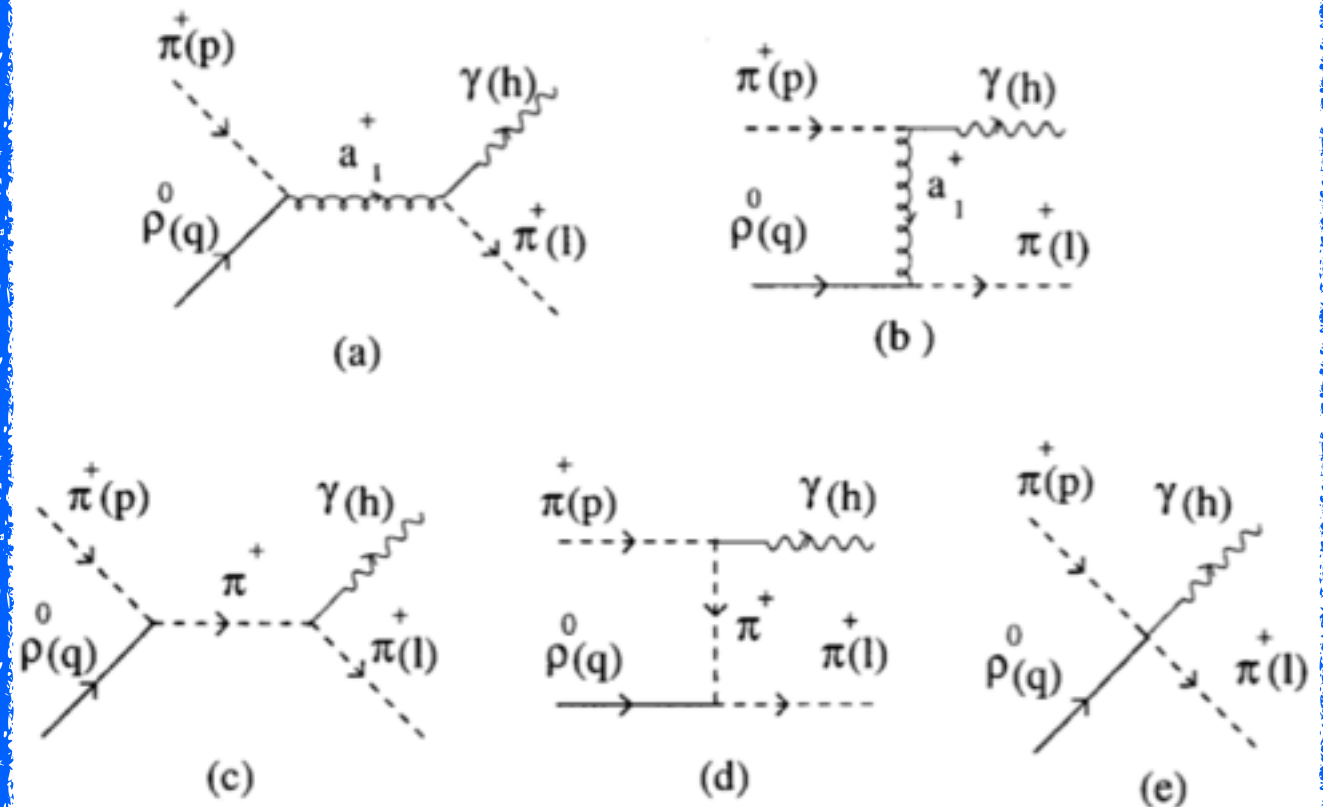
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Equilibrium rates

QGP



Hadron Gas



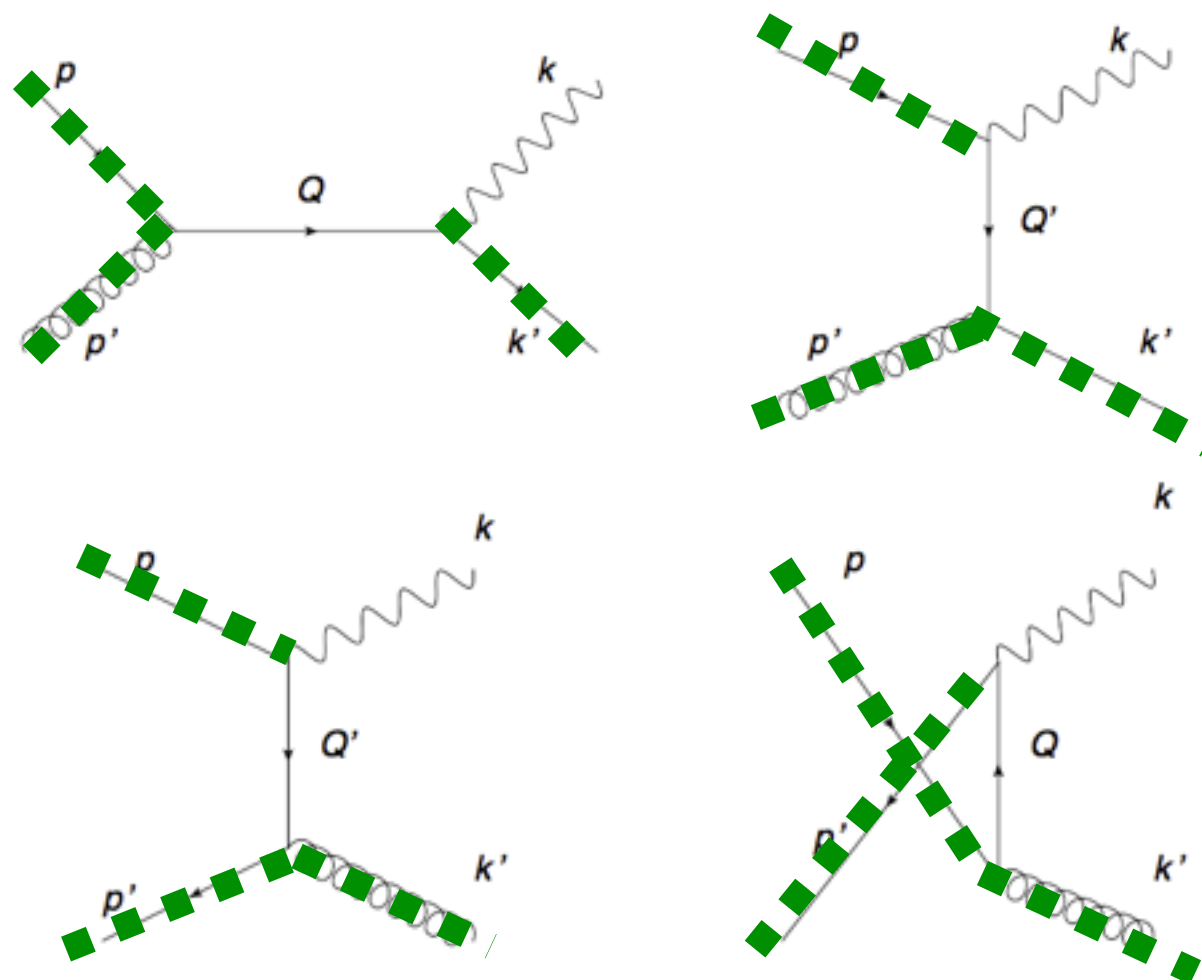
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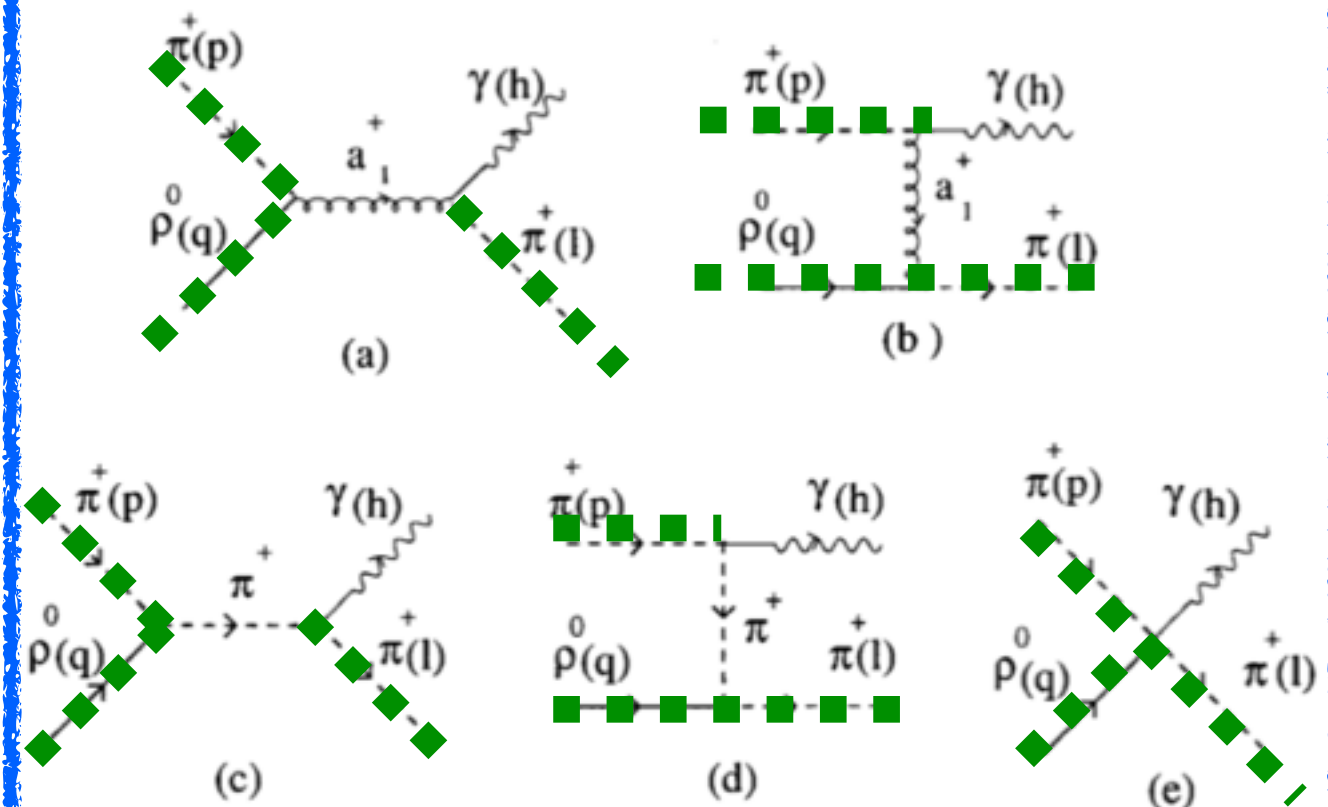
Equilibrium rates

off-equilibrium δf corrections

QGP



Hadron Gas



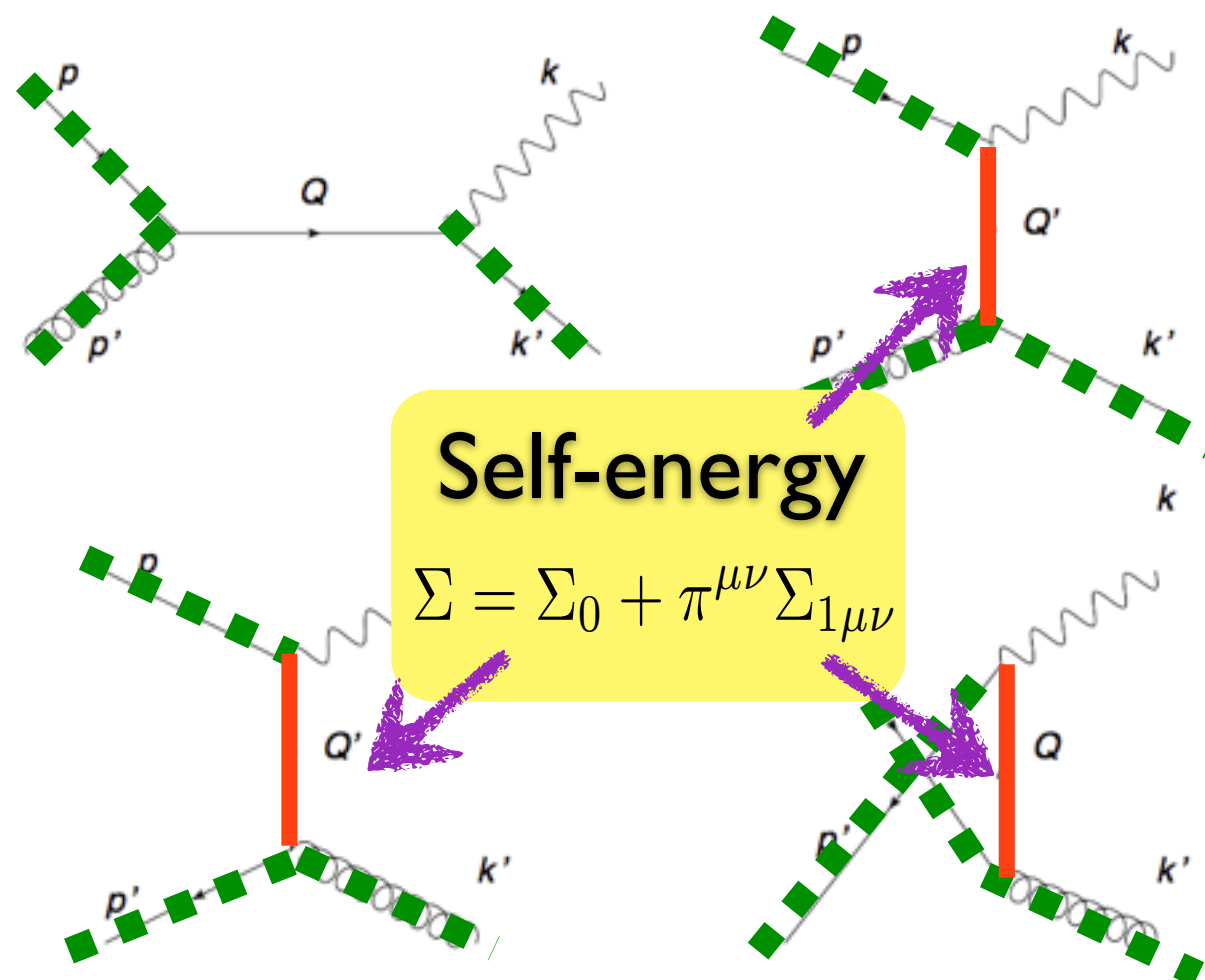
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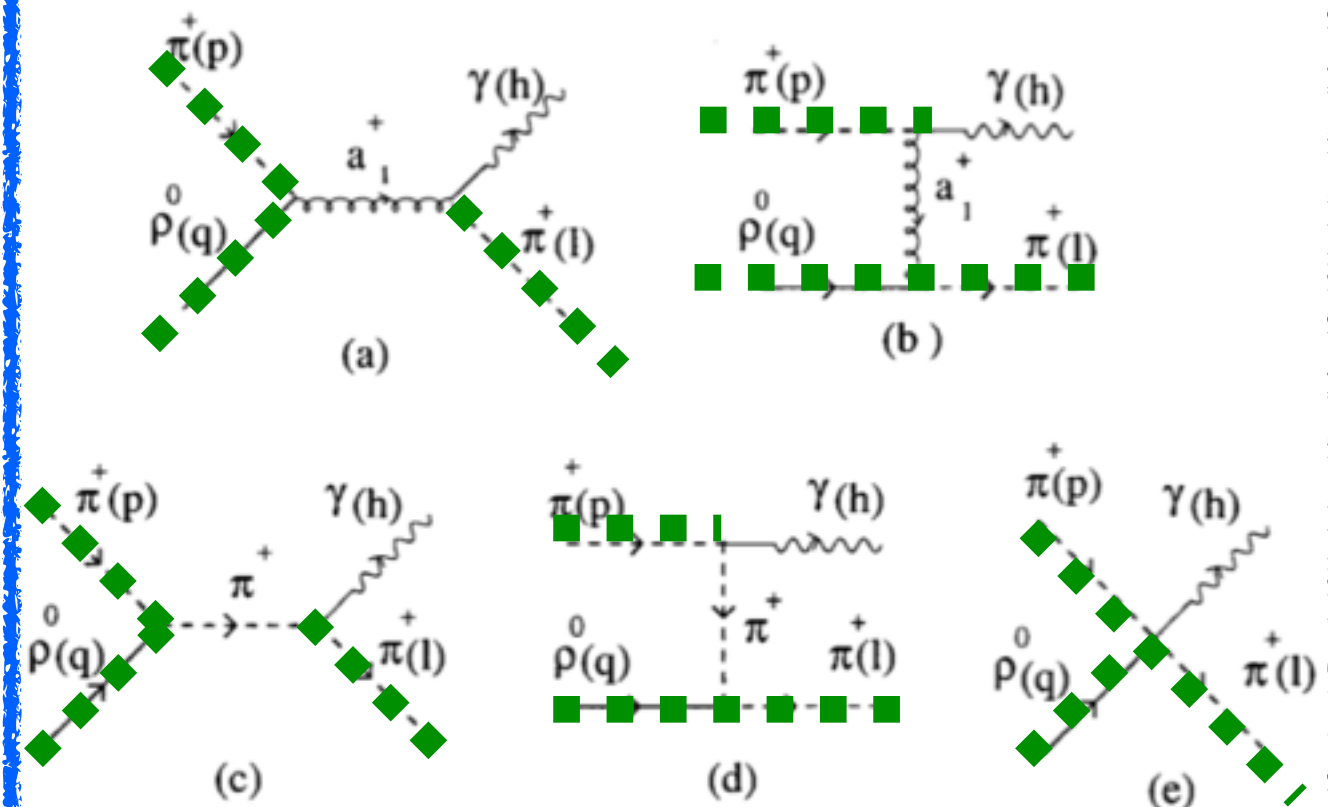
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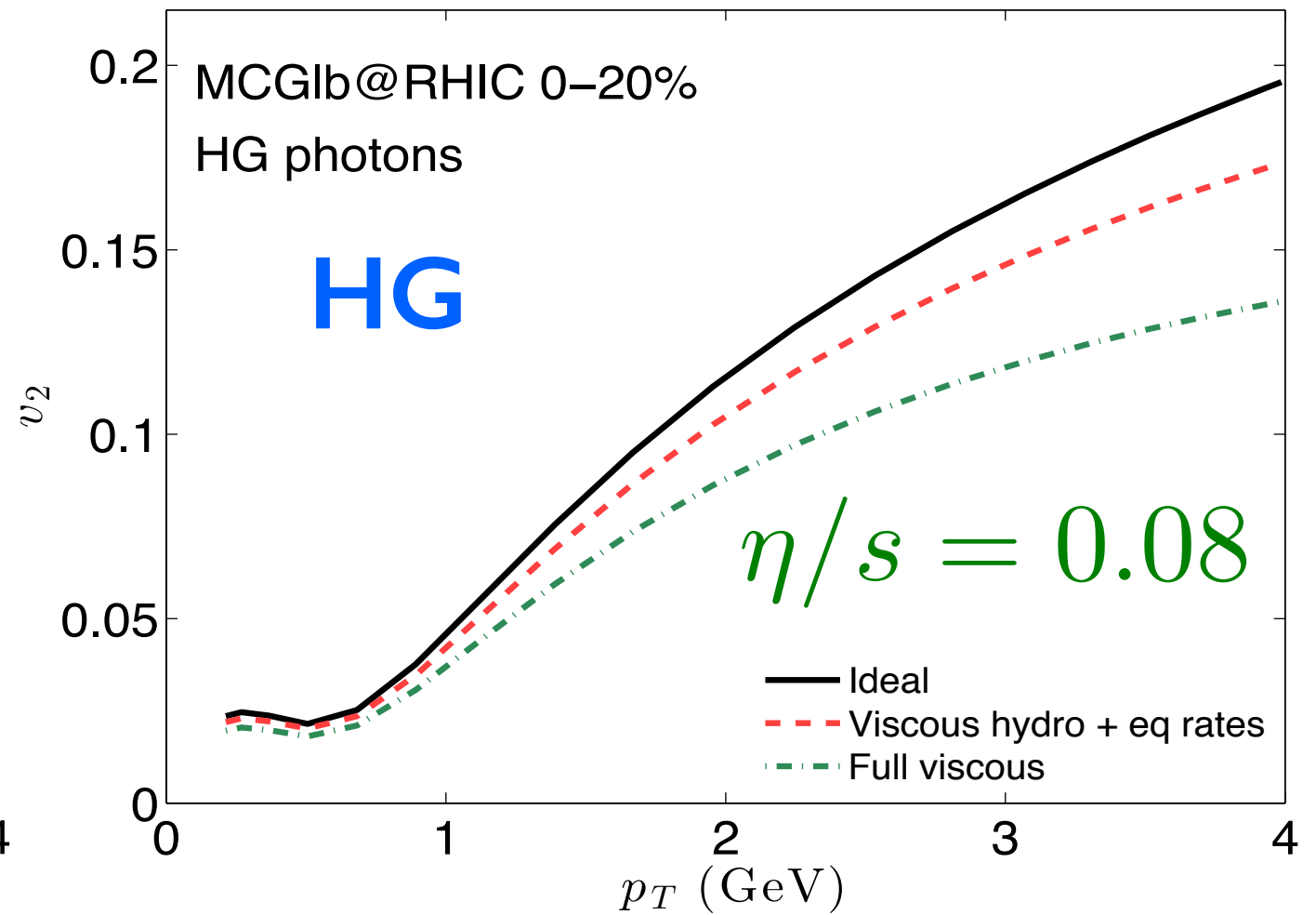
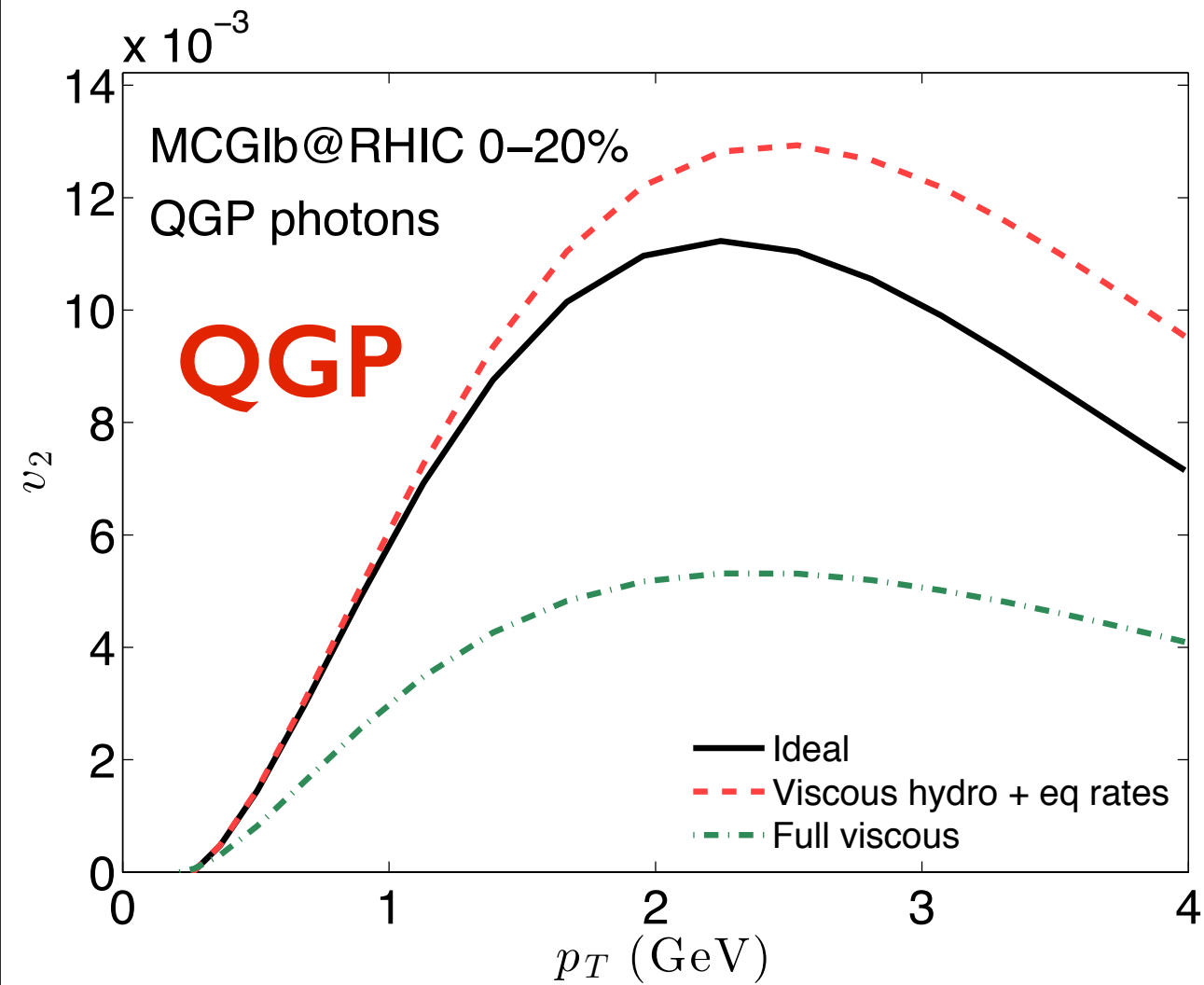
QGP



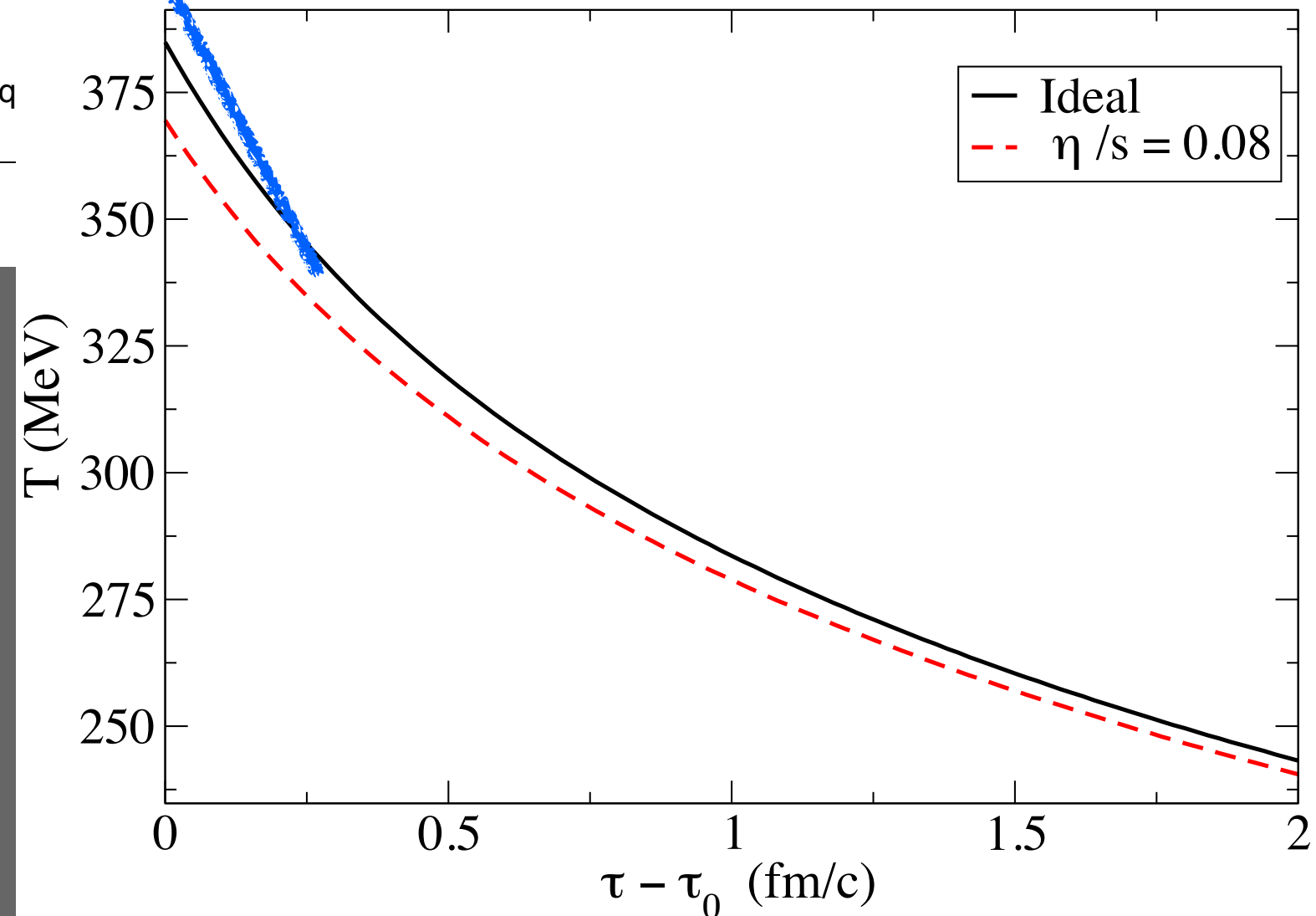
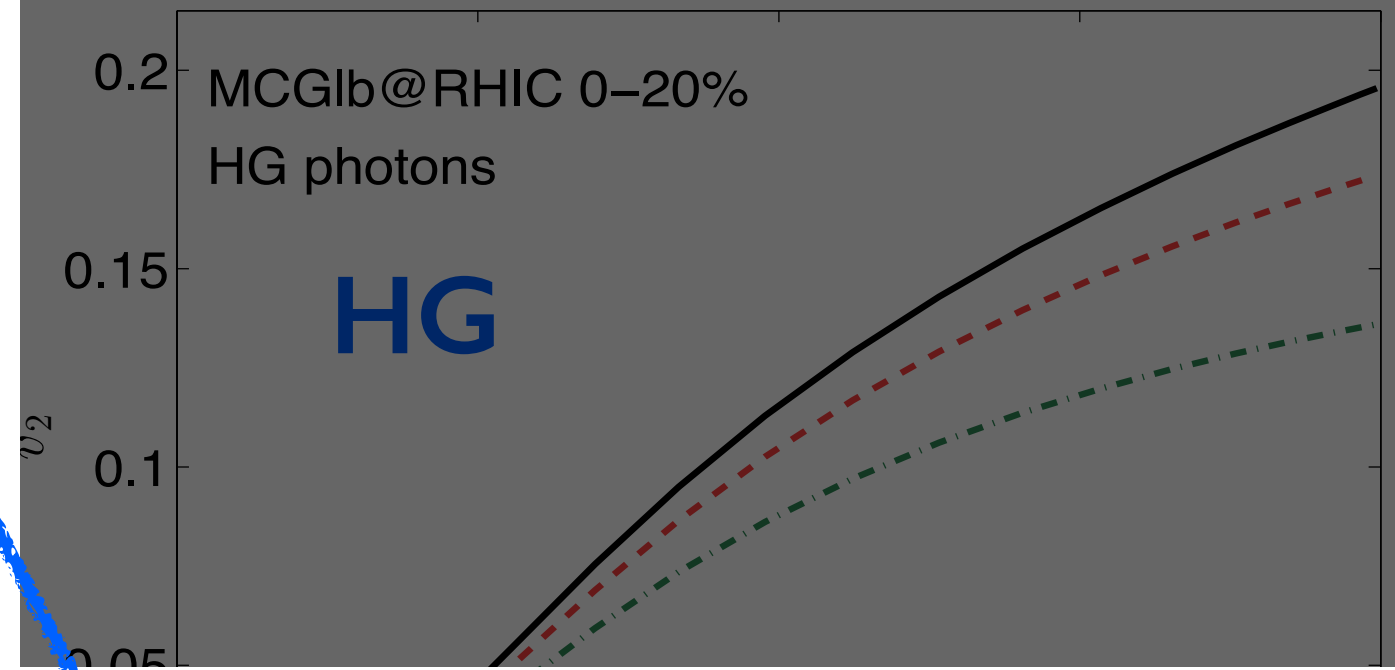
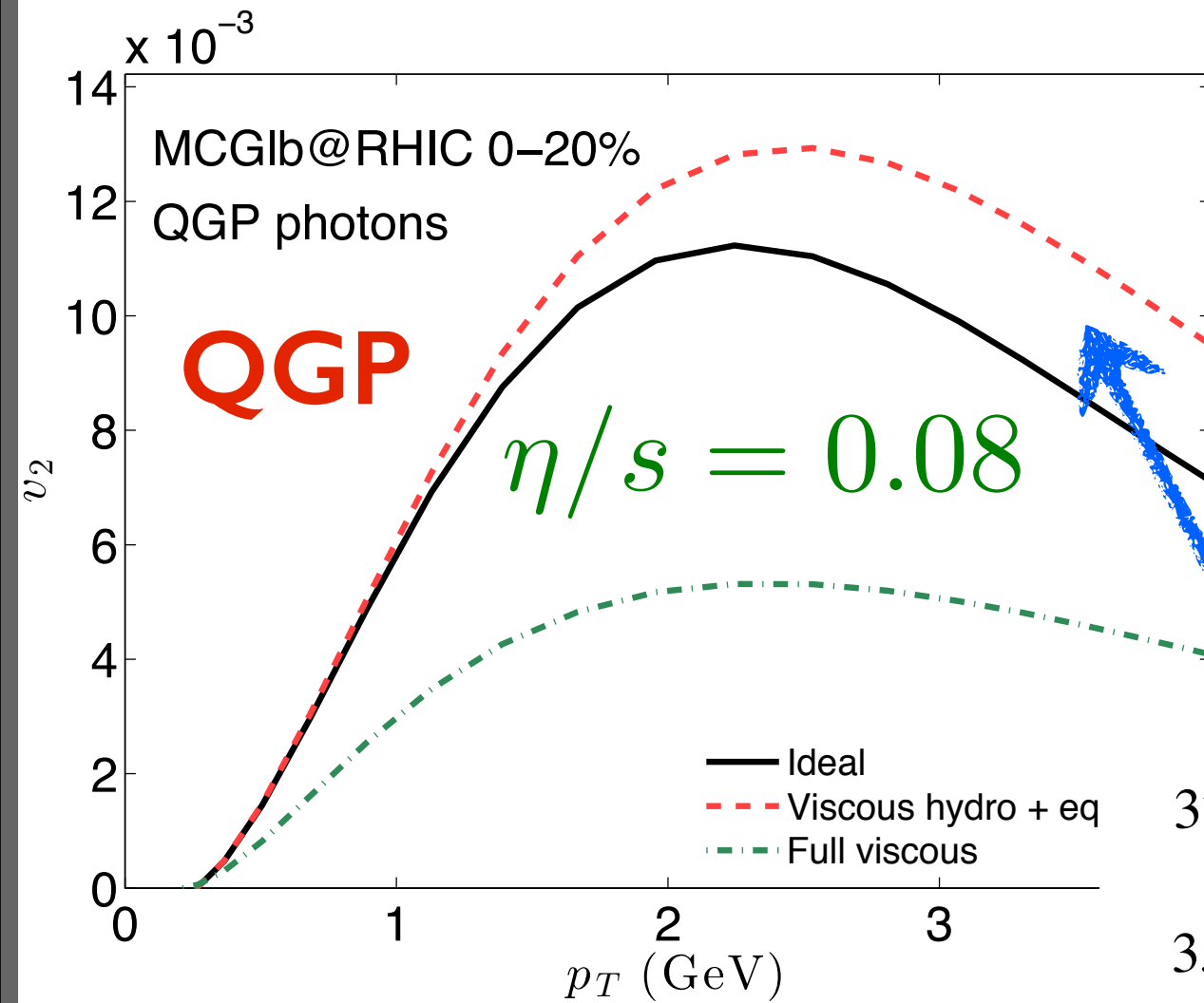
Hadron Gas



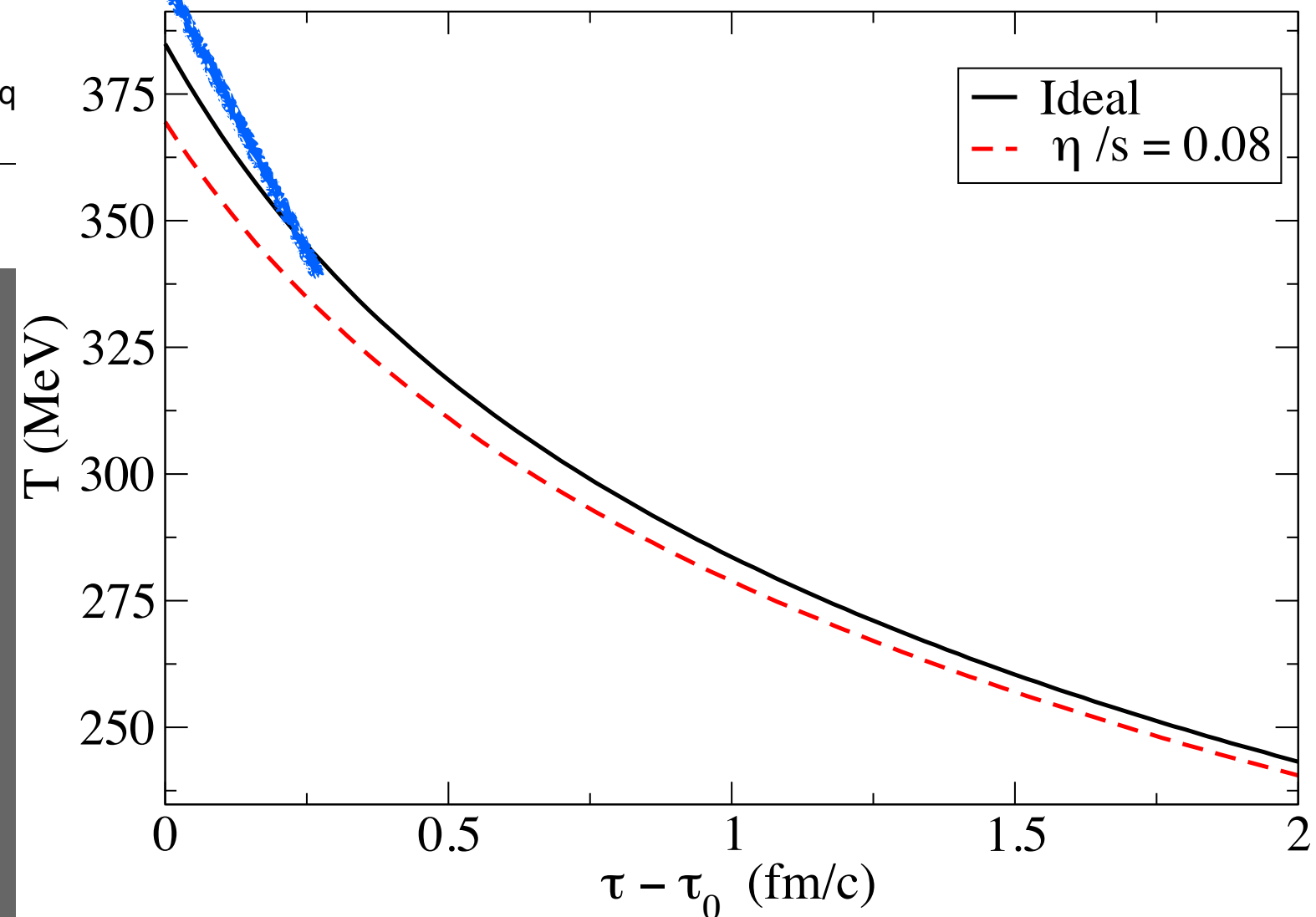
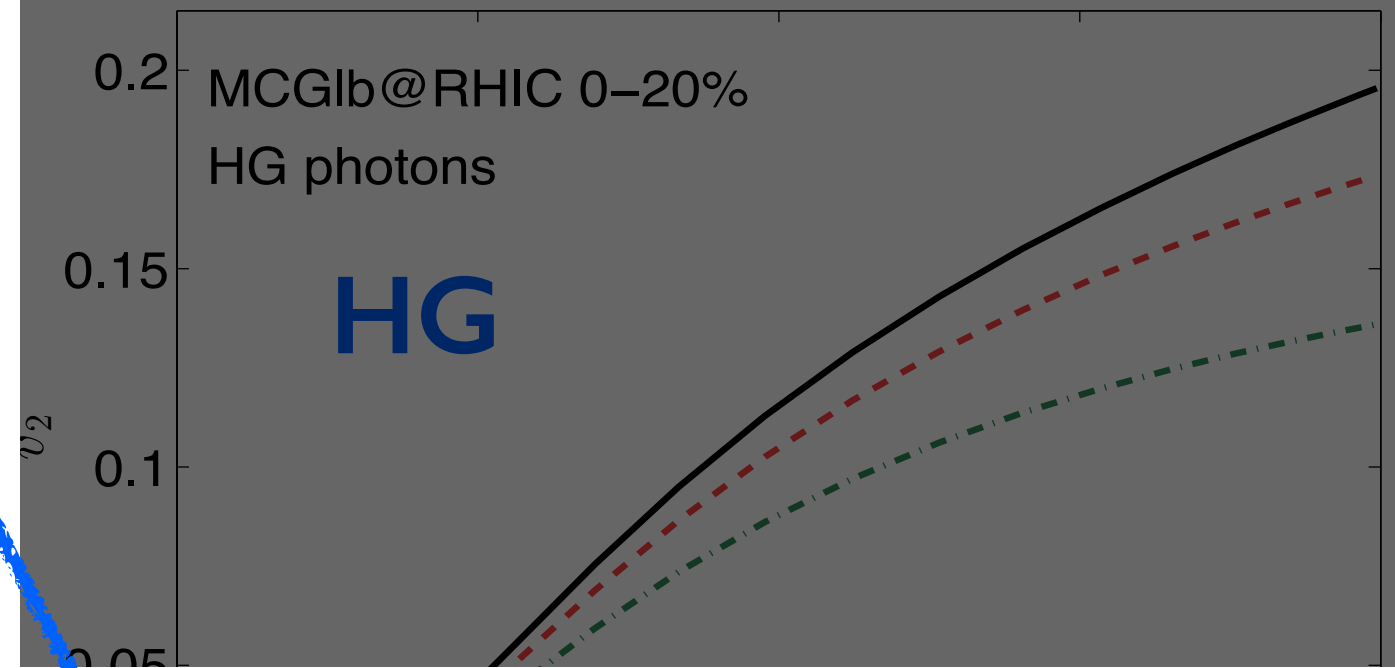
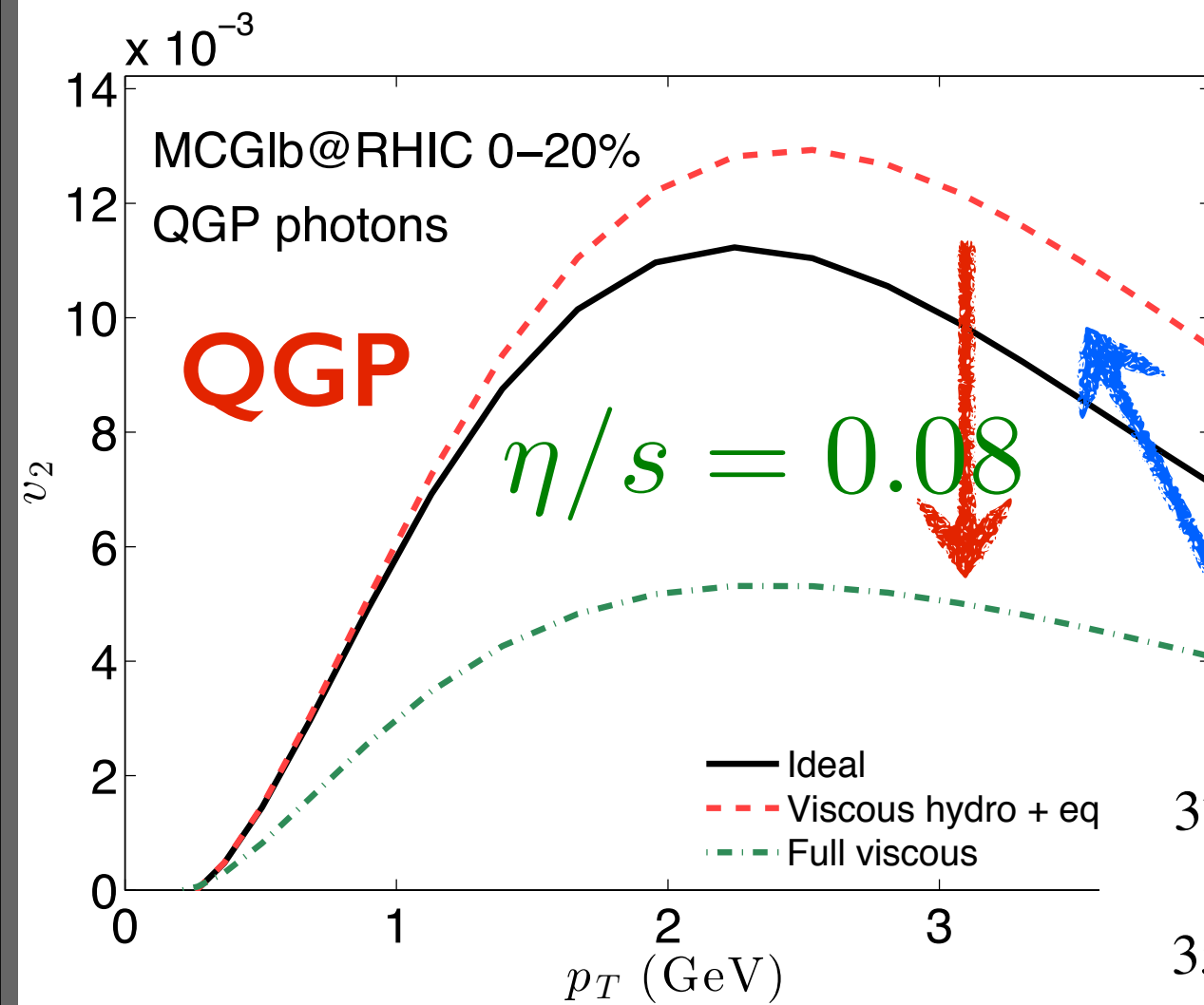
Viscous effects on photon elliptic flow



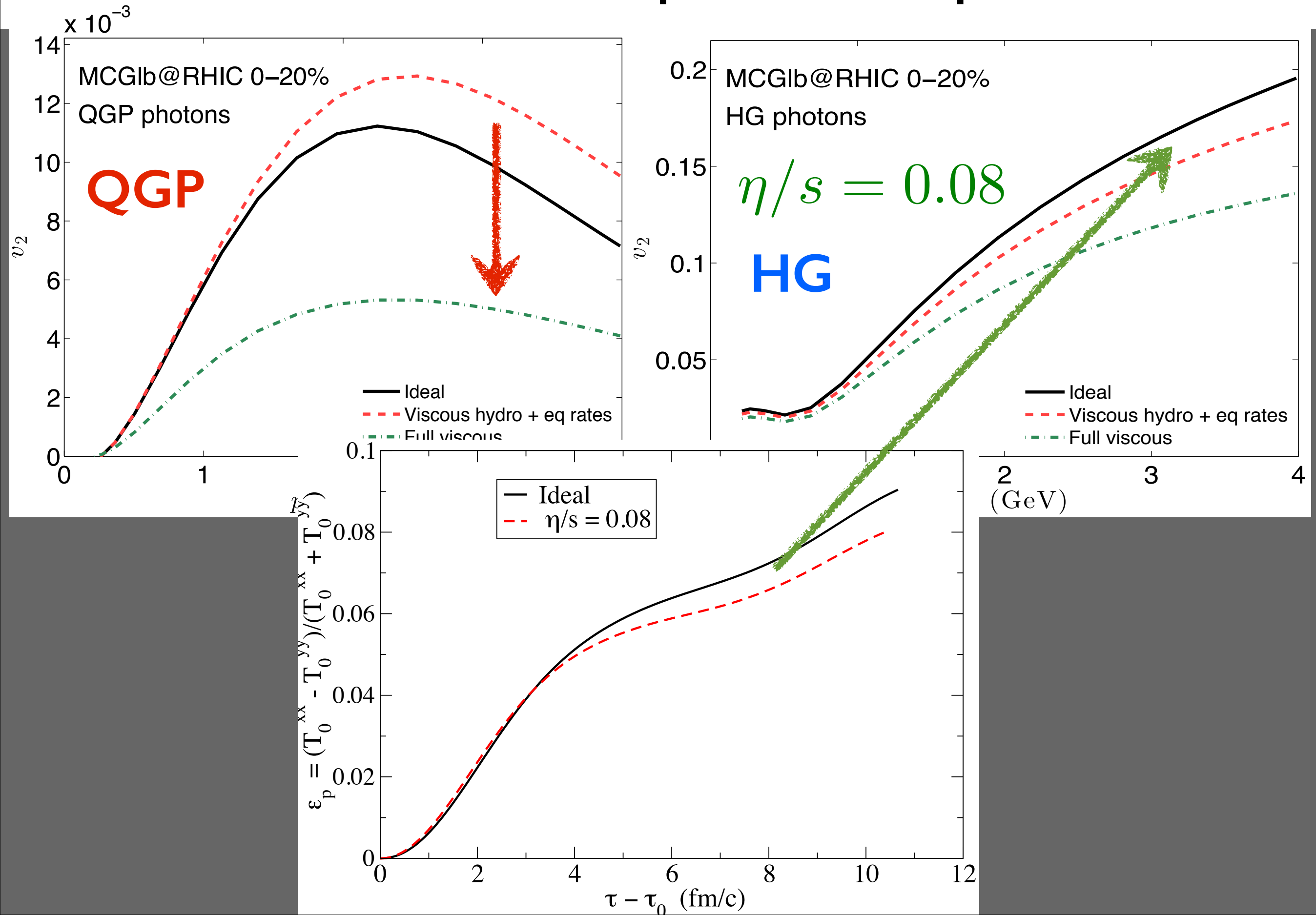
Viscous effects on photon elliptic flow



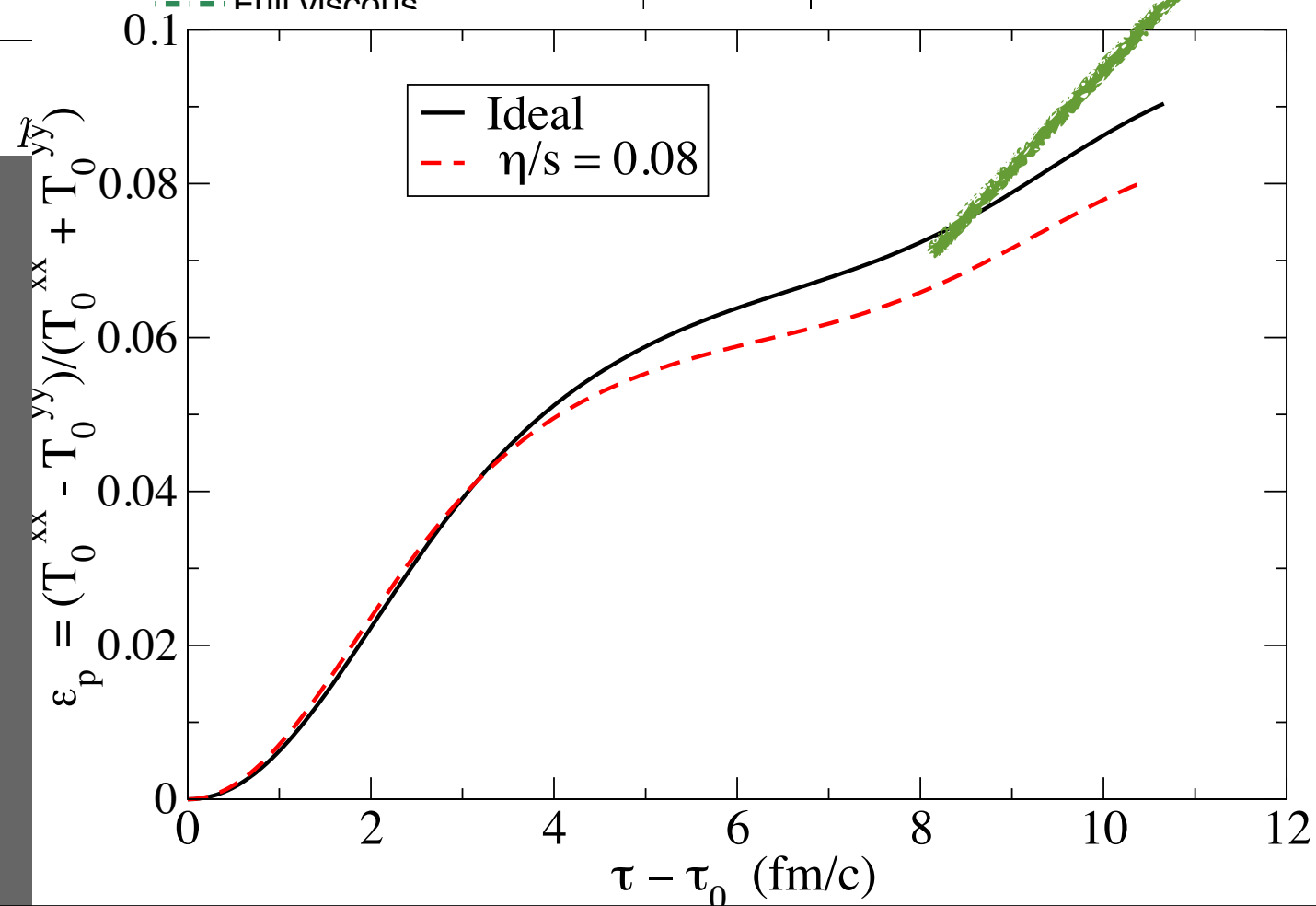
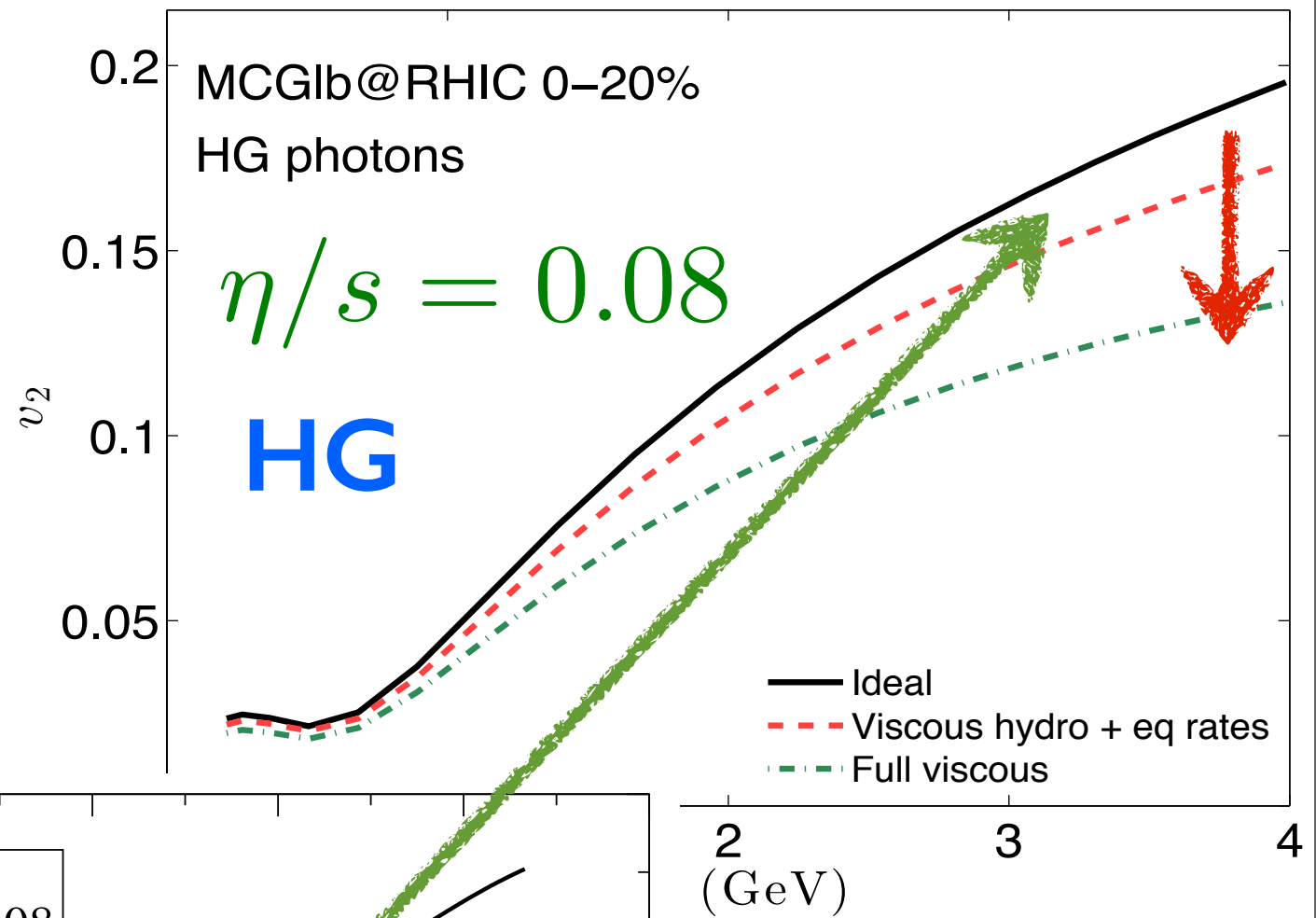
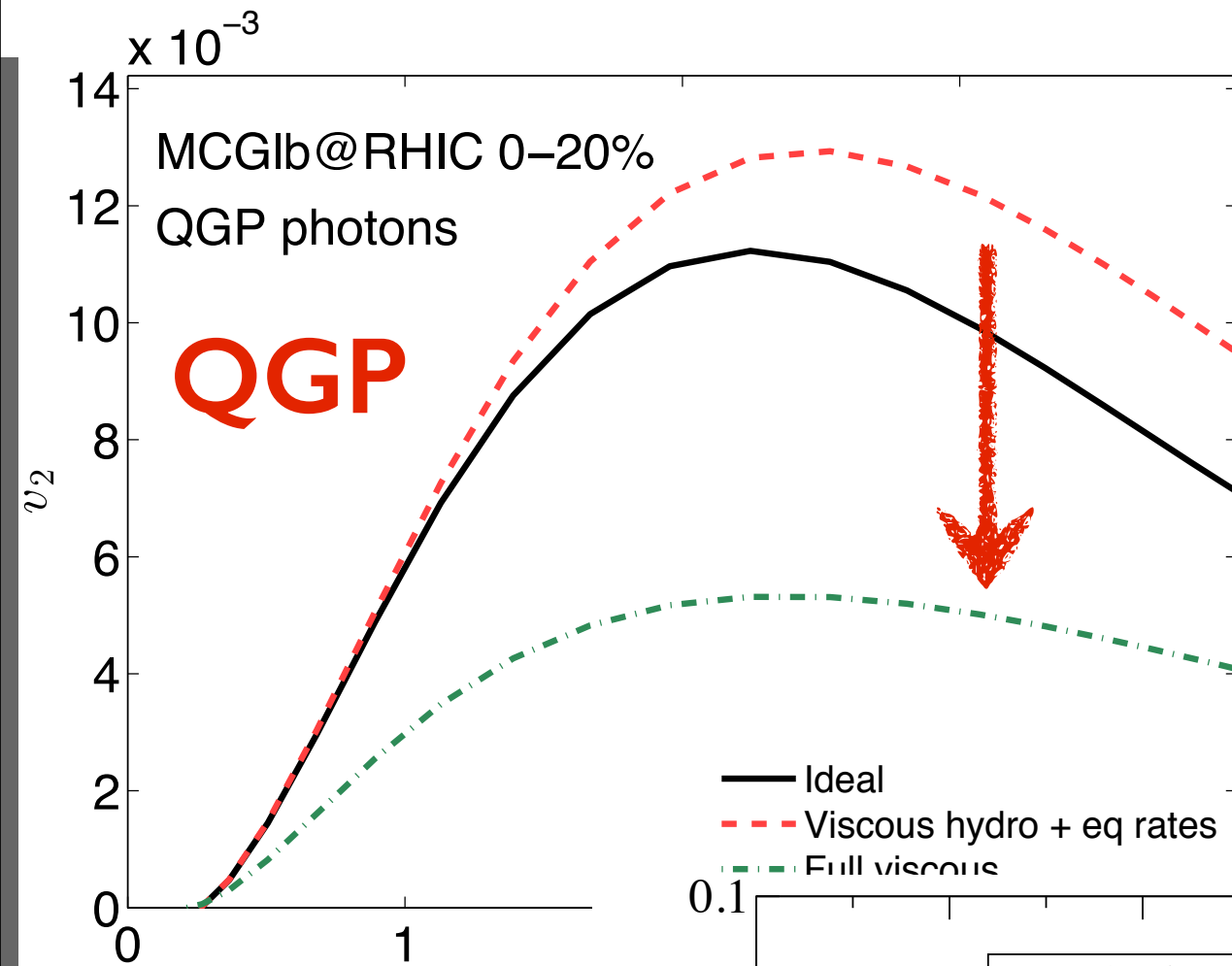
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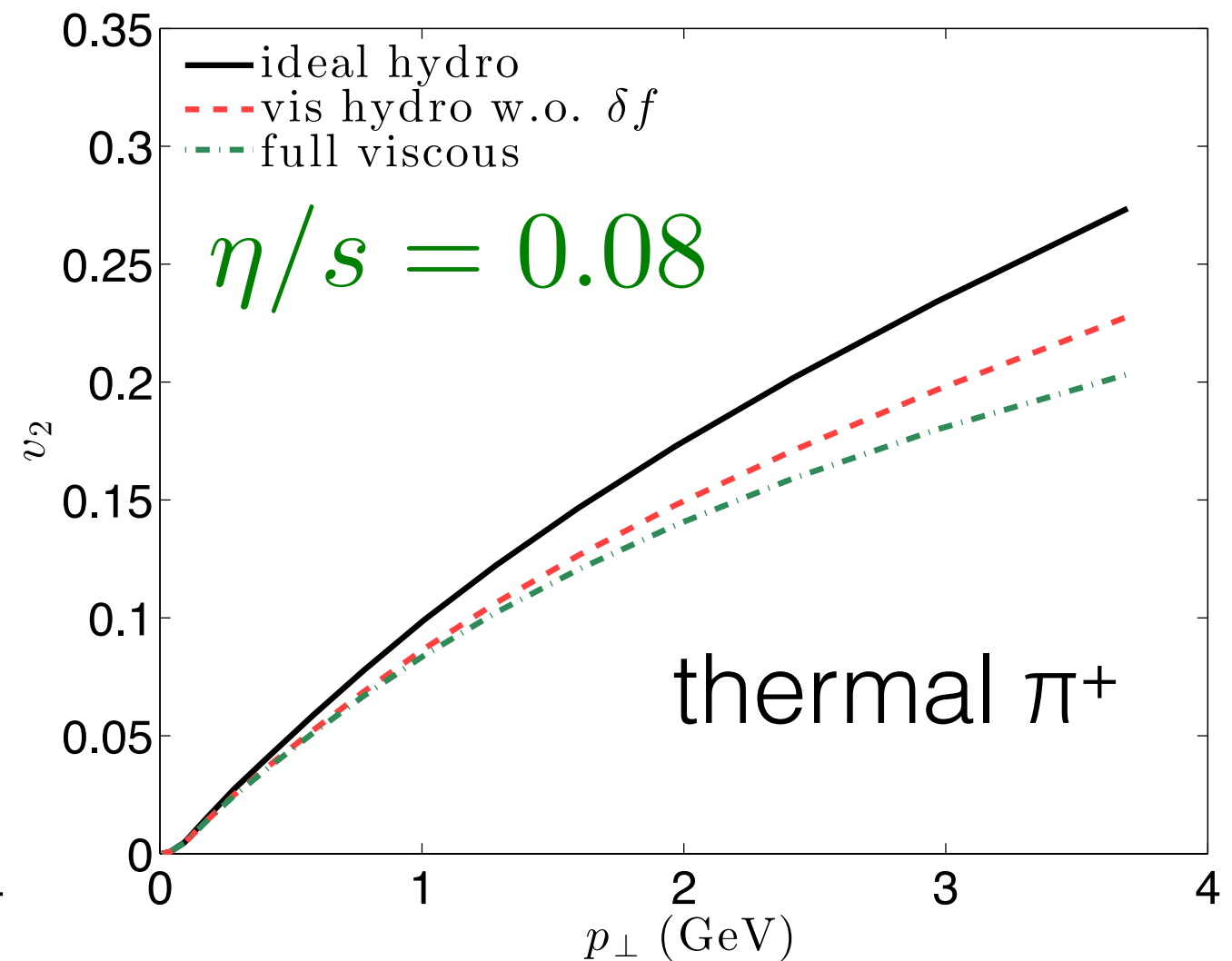
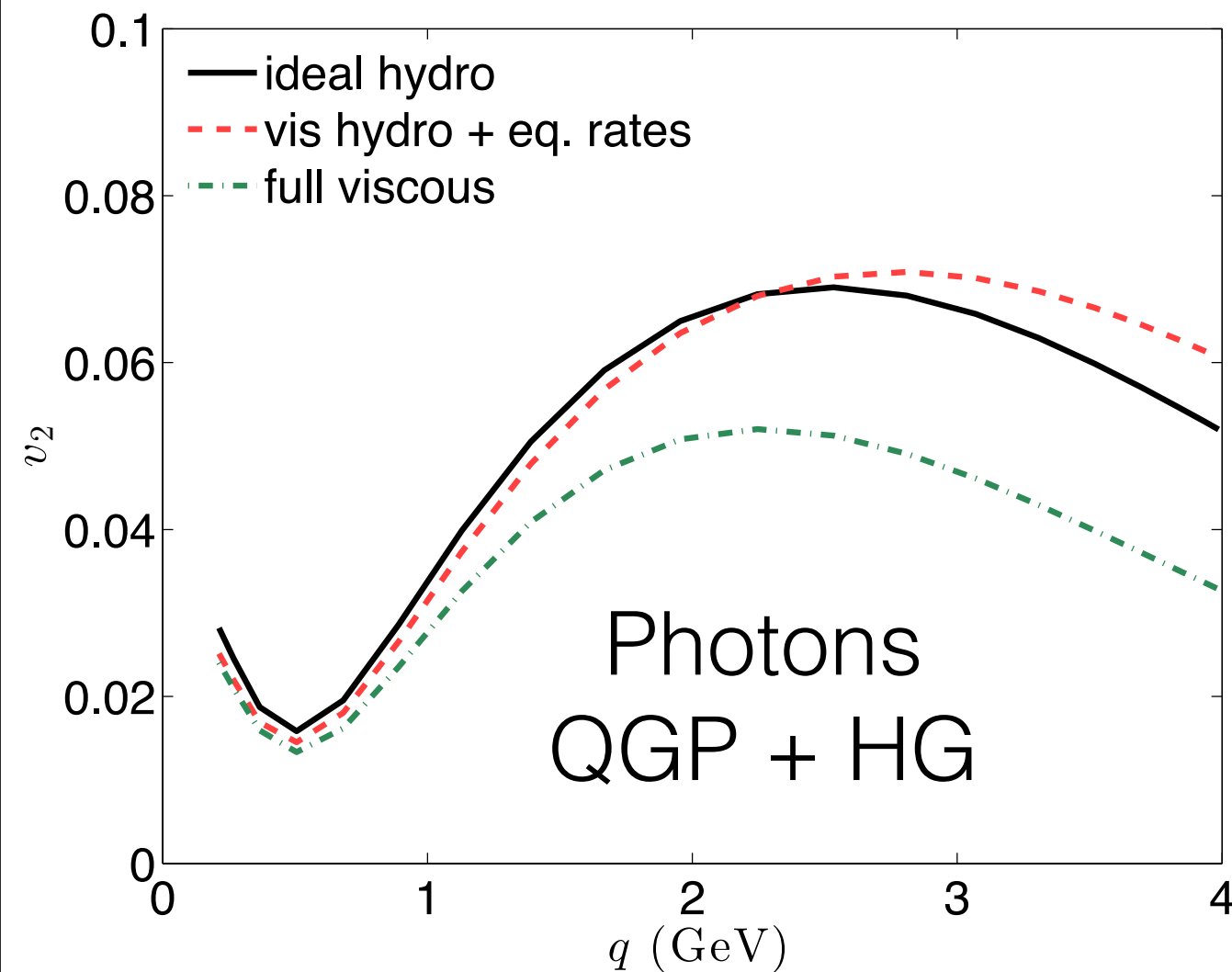
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Viscous effects on photon elliptic flow

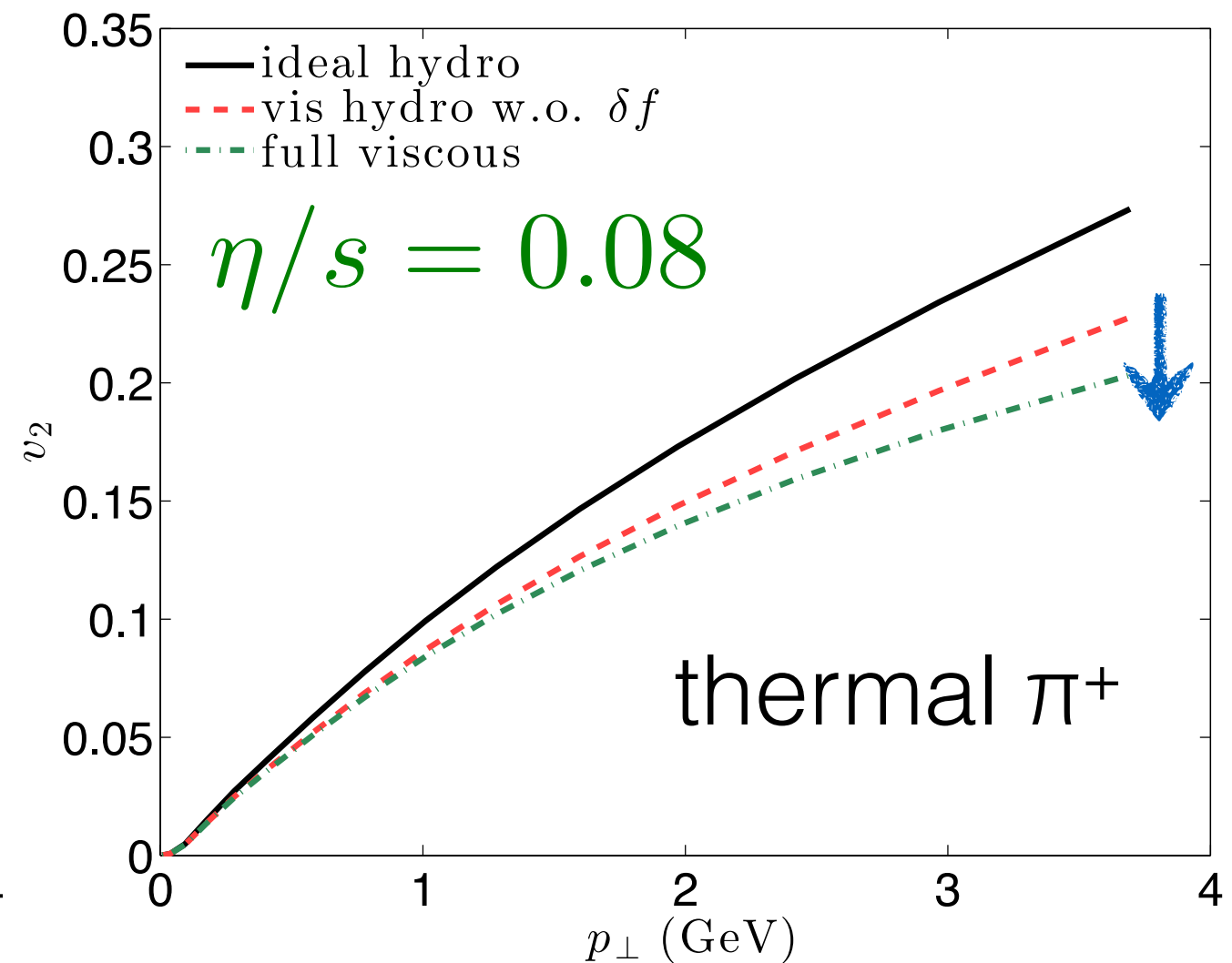
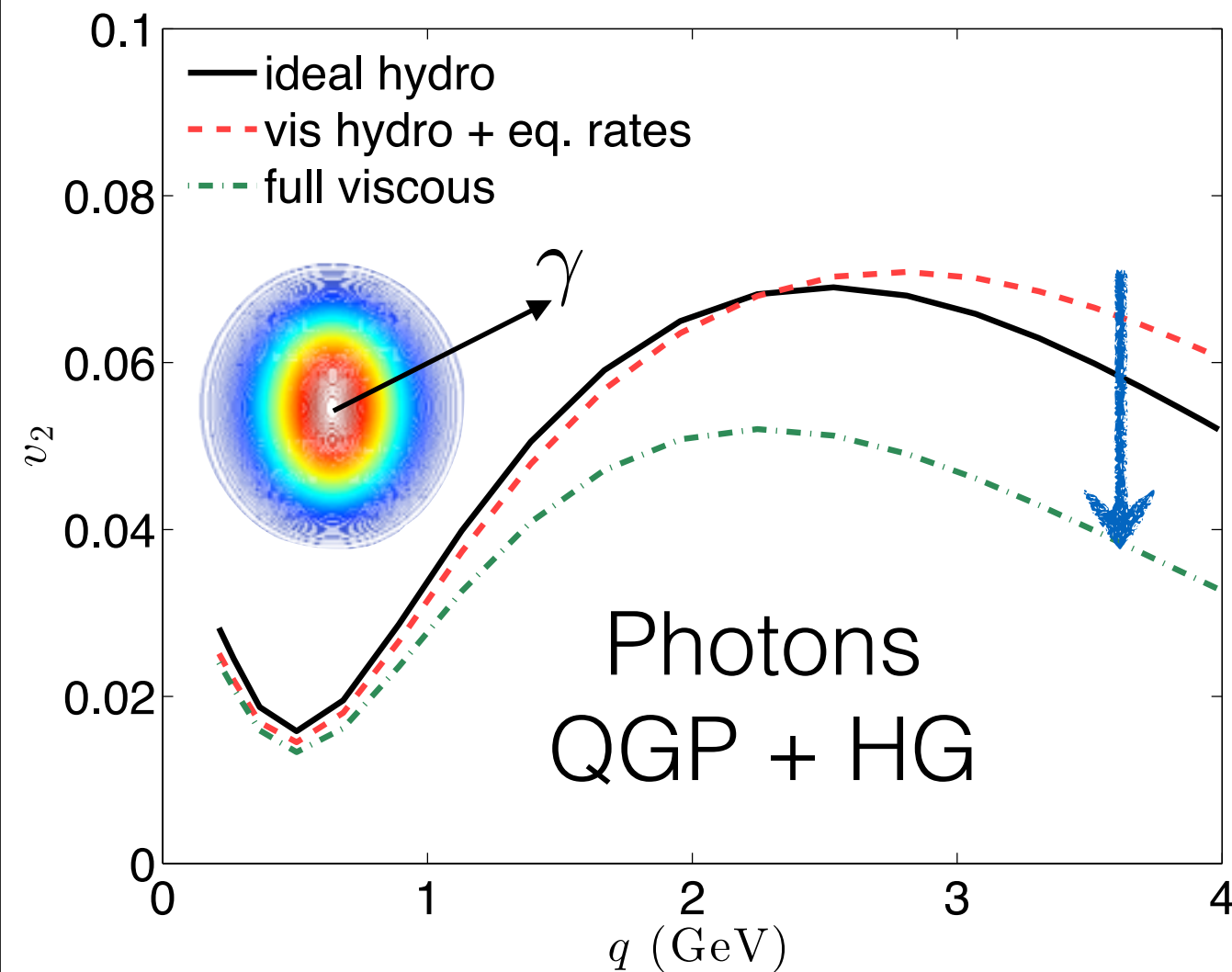


Viscous effects on photon elliptic flow



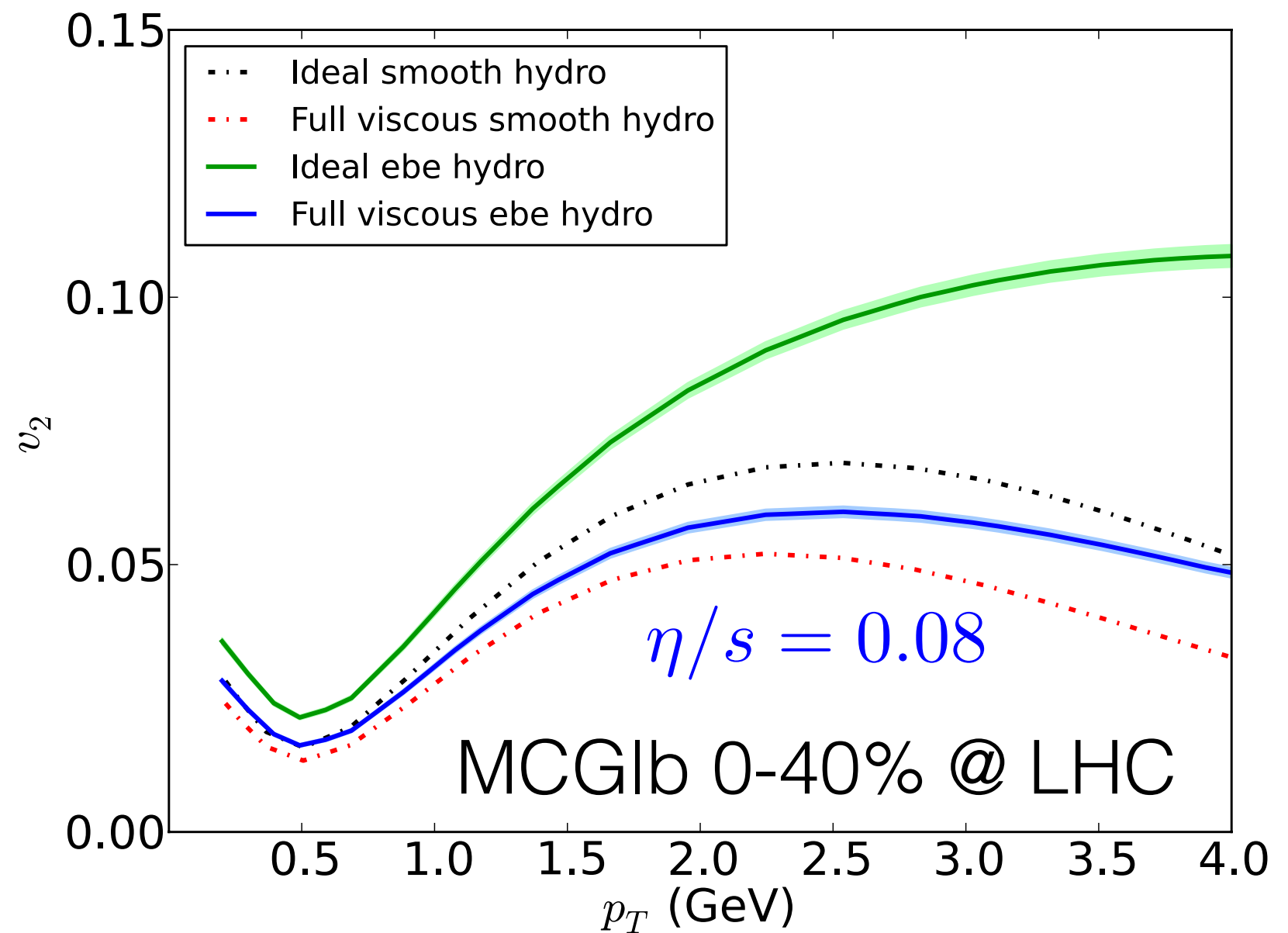
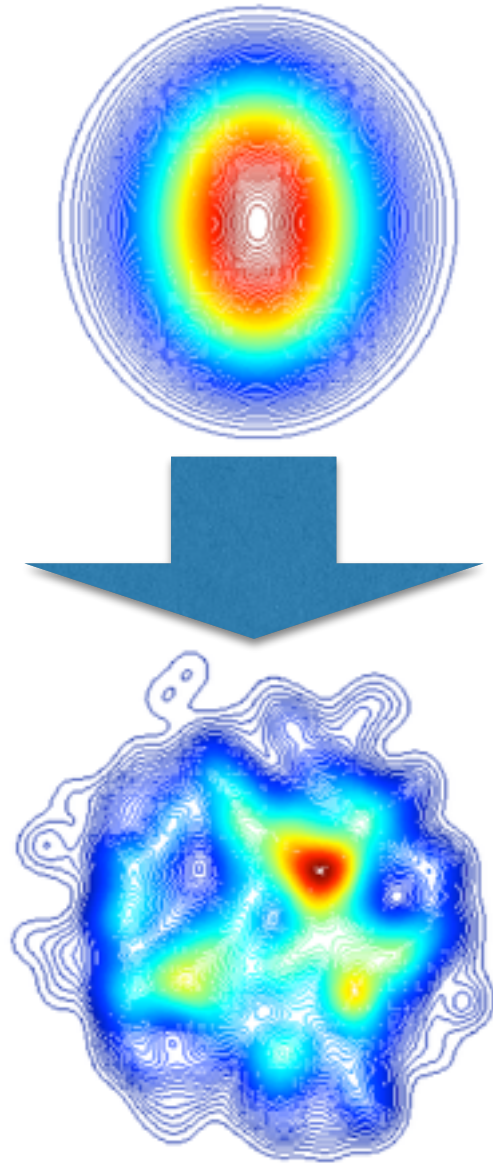
- Shear viscous suppression of photon v_2 is dominated by the viscous corrections to photon emission rates
- Photon elliptic flow is more sensitive to the evolution of shear stress tensor during the early time

Viscous effects on photon elliptic flow

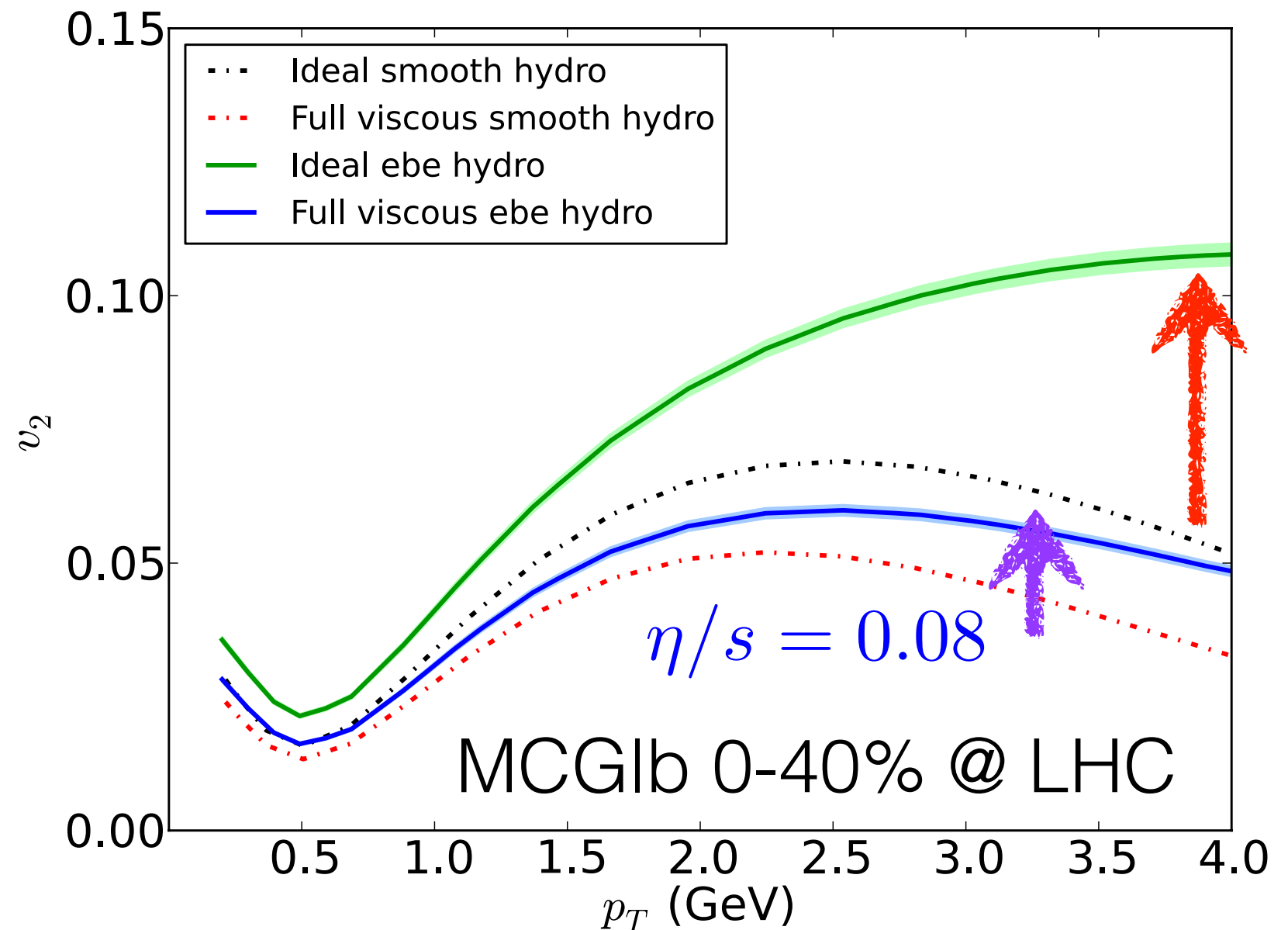
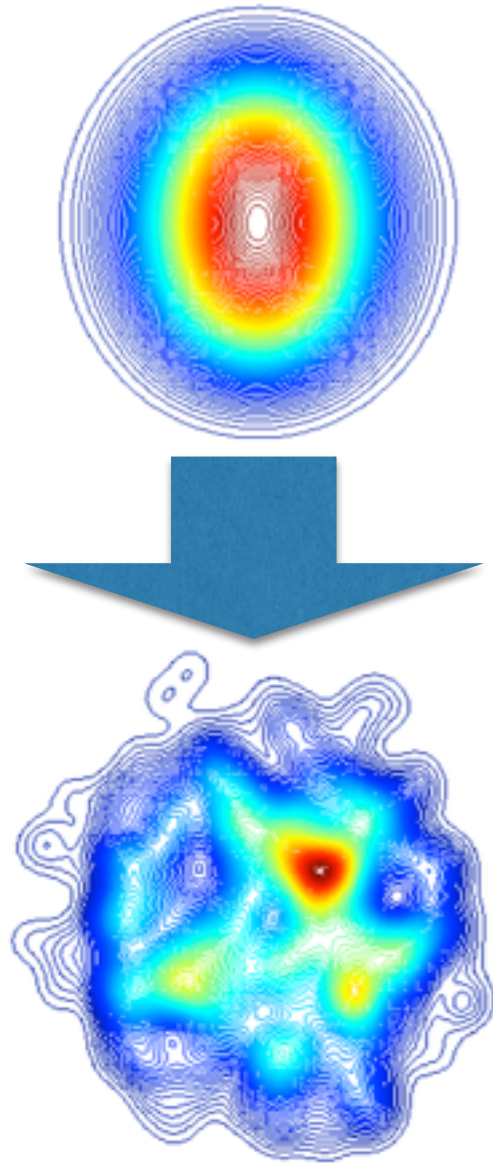


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Fluctuation effects on photon elliptic flow

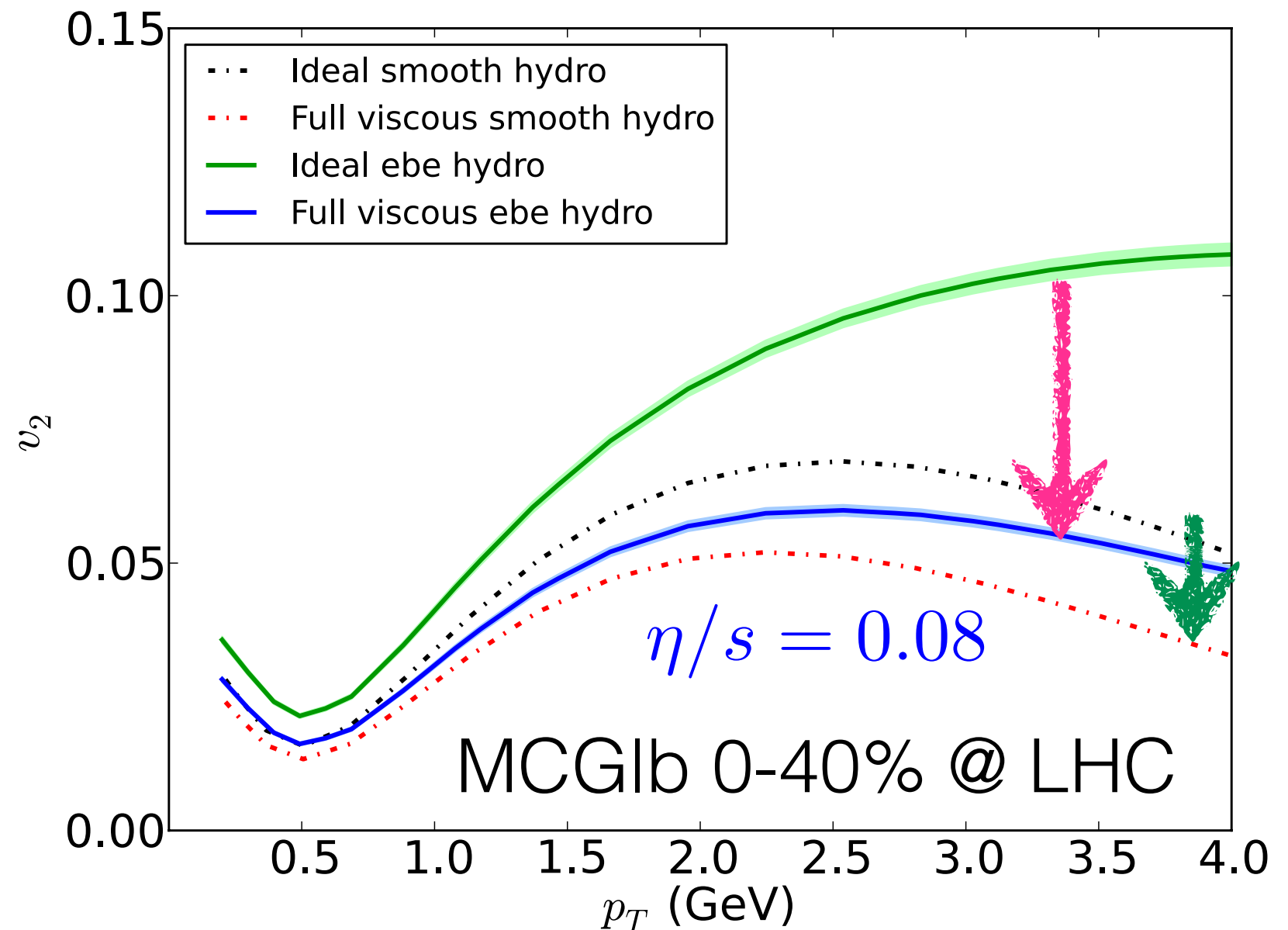
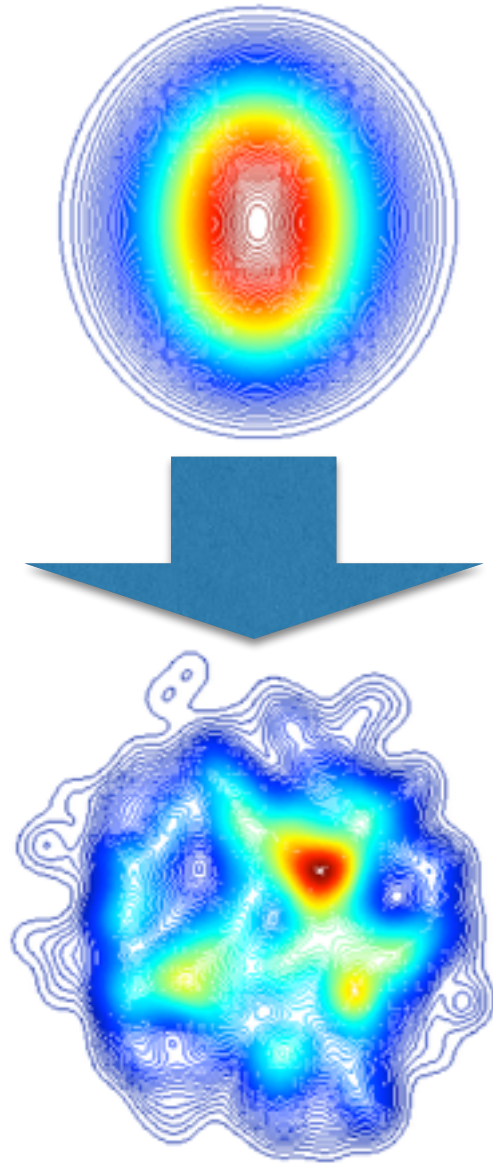


Fluctuation effects on photon elliptic flow



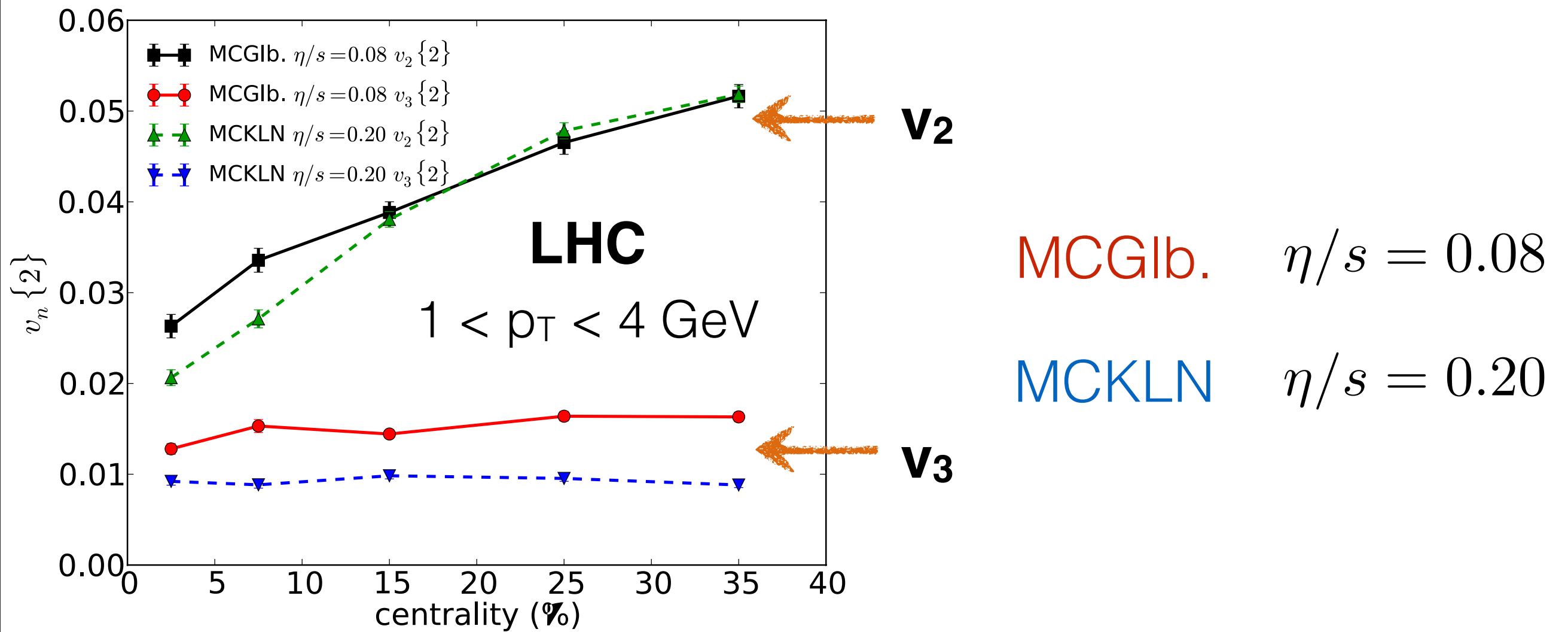
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Fluctuation effects on photon elliptic flow



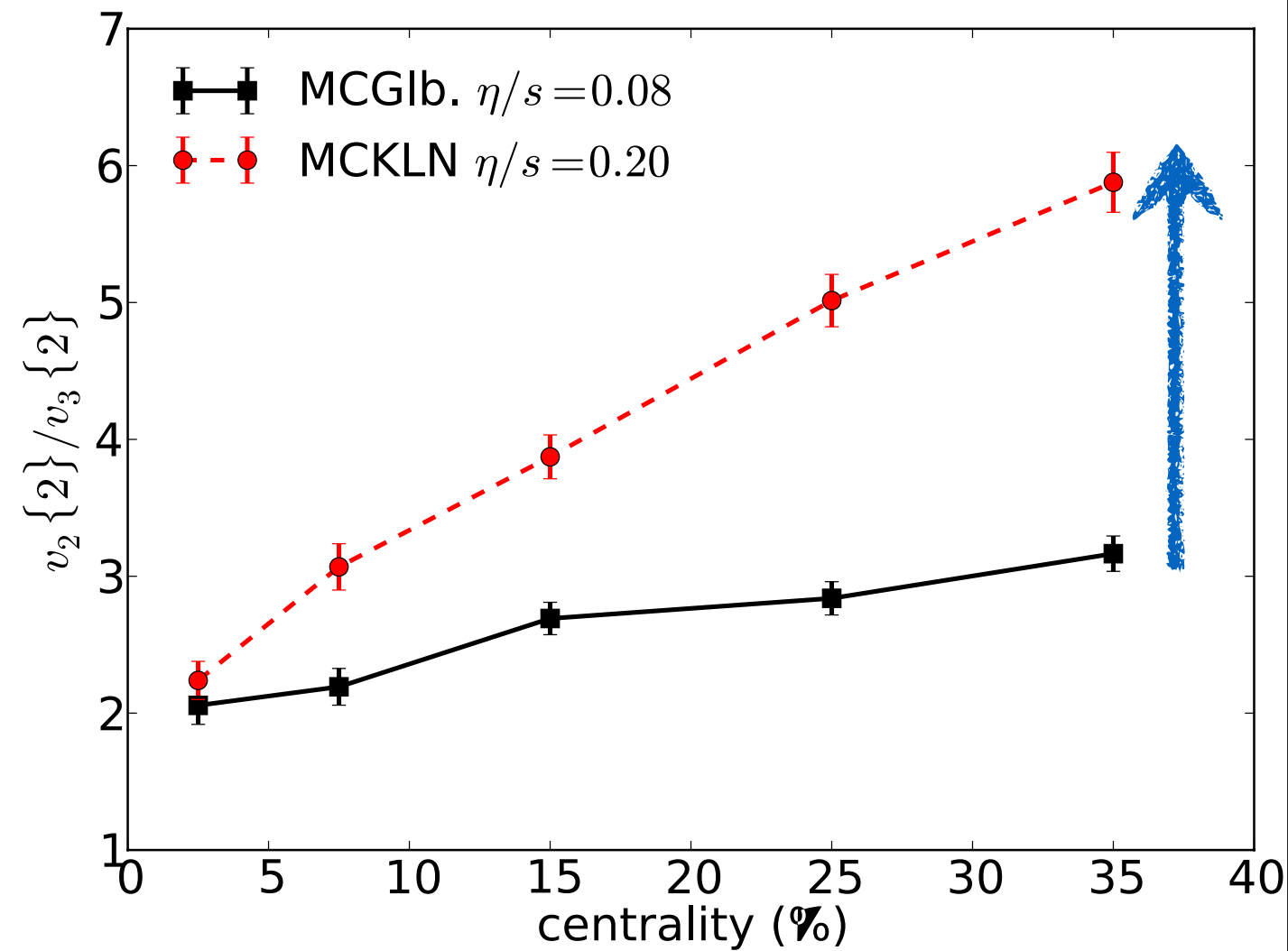
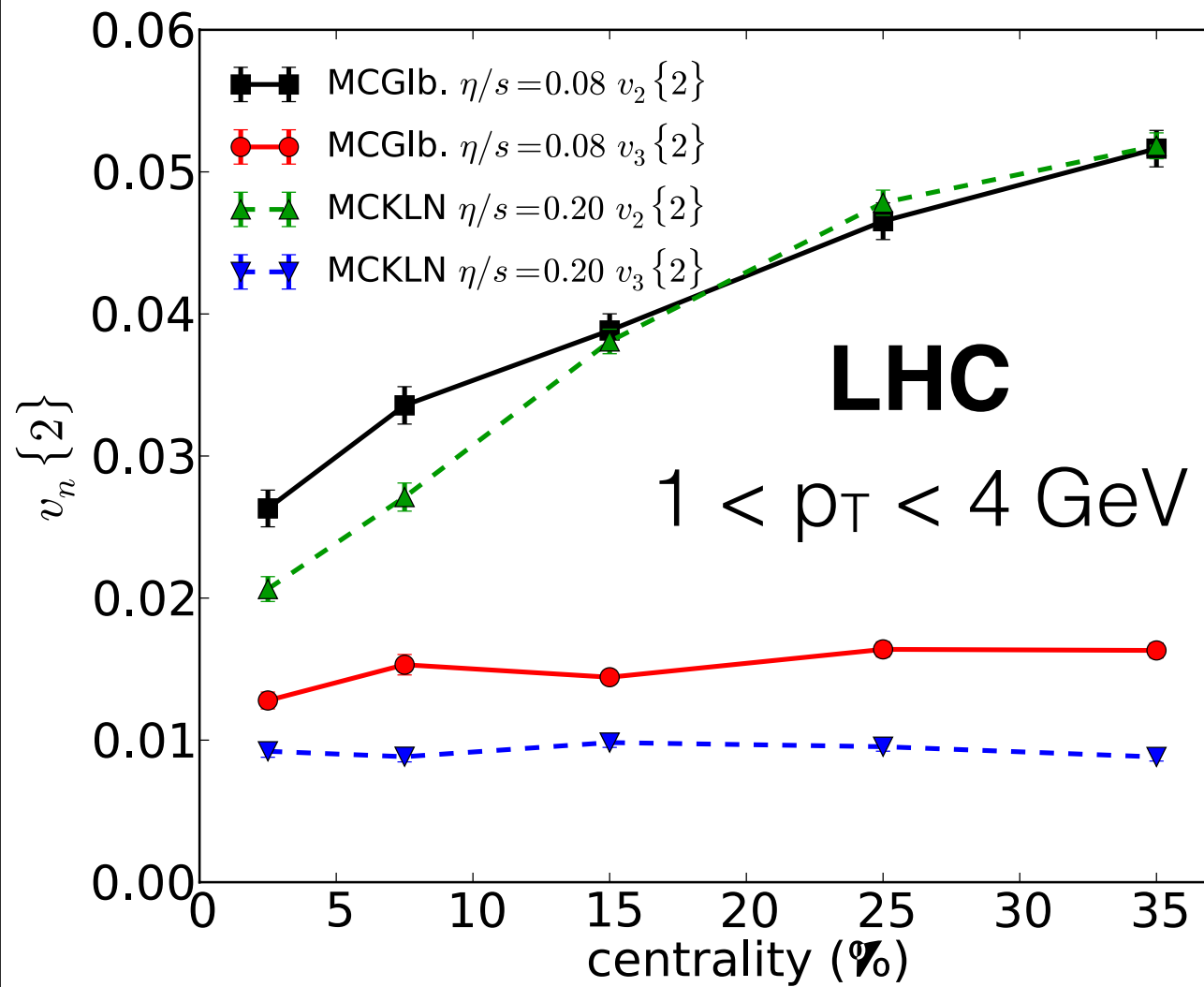
- ▶ Initial fluctuations increase photon's elliptic flow
- ▶ Viscous suppression is larger in the event-by-event runs

Event-by-Event Full Viscous Photon v_n



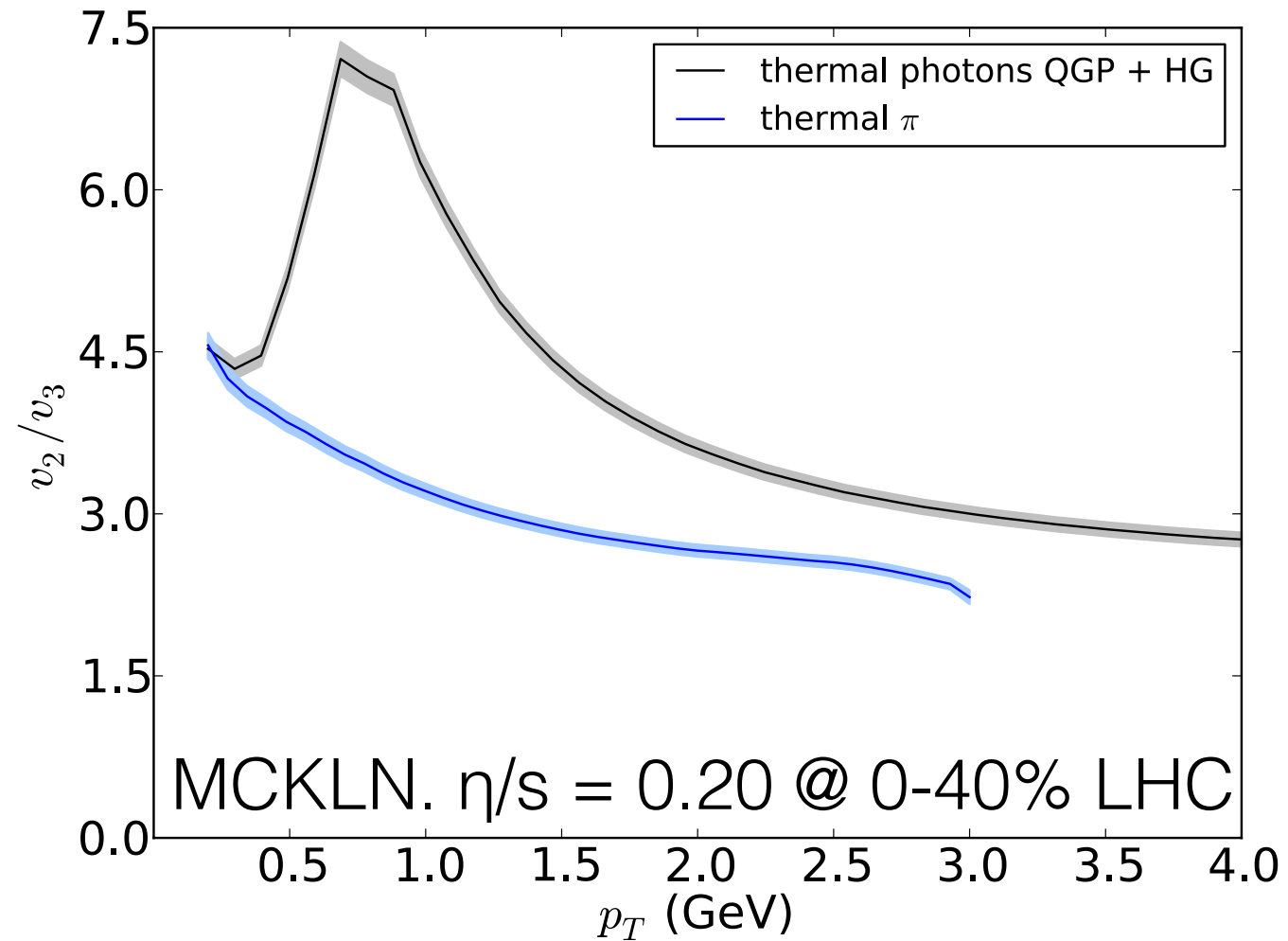
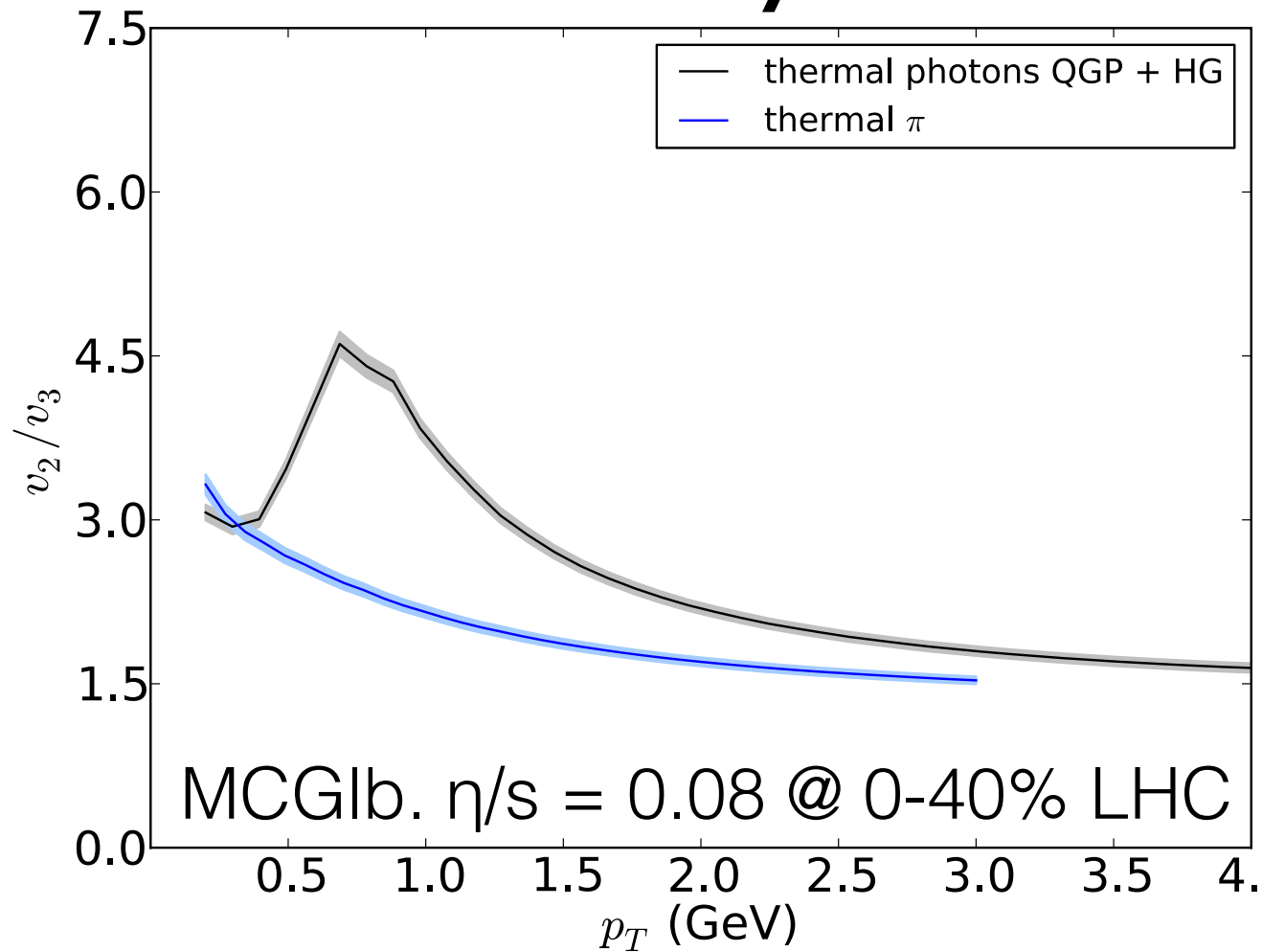
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Event-by-Event Full Viscous Photon v_n



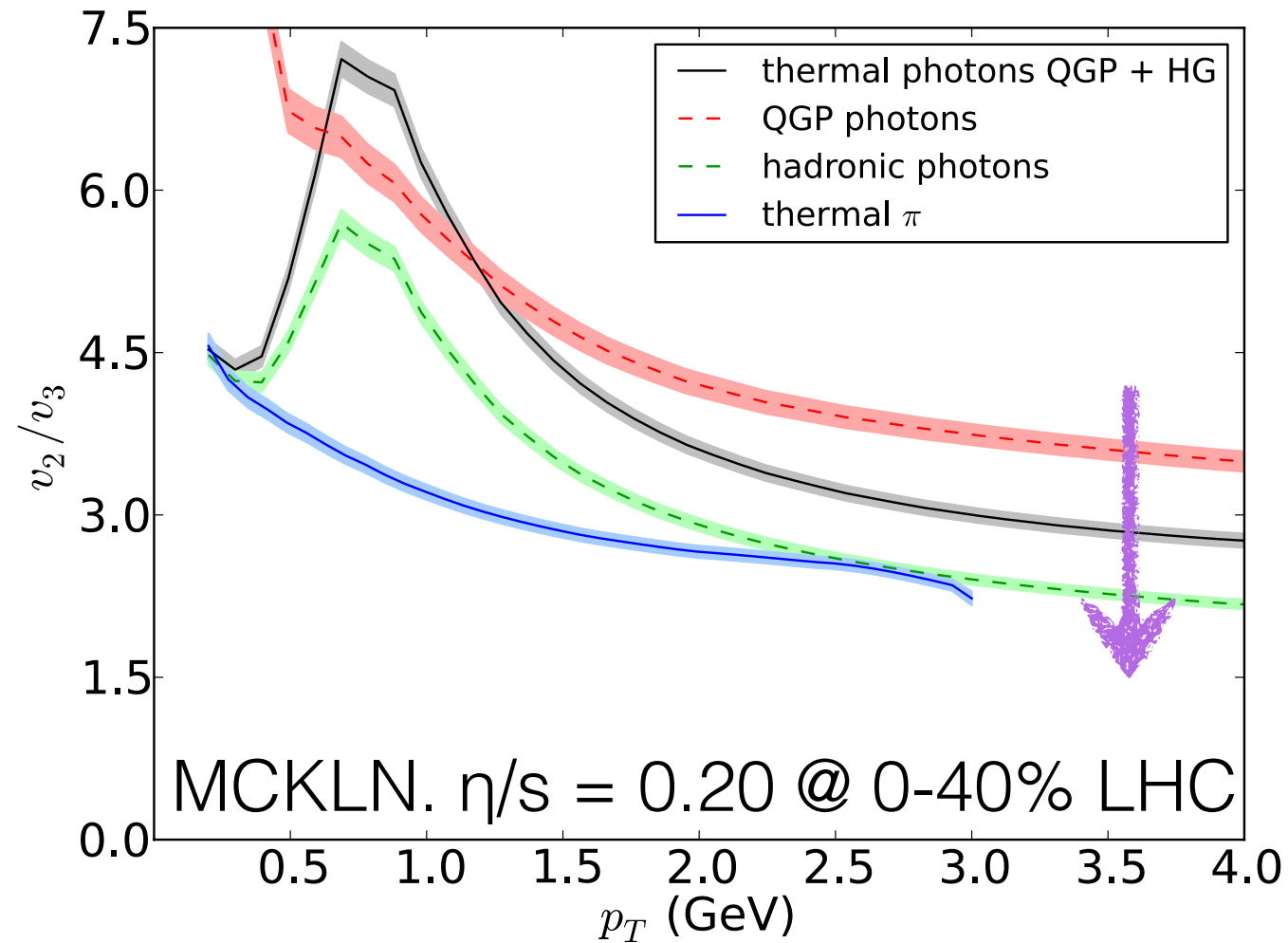
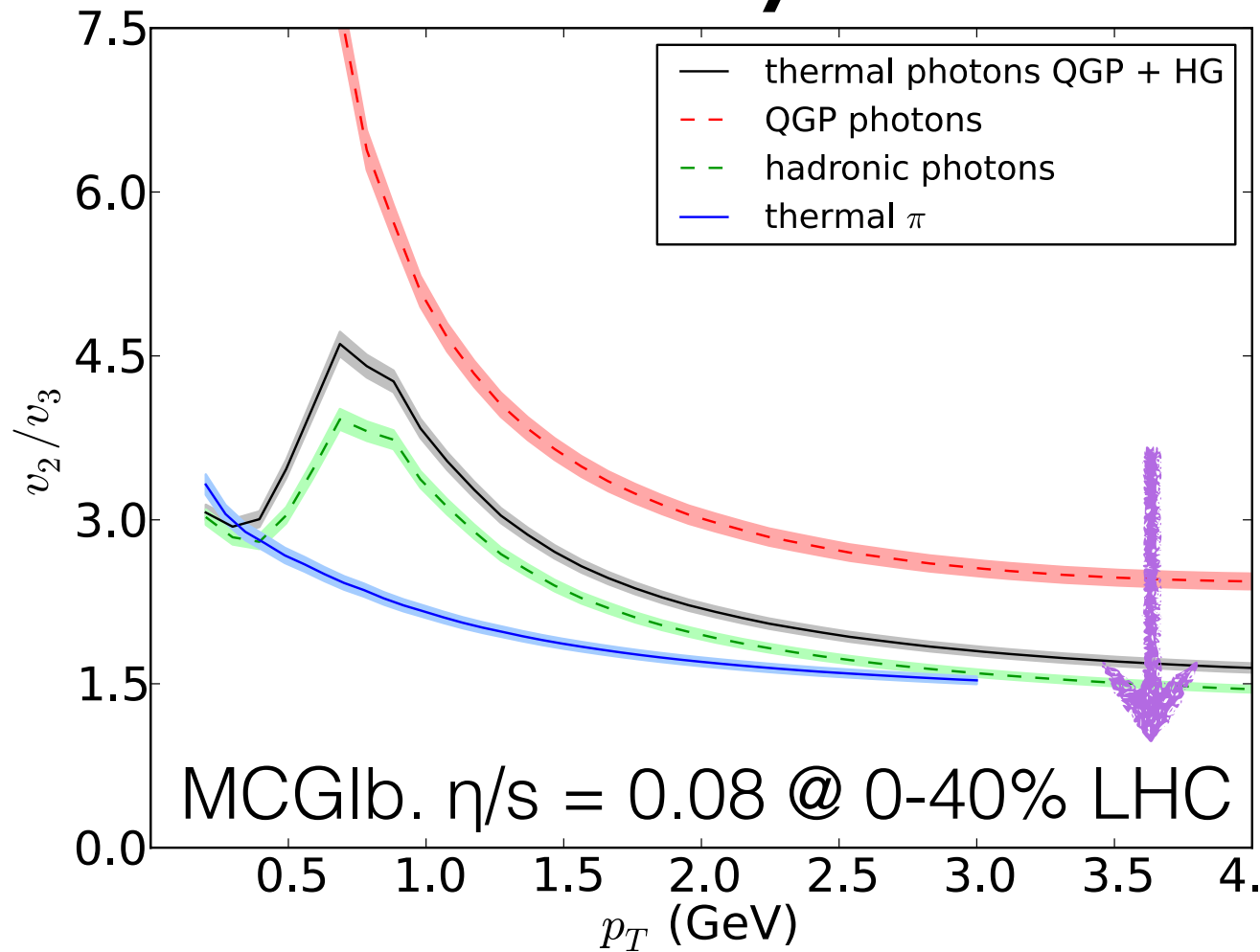
- The anisotropic flows of photons show similar centrality behavior as hadrons v_n
- The ratio of v_2/v_3 increase with the shear viscosity.
- The centrality dependence of this ratio is stronger for MCKLN model

Event-by-Event Full Viscous Photon v_n



- The ratio of v_2/v_3 of photons is larger than the ratio of thermal pions

Event-by-Event Full Viscous Photon v_n



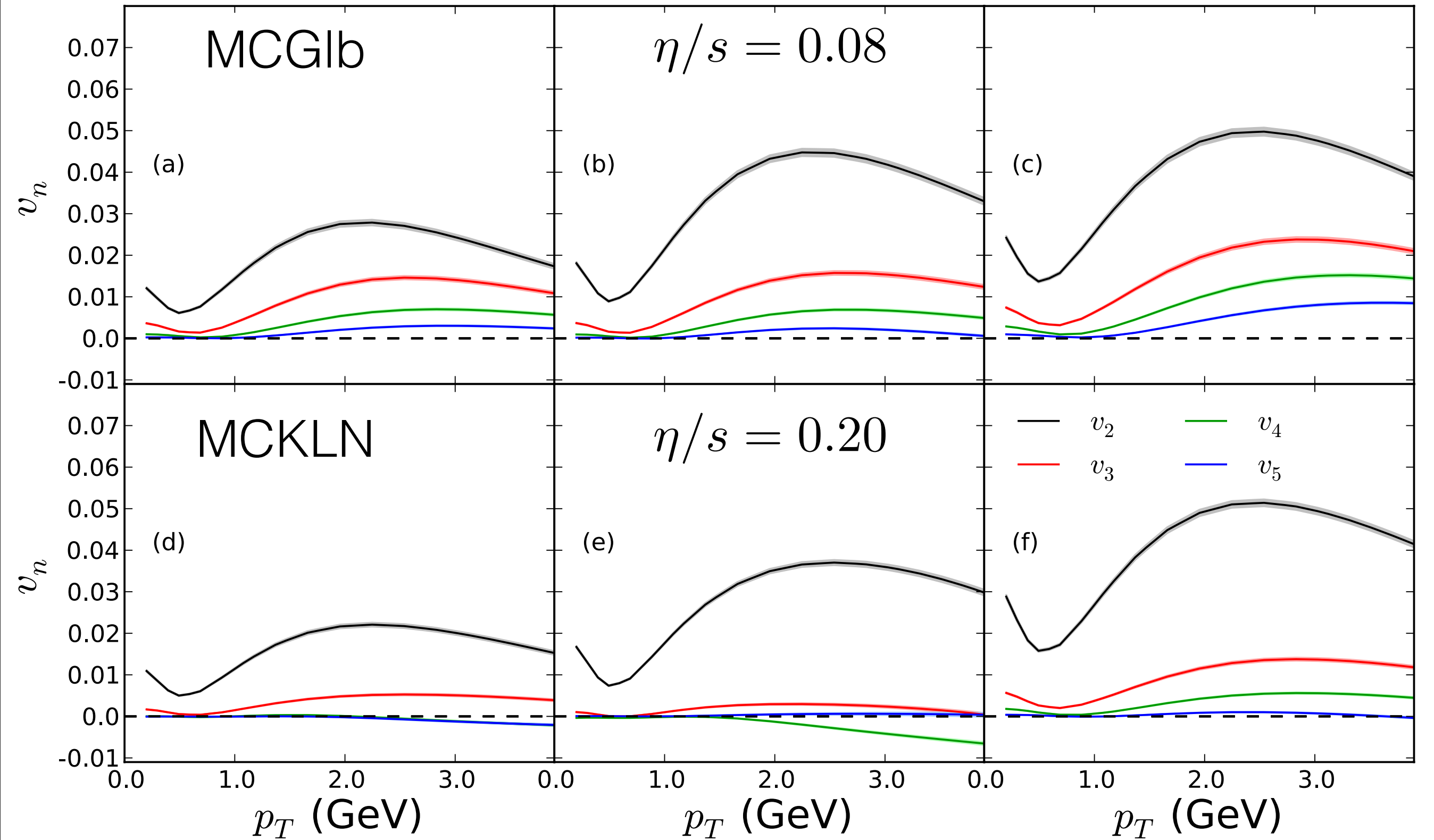
- The ratio of v_2/v_3 of photons is larger than the ratio of thermal pions
- The ratio of v_2/v_3 is larger for QGP photons compared to hadronic photons which indicates triangular flow develops faster than elliptic flow during the late stage of hydrodynamic evolution

Event-by-Event Full Viscous Photon v_n

0-20% @ RHIC

20-40% @ RHIC

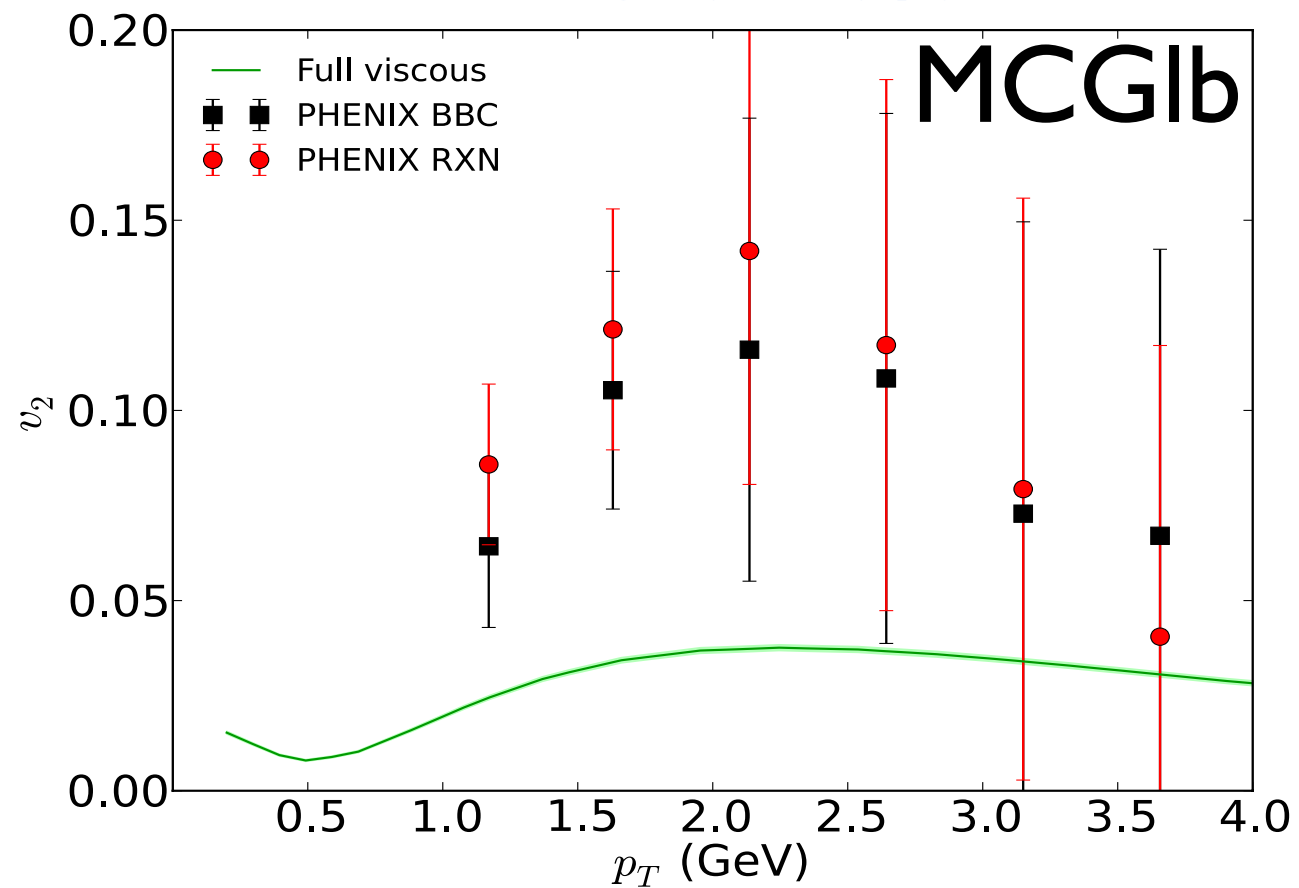
0-40% @ LHC



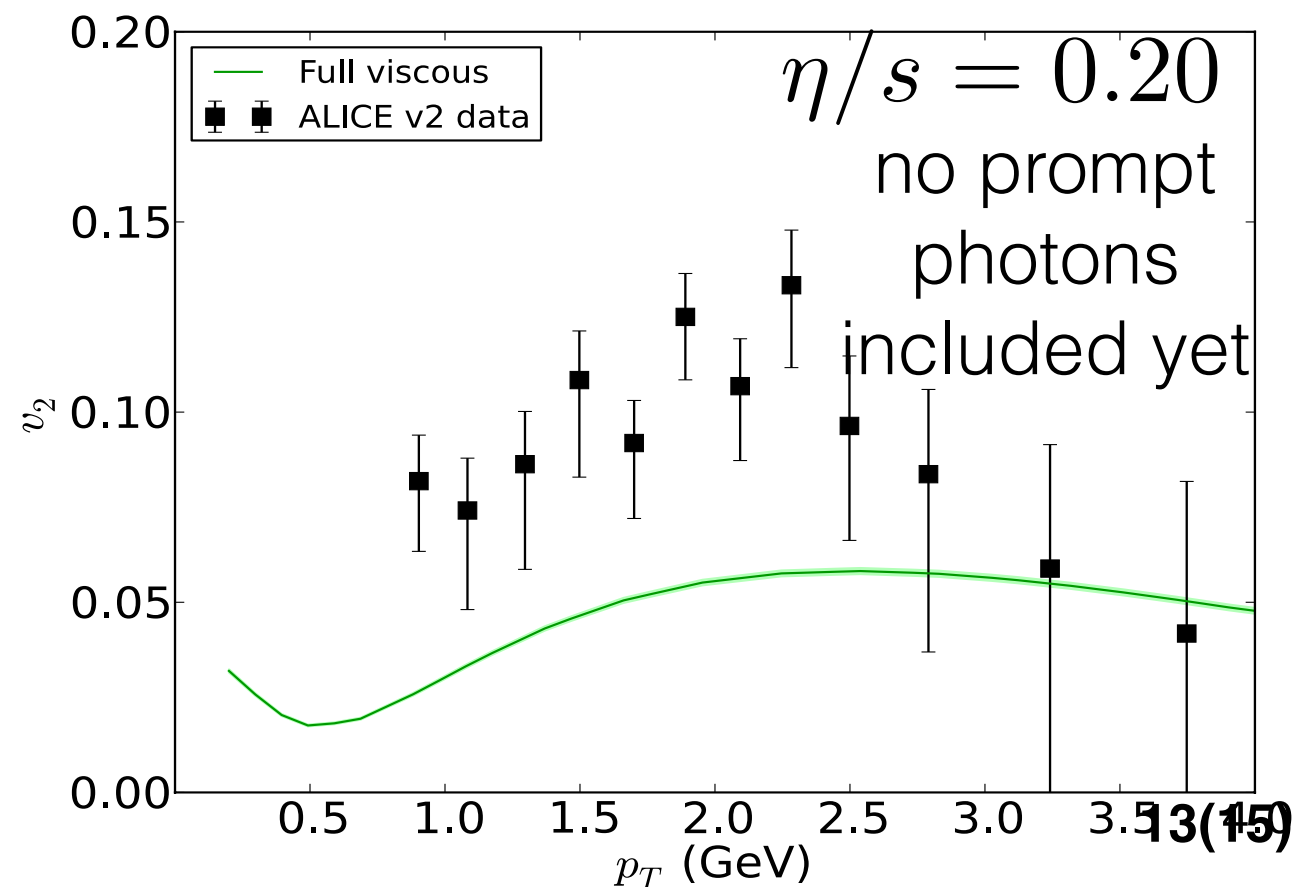
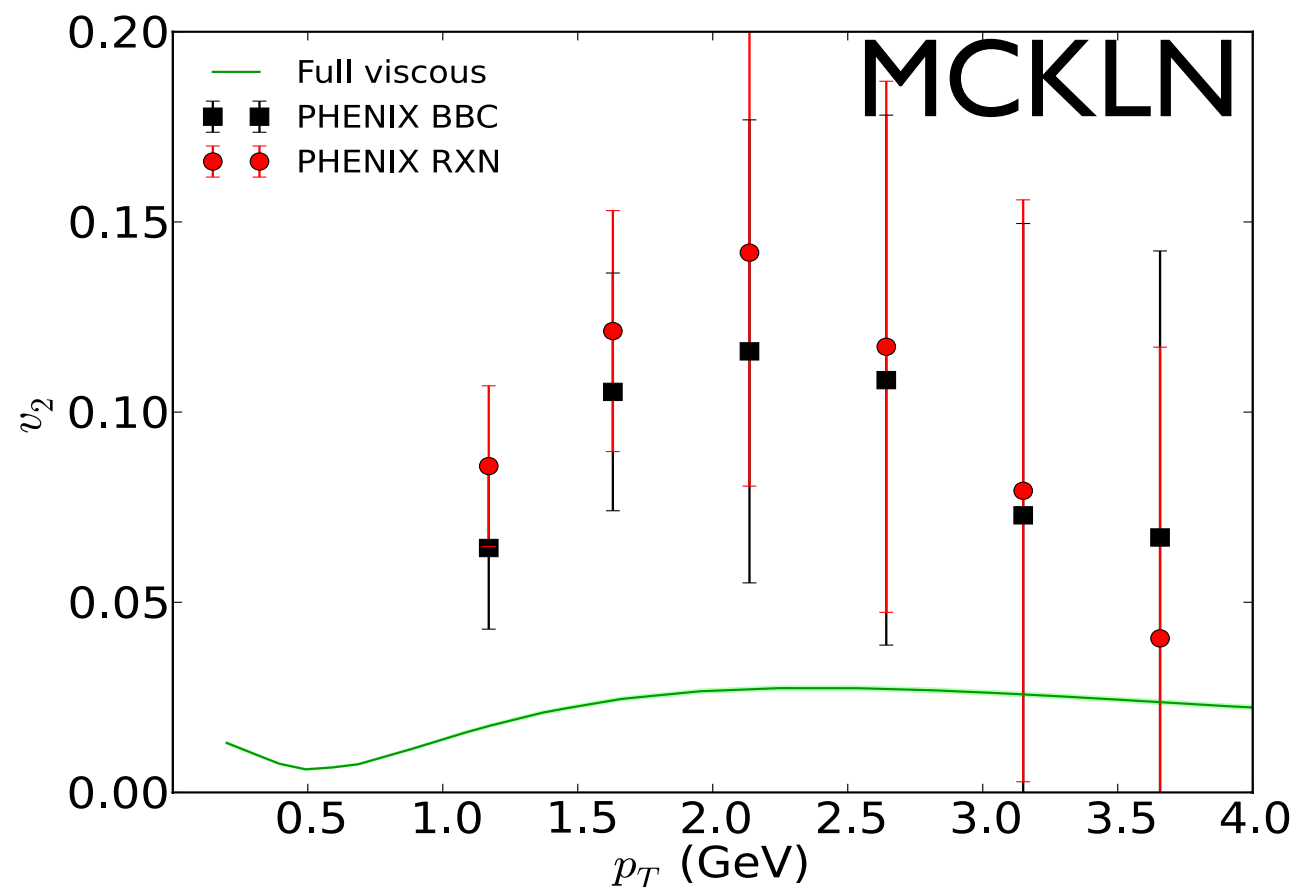
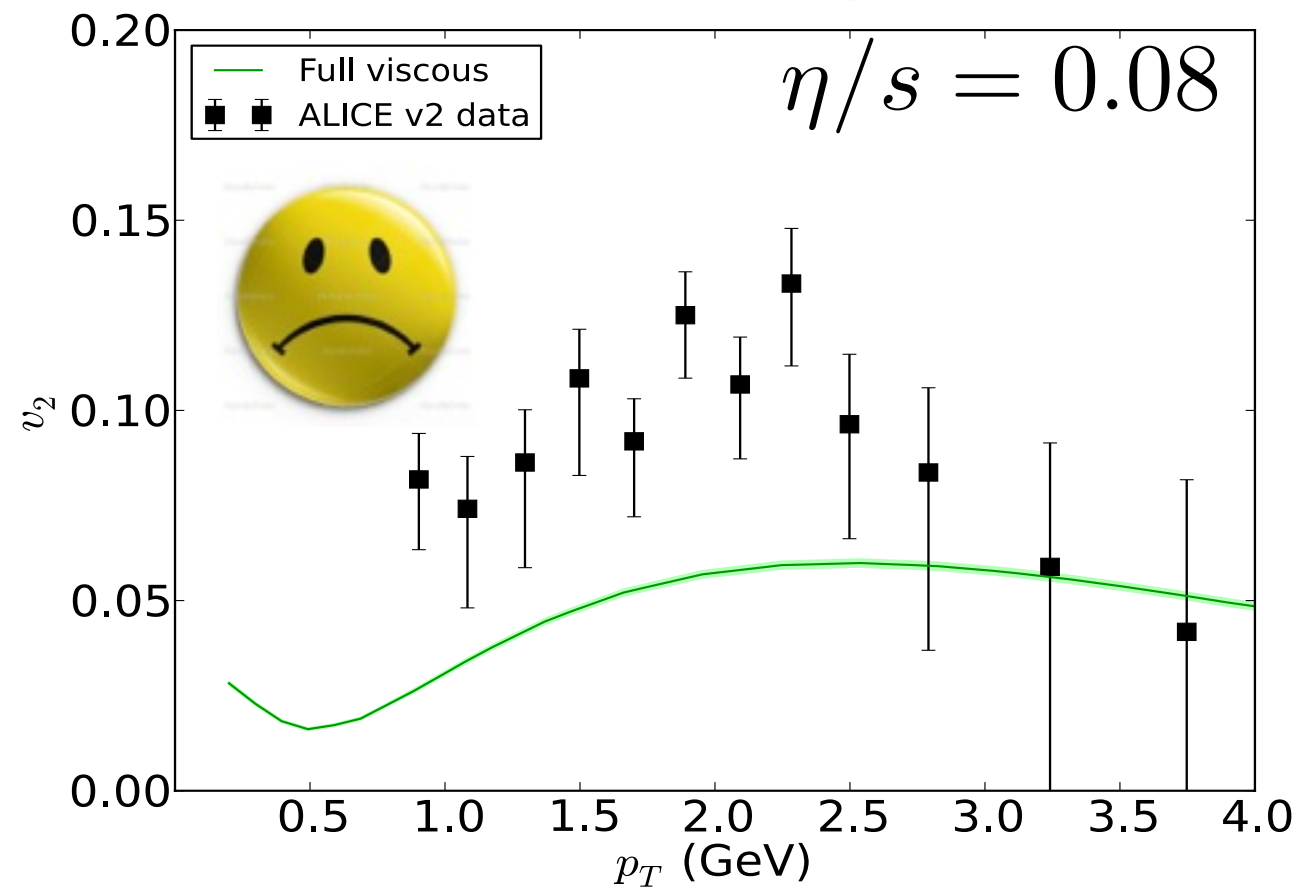
1000 events each centrality ¹²⁽¹⁵⁾

Comparisons with exp. data

RHIC 0-20%

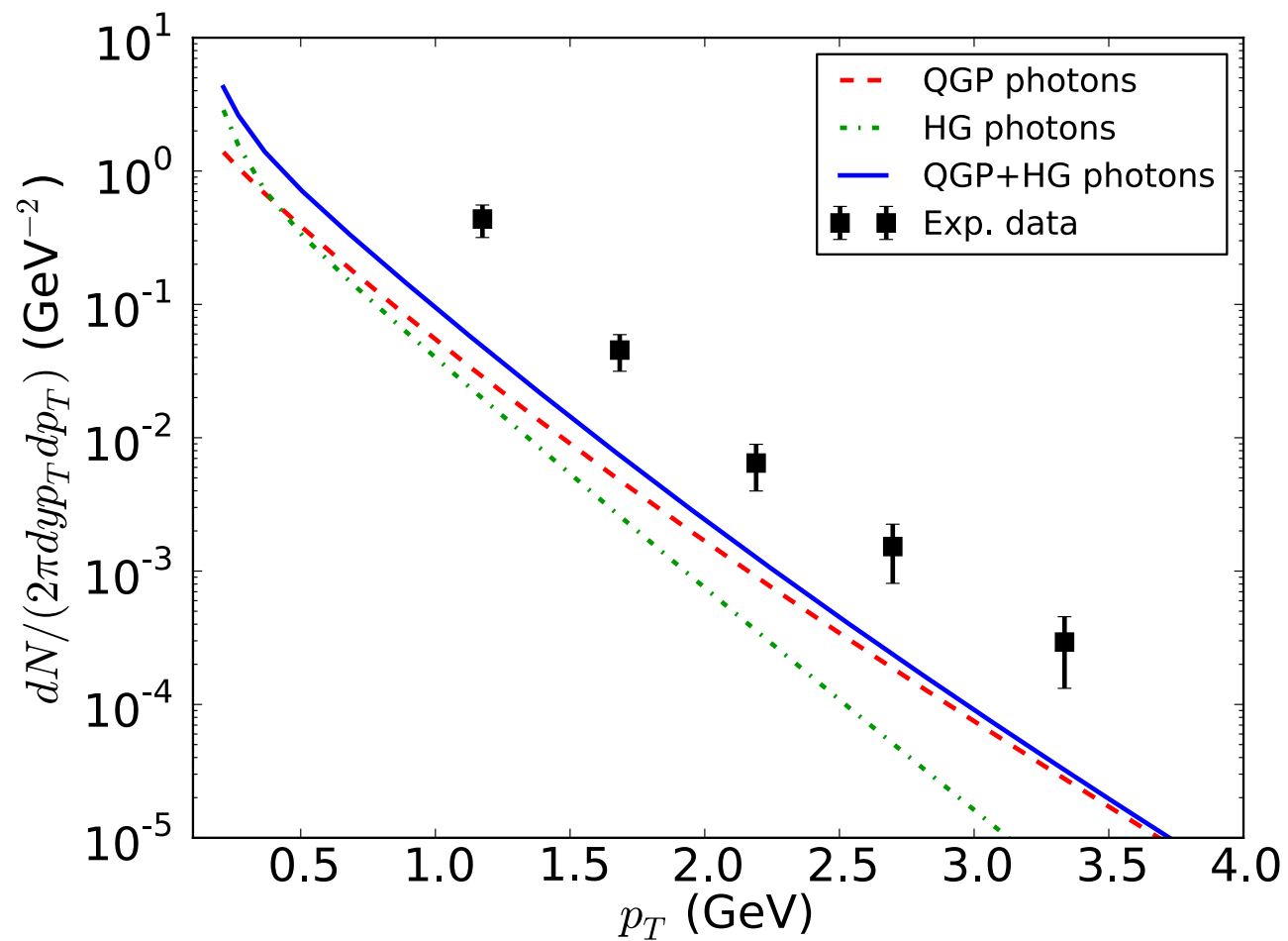
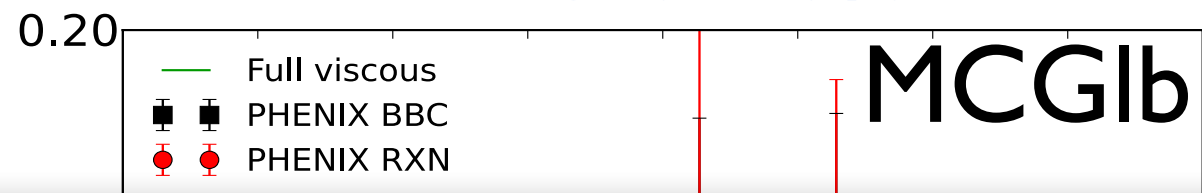


LHC 0-40%

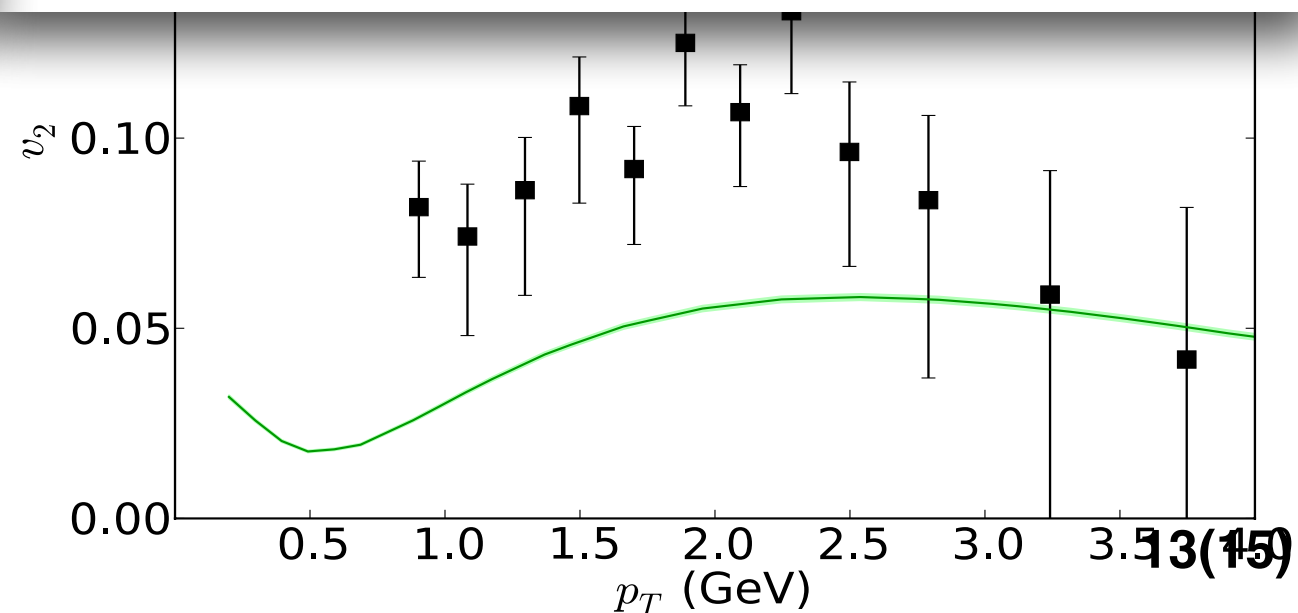
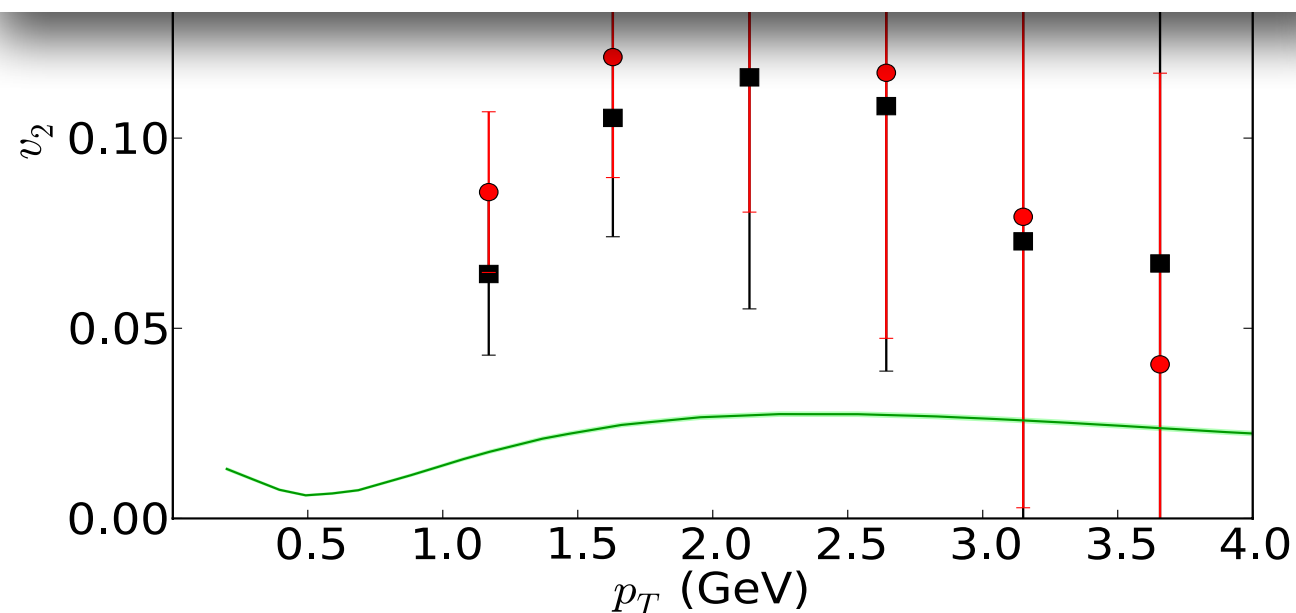
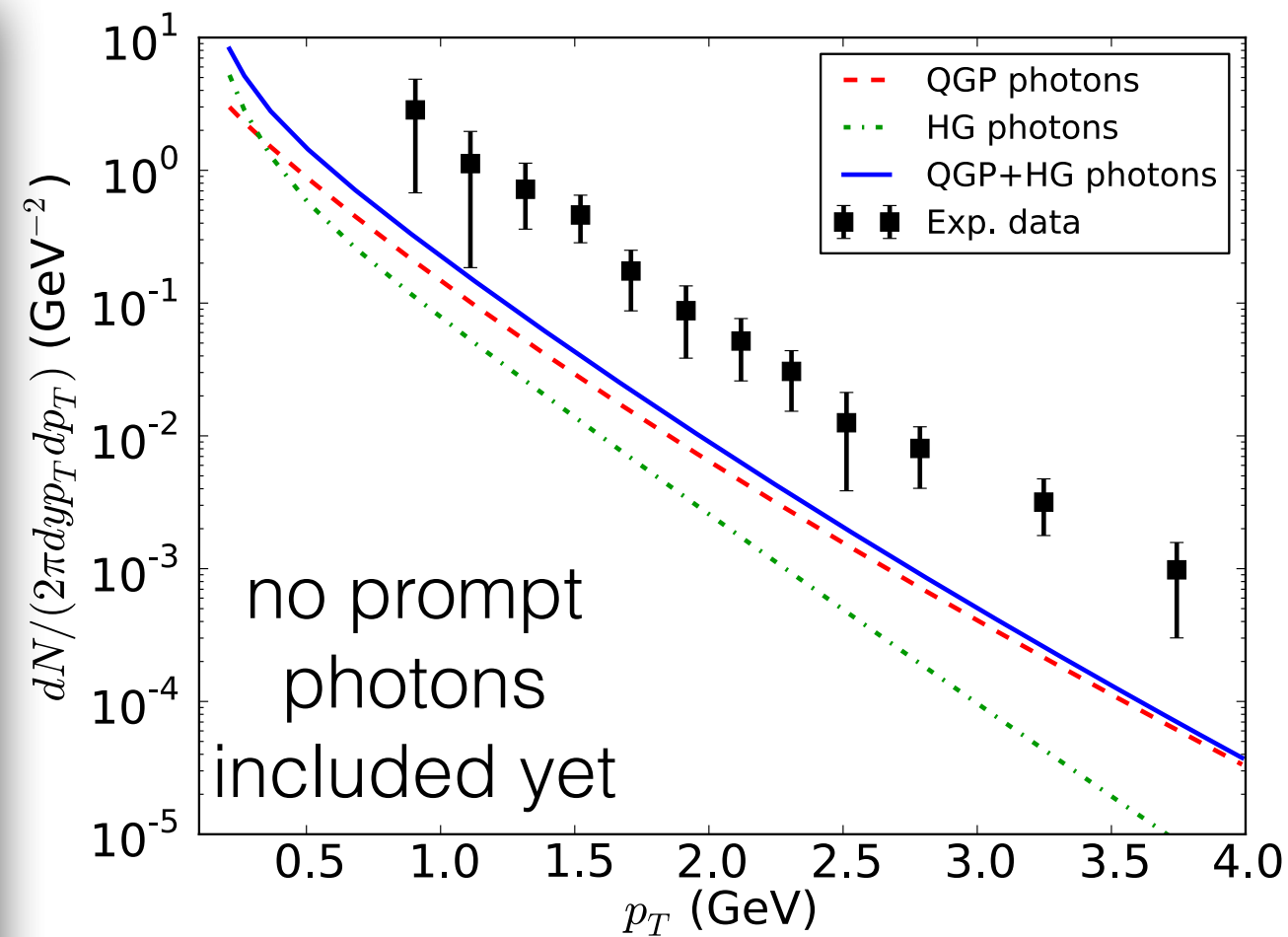
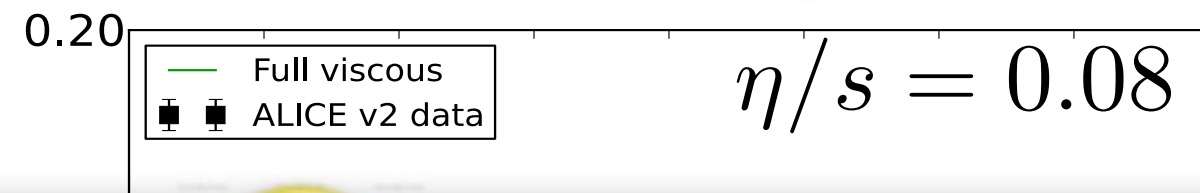


Comparisons with exp. data

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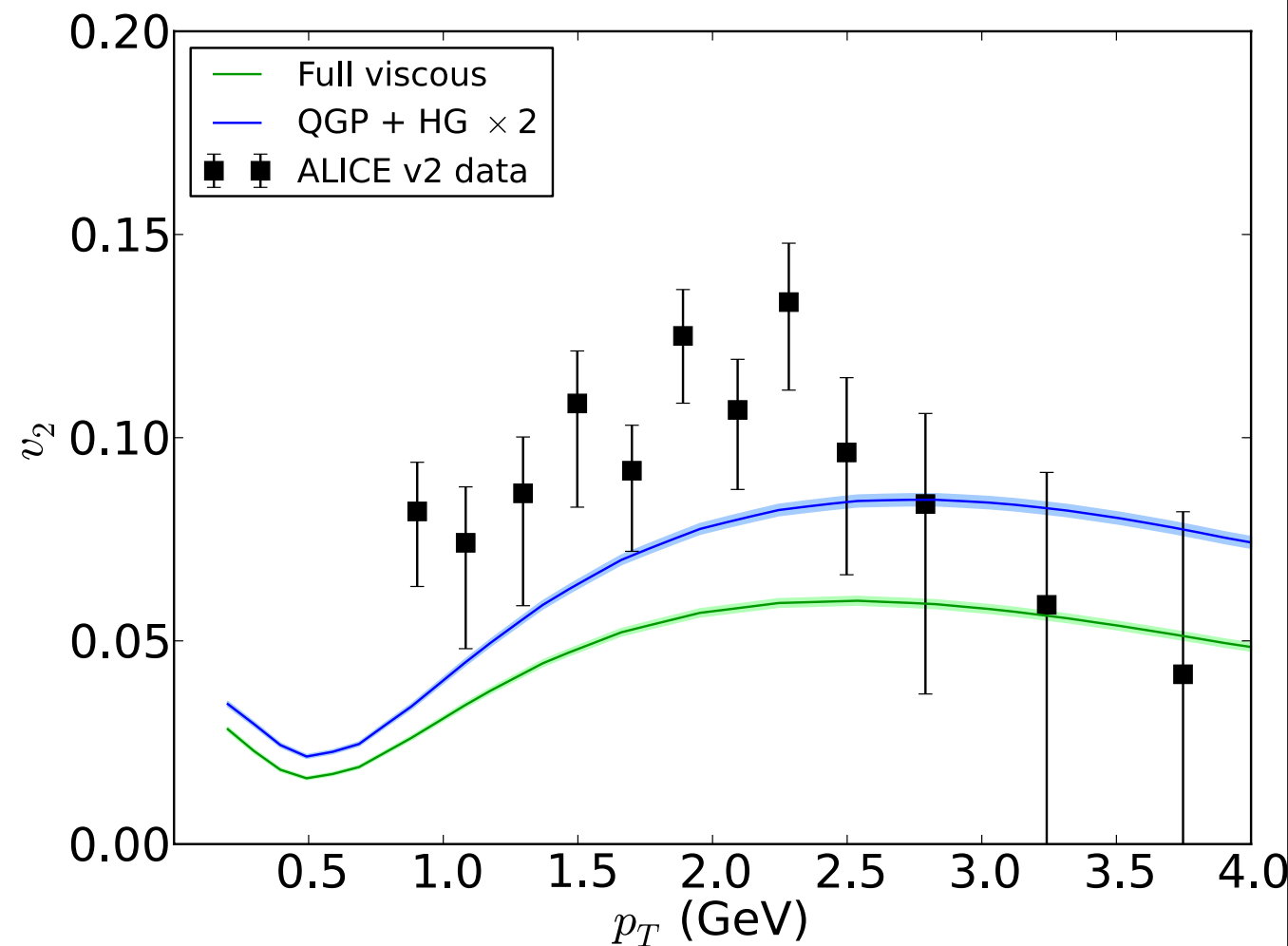
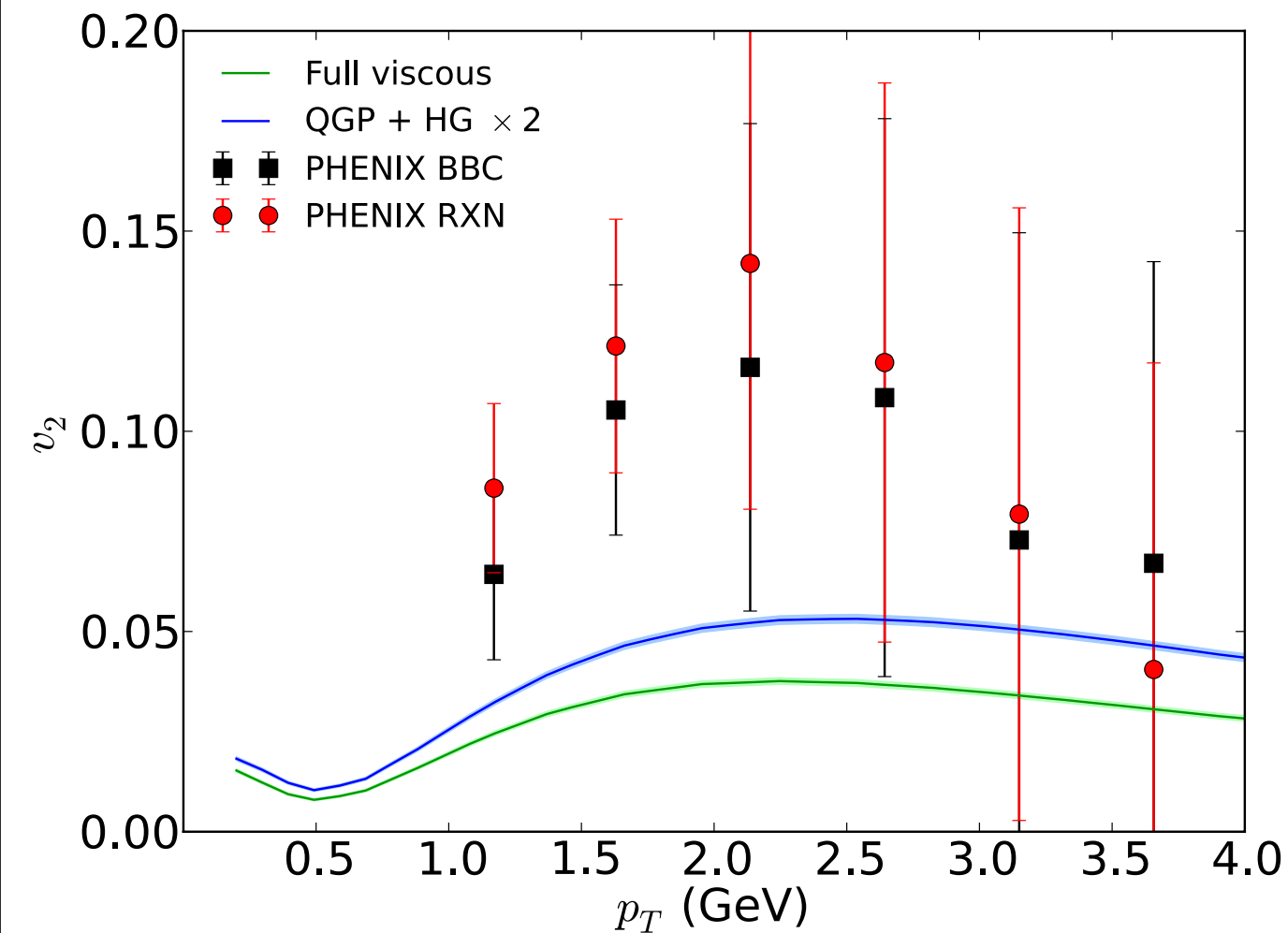


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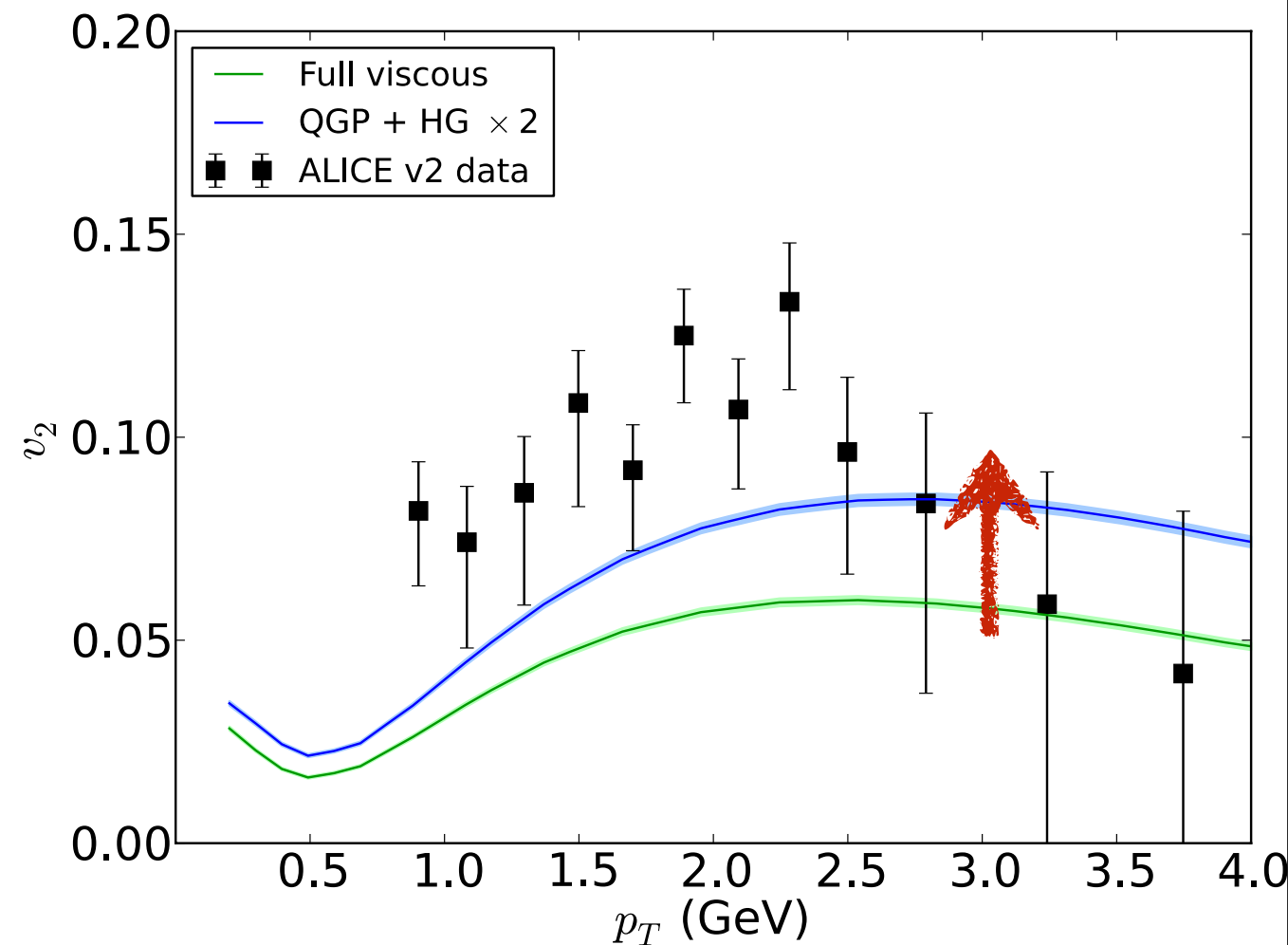
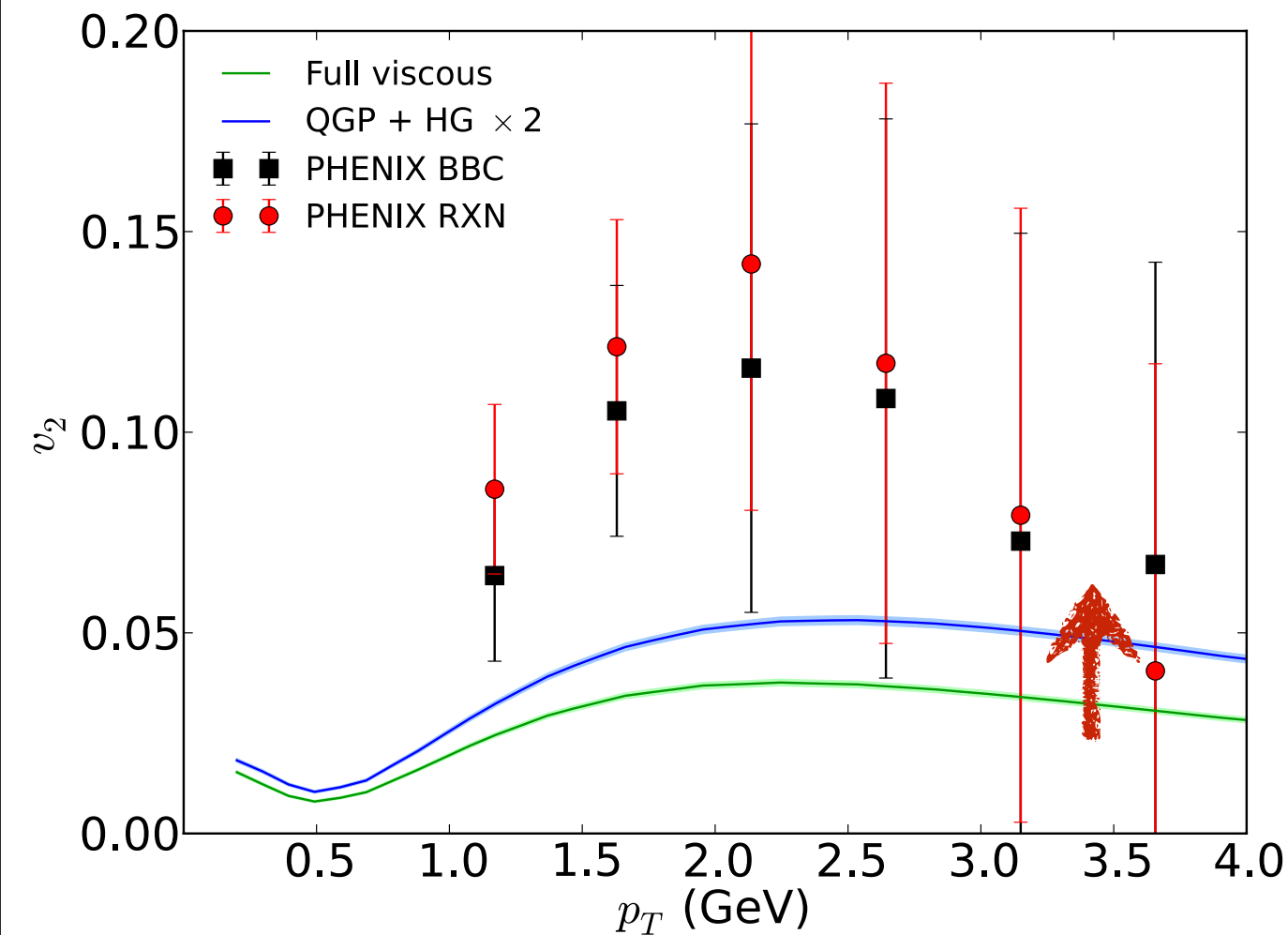
Missing rates in hadronic phase

Photon production rates from **baryonic channels** are missing in the hadronic phase. We can estimate this by increase photon emission rates in hadronic phase by a **factor of 2**,



Missing rates in hadronic phase

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- it increases total photon v_2 by **$\sim 45\%$** at both RHIC and LHC energies

Conclusion

- We calculate photon anisotropic flows $\mathbf{v}_n$ from *event-by-event* viscous hydrodynamic medium
- Shear viscosity *suppresses* photon v_n . Dominant suppression comes not from flow, but from the viscous correction to the production rates.
- **Elliptic** and **triangular** flow of photons are *more sensitive* than hadrons to the shear stress tensor at early time and the initial state fluctuations.
- Photon production from missing sources (e.g. *baryonic channels*) are needed to improve the agreement between experiment and theory.

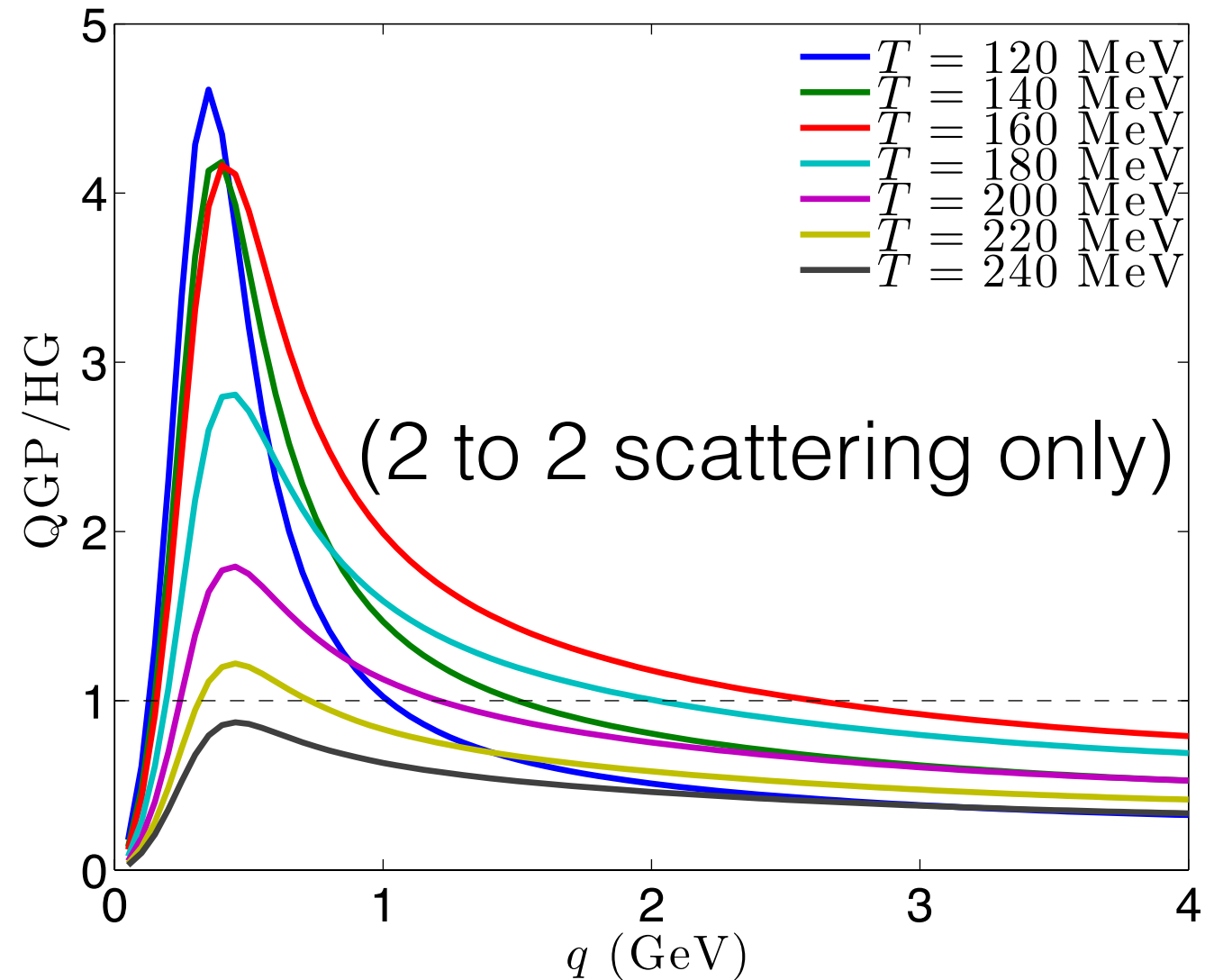
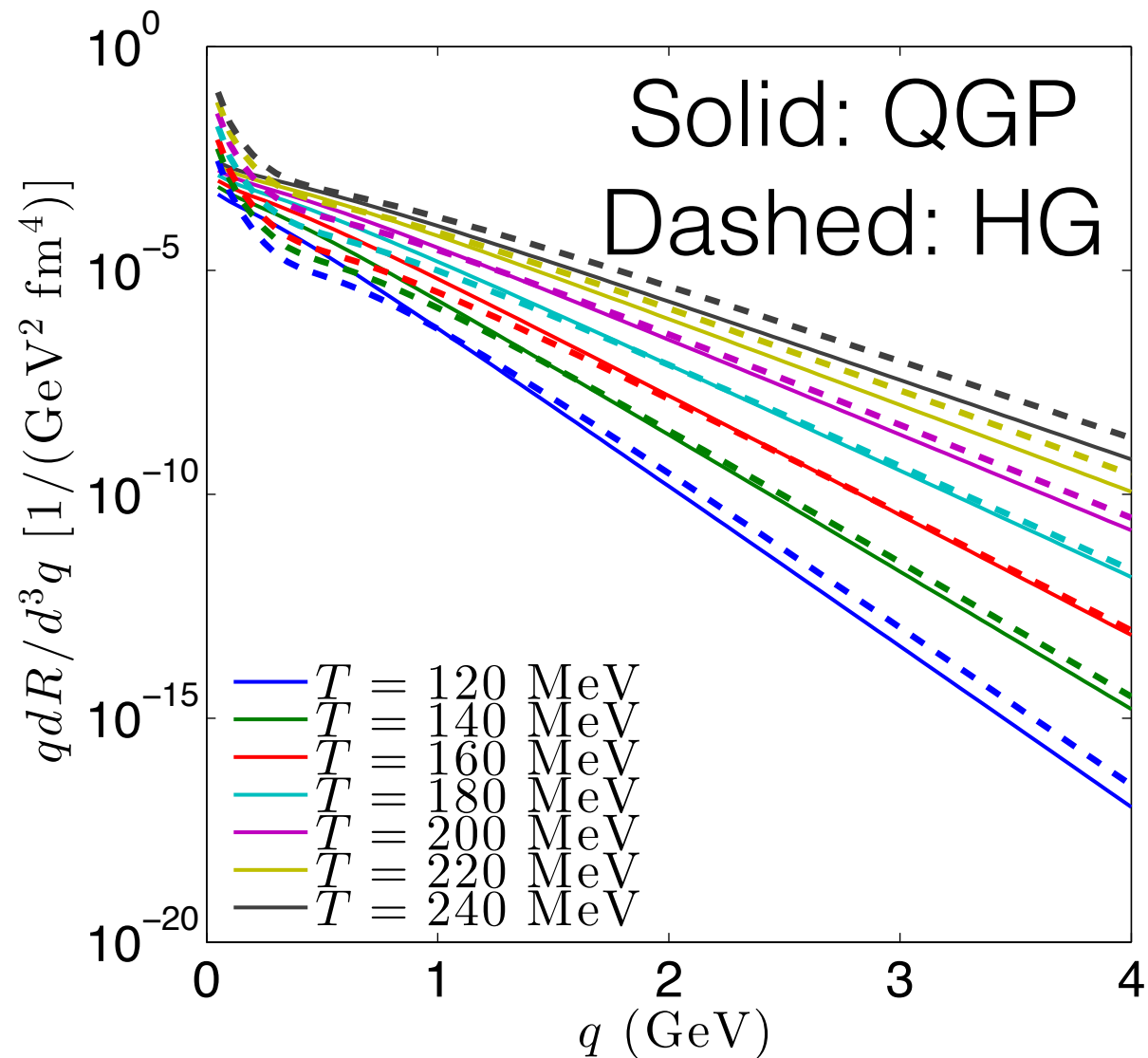
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Viscous effects on dilepton flow:
Gojko Vujanovic Tue 2:30

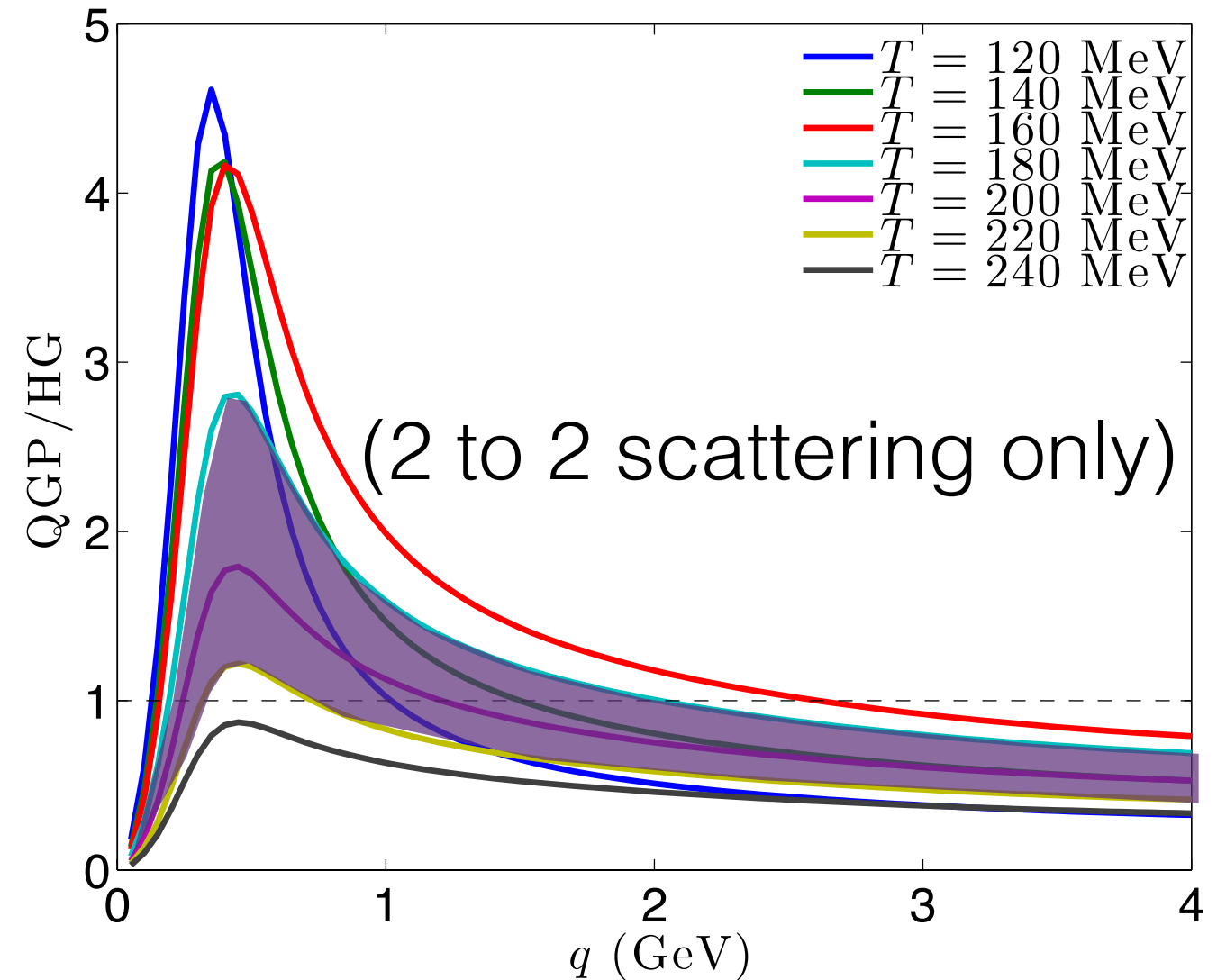
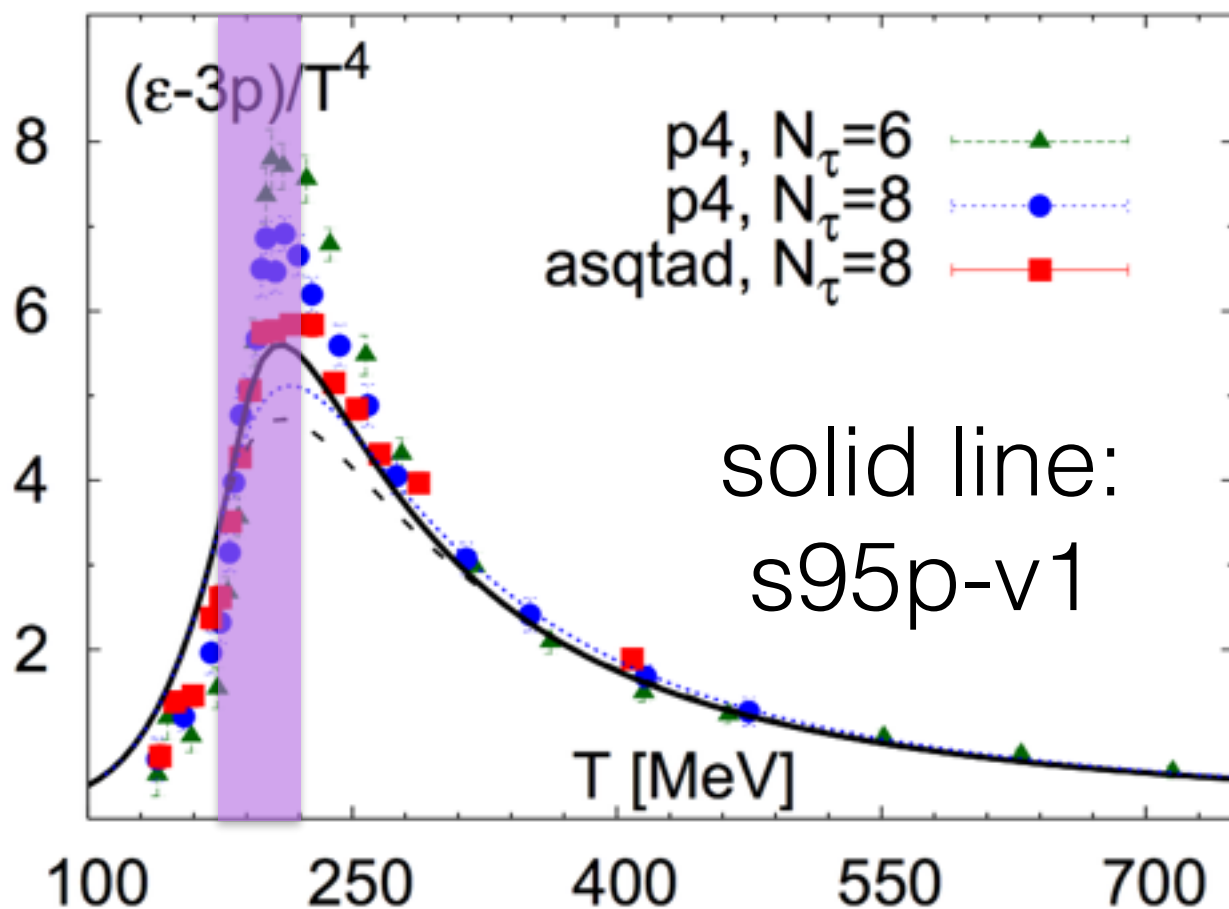
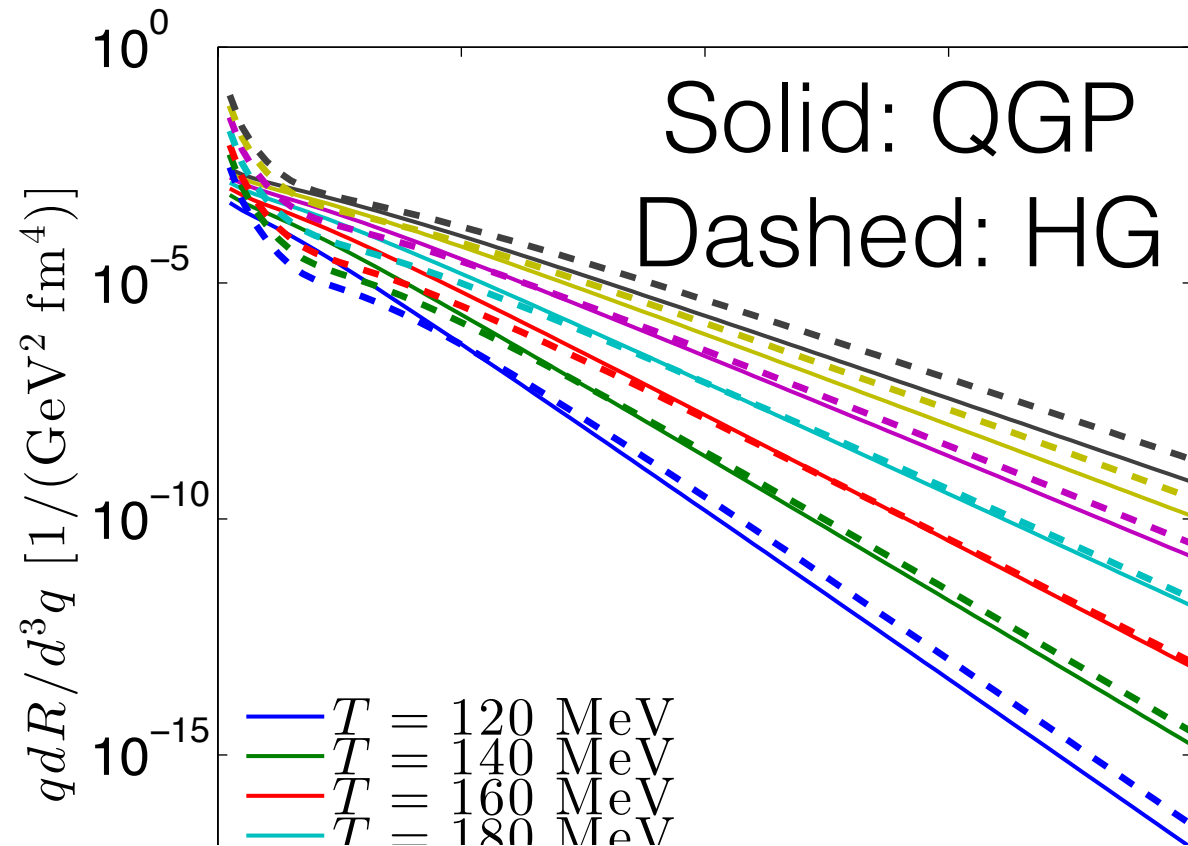
Back ups ...

Photon Emission Rates QGP vs HG



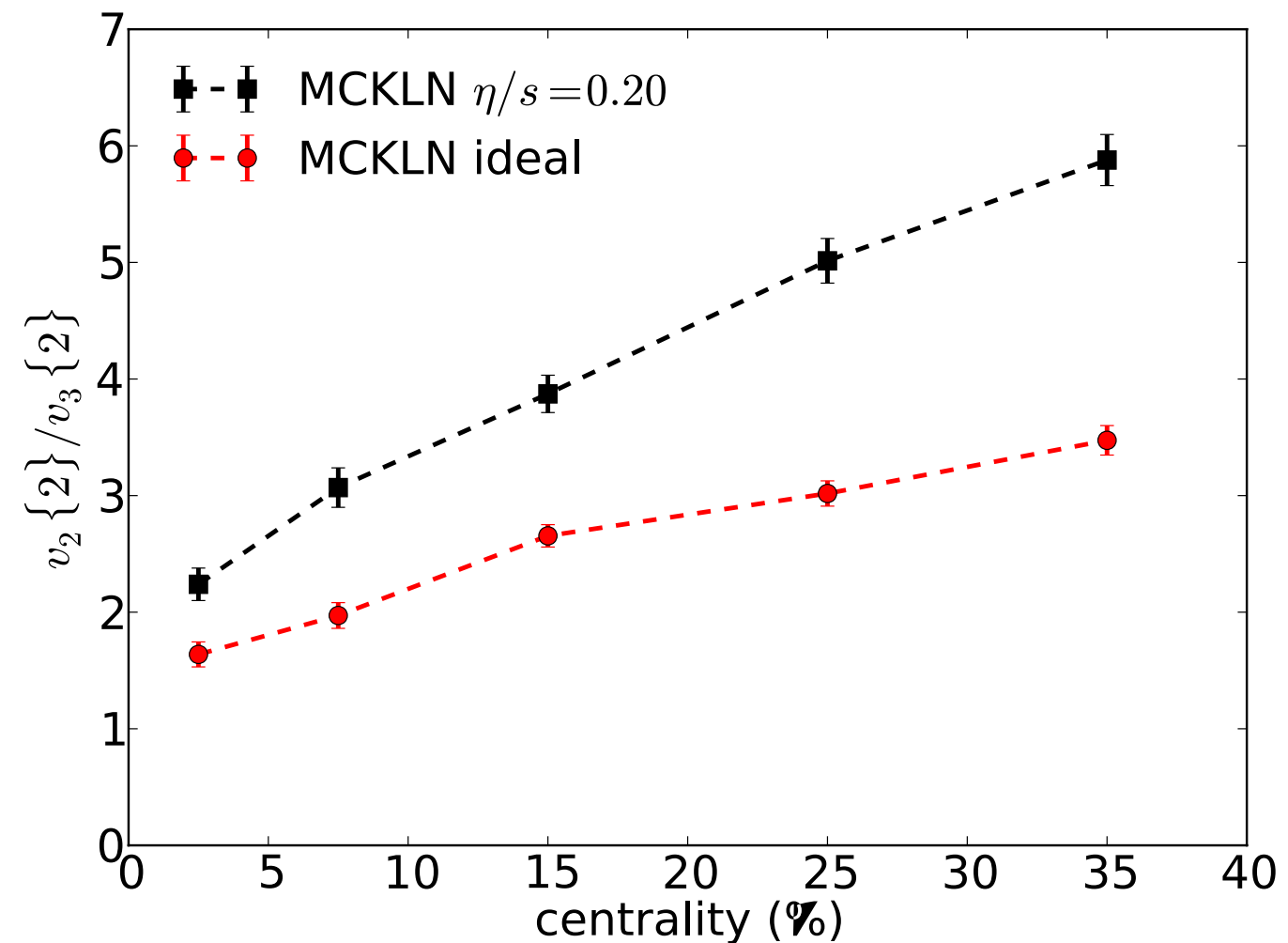
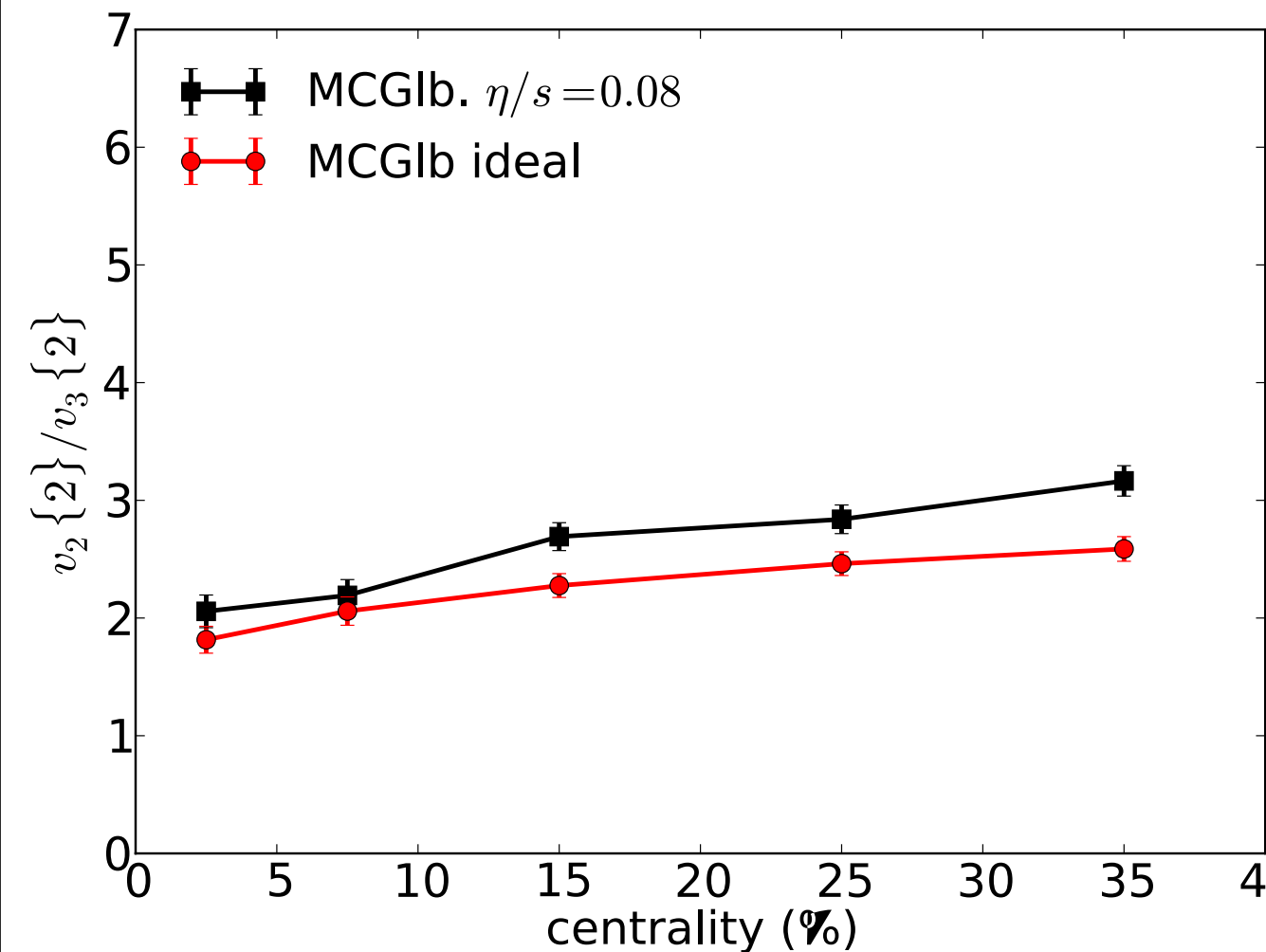
- QGP rates have very different p_T dependence compared to HG rates

Photon Emission Rates QGP vs HG



- QGP rates have very different p_T dependence compared to HG rates
- Estimated transition region for production rates, $T \sim \mathbf{184 - 220}$ MeV

Event-by-Event Full Viscous Photon v_n



- Comparing with ideal hydro runs, the v_2/v_3 ratio increases with shear viscosity
- MCKLN model shows stronger centrality dependence than MCGIb model