

# The Energy Momentum Tensor associated with Hard Parton production in Finite Time

Ben Meiring  
& W A Horowitz

University of Cape Town  
*mrnben002@myuct.ac.za*

1 - 5 Dec, 2014

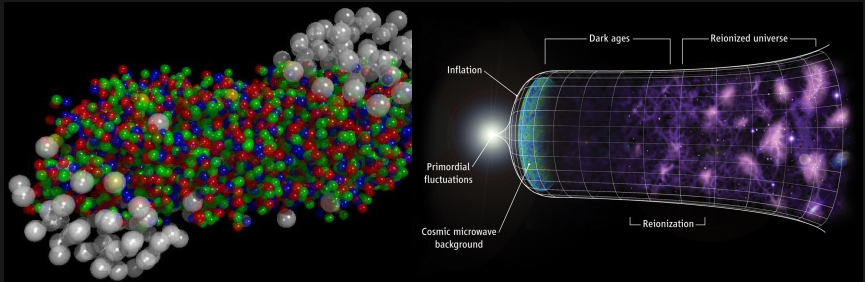


SA CERN



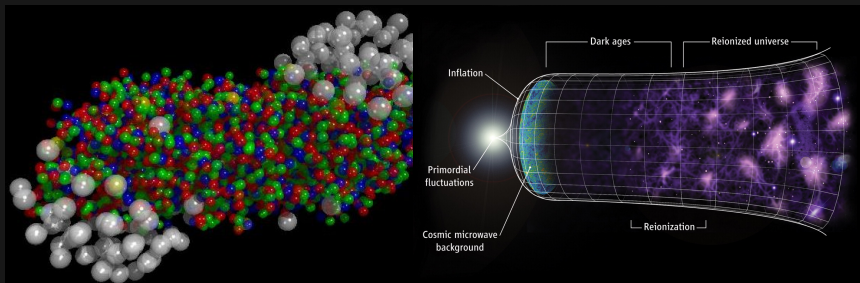
# What is the Quark Gluon Plasma?

- ▶ A new Fundamental state of matter, that we don't yet fully understand
- ▶ Dominated the universe at  $\sim 10^{-34}$  seconds after the big bang



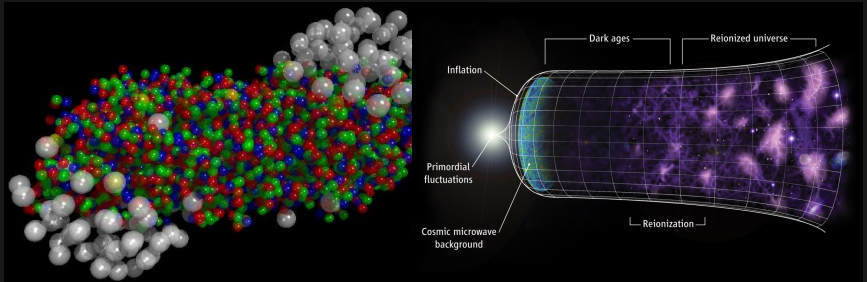
# What is the Quark Gluon Plasma?

- ▶ A new Fundamental state of matter, that we don't yet fully understand
- ▶ Dominated the universe at  $\sim 10^{-34}$  seconds after the big bang



# What is the Quark Gluon Plasma?

- ▶ A new Fundamental state of matter, that we don't yet fully understand
- ▶ Dominated the universe at  $\sim 10^{-34}$  seconds after the big bang



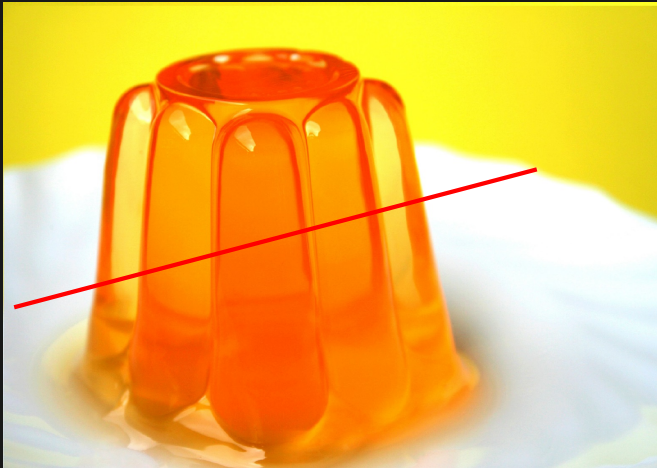
# How do we study the QGP?

Design an experiment to probe this Jello.



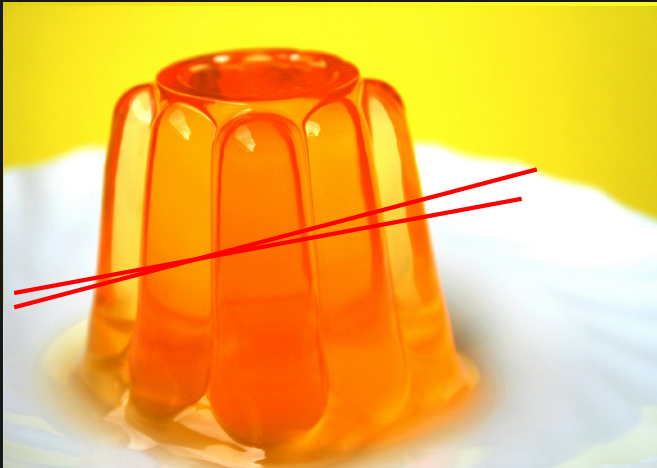
# How do we study the QGP?

Design an experiment to probe this Jello.



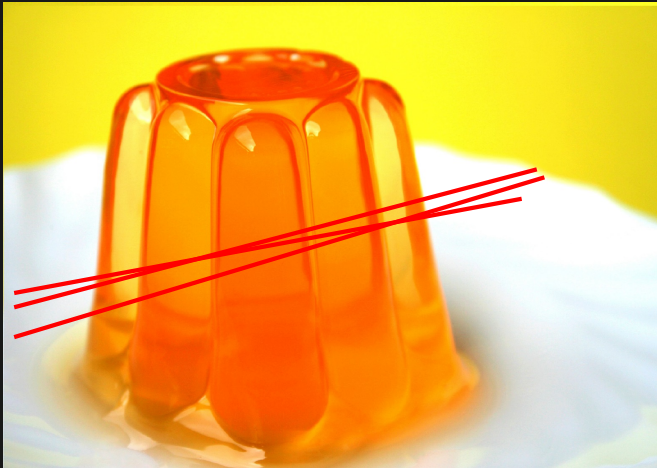
# How do we study the QGP?

Design an experiment to probe this Jello.



# How do we study the QGP?

Design an experiment to probe this Jello.





# How do we study the QGP?

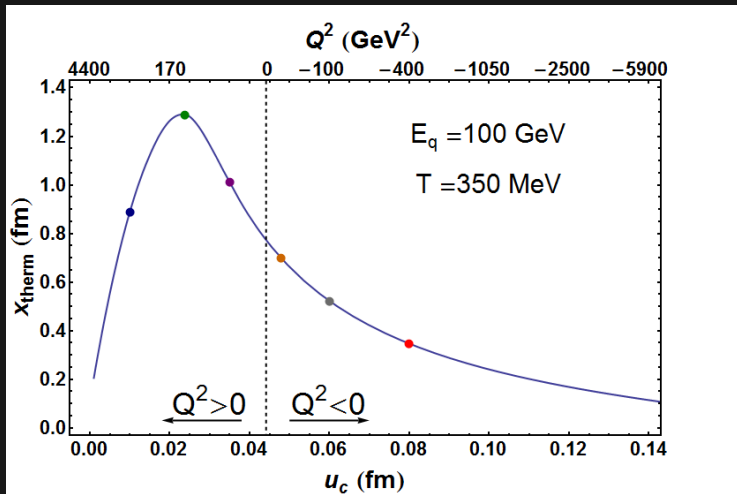


Figure: Stopping Distance as a function of radial  $u_c$ . R. Morad & W A Horowitz (1409.7545)

Toy Problem: Find the  $T_{\mu\nu}$  of a “gluon” emitted from a static “quark” in scalar QCD

Toy Problem: Find the  $T_{\mu\nu}$  of a “gluon” emitted from a static “quark” in scalar QCD

$$\langle \psi | T_{\mu\nu}(x) | \psi \rangle$$

Toy Problem: Find the  $T_{\mu\nu}$  of a “gluon” emitted from a static “quark” in scalar QCD

$$\langle \psi | T_{\mu\nu}(x) | \psi \rangle$$

The **expectation value** of the Energy Momentum Tensor, given an initial state of a single  $\psi$  particle at  $t = -\infty$ .

Toy Problem: Find the  $T_{\mu\nu}$  of a “gluon” emitted from a static “quark” in scalar QCD

$$\langle \psi | T_{\mu\nu}(x) | \psi \rangle$$

The **expectation value** of the Energy Momentum Tensor, given an initial state of a single  $\psi$  particle at  $t = -\infty$ .

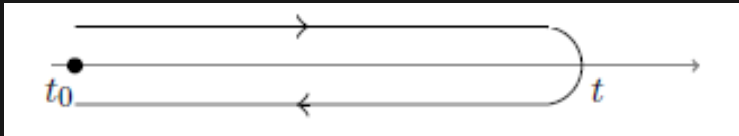


Figure: Schwinger-Keldysh Contour

Toy Problem: Find the  $T_{\mu\nu}$  of a “gluon” emitted from a static “quark” in scalar QCD

$$\langle \psi | T_{\mu\nu}(x) | \psi \rangle$$

The **expectation value** of the Energy Momentum Tensor, given an initial state of a single  $\psi$  particle at  $t = -\infty$ .

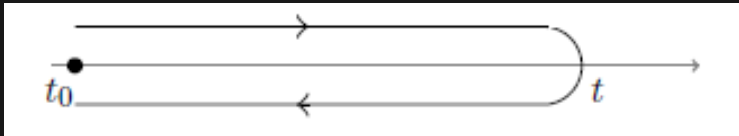


Figure: Schwinger-Keldysh Contour

$$T_{\mu\nu} = \partial_\mu \phi \partial_\nu \phi + \partial_\mu \psi \partial_\nu \psi - g_{\mu\nu} \mathcal{L}$$

$$\mathcal{L} = \frac{1}{2} (\partial\psi)^2 - \frac{1}{2} m_\psi^2 \psi^2 + \frac{1}{2} (\partial\phi)^2 - \frac{1}{2} m_\phi^2 \phi^2 - g\psi\phi\psi$$

Complications:  $\langle \psi | T_{\mu\nu}(x) | \psi \rangle$  is divergent (even after renormalization)

Complications:  $\langle \psi | T_{\mu\nu}(x) | \psi \rangle$  is divergent (even after renormalization)

We need to include an **improvement term**

$$\Theta_{\mu\nu} = T_{\mu\nu} - \frac{1}{4} \frac{(n-2)}{(n-1)} (\partial_\mu \partial_\nu - g_{\mu\nu}) \phi^2$$

Callan, Coleman, Jackiw. ANNALS OF PHYSICS: 59, 42-73 (1970)



Complications:  $\langle \psi | T_{\mu\nu}(x) | \psi \rangle$  is divergent (even after renormalization)

We need to include an **improvement term**

$$\Theta_{\mu\nu} = T_{\mu\nu} - \frac{1}{4} \frac{(n-2)}{(n-1)} (\partial_\mu \partial_\nu - g_{\mu\nu}) \phi^2$$

Callan, Coleman, Jackiw. ANNALS OF PHYSICS: 59, 42-73 (1970)

But we will use a **sneaky trick**

Complications:  $\langle \psi | T_{\mu\nu}(x) | \psi \rangle$  is divergent (even after renormalization)

We need to include an **improvement term**

$$\Theta_{\mu\nu} = T_{\mu\nu} - \frac{1}{4} \frac{(n-2)}{(n-1)} (\partial_\mu \partial_\nu - g_{\mu\nu}) \phi^2$$

Callan, Coleman, Jackiw. ANNALS OF PHYSICS: 59, 42-73 (1970)

But we will use a **sneaky trick**

$$\begin{aligned} \langle \phi(x) \phi(x) \rangle &= \langle \phi(x) \rangle^2 + (\Delta \phi(x))^2 \\ &\approx \langle \phi(x) \rangle^2 \end{aligned}$$

Complications:  $\langle \psi | T_{\mu\nu}(x) | \psi \rangle$  is divergent (even after renormalization)

We need to include an **improvement term**

$$\Theta_{\mu\nu} = T_{\mu\nu} - \frac{1}{4} \frac{(n-2)}{(n-1)} (\partial_\mu \partial_\nu - g_{\mu\nu}) \phi^2$$

Callan, Coleman, Jackiw. ANNALS OF PHYSICS: 59, 42-73 (1970)

But we will use a **sneaky trick**

$$\begin{aligned} \langle \phi(x) \phi(x) \rangle &= \langle \phi(x) \rangle^2 + (\Delta \phi(x))^2 \\ &\approx \langle \phi(x) \rangle^2 \end{aligned}$$

So we find  $\langle \phi(x) \rangle$  and plug in

$$\langle T_{\mu\nu}(x) \rangle \approx T_{\mu\nu} (\langle \phi(x) \rangle)$$

Approximation: Calculate  $\langle \phi(x) \rangle$  and Plug into

$$T_{\mu\nu}$$

# Approximation: Calculate $\langle \phi(x) \rangle$ and Plug into $T_{\mu\nu}$

We found a **general expression** for  $\langle \phi(x) \rangle$  with  $H_{\text{int}} = g\psi\phi\psi$

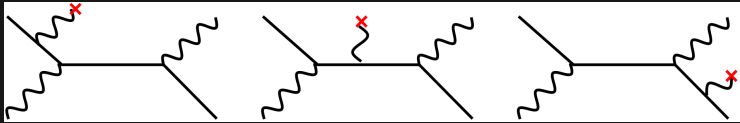


Figure: Typical Diagrams for  $\langle \phi(x) \rangle$  for  $|\text{in}\rangle = |\psi\phi\rangle$

# Approximation: Calculate $\langle \phi(x) \rangle$ and Plug into $T_{\mu\nu}$

We found a **general expression** for  $\langle \phi(x) \rangle$  with  $H_{\text{int}} = g\psi\phi\psi$

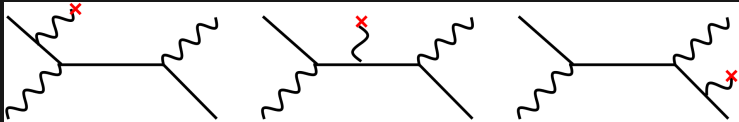


Figure: Typical Diagrams for  $\langle \phi(x) \rangle$  for  $|\text{in}\rangle = |\psi\phi\rangle$

$$\langle \phi(x) \rangle = -ig \int d^4z D_R(x-z) \langle \text{in} | \psi(z) \psi(z) | \text{in} \rangle$$

(Generalizes for arbitrary  $H_{\text{int}}$ )

Approximation: Calculate  $\langle \phi(x) \rangle$  and Plug into

$T_{\mu\nu}$

Choosing gaussian smeared wavepackets for the  $|\text{in}\rangle = |\psi\rangle$  state with momentum width  $1/\alpha$ . We take  $m_\psi \rightarrow \infty$

$$\langle \phi(x) \rangle = \frac{g}{m_\psi} \int d^3z \left( \frac{e^{-m_\phi |\vec{z}|}}{|\vec{z}|} \right) \frac{e^{-\frac{(\vec{x}-\vec{z})^2}{\alpha}}}{\sqrt{4\pi\alpha}^3}$$

Approximation: Calculate  $\langle \phi(x) \rangle$  and Plug into

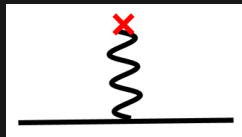
$T_{\mu\nu}$

Choosing gaussian smeared wavepackets for the  $|\text{in}\rangle = |\psi\rangle$  state with momentum width  $1/\alpha$ . We take  $m_\psi \rightarrow \infty$

$$\langle \phi(x) \rangle = \frac{g}{m_\psi} \int d^3 z \left( \frac{e^{-m_\phi |\vec{z}|}}{|\vec{z}|} \right) \frac{e^{-\frac{(\vec{x}-\vec{z})^2}{\alpha}}}{\sqrt{4\pi\alpha}^3}$$

Solution to the **Diffusion equation**

$$\partial_\alpha \langle \phi(x, \alpha) \rangle = \partial_{\vec{x}}^2 \langle \phi(x, \alpha) \rangle$$





Approximation: Calculate  $\langle \phi(x) \rangle$  and Plug into

$T_{\mu\nu}$

Choosing gaussian smeared wavepackets for the  $|\text{in}\rangle = |\psi\rangle$  state with momentum width  $1/\alpha$ . We take  $m_\psi \rightarrow \infty$

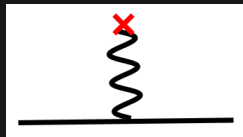
$$\langle \phi(x) \rangle = \frac{g}{m_\psi} \int d^3z \left( \frac{e^{-m_\phi |\vec{z}|}}{|\vec{z}|} \right) \frac{e^{-\frac{(\vec{x}-\vec{z})^2}{\alpha}}}{\sqrt{4\pi\alpha}^3}$$

Solution to the **Diffusion equation**

$$\partial_\alpha \langle \phi(x, \alpha) \rangle = \partial_{\vec{x}}^2 \langle \phi(x, \alpha) \rangle$$

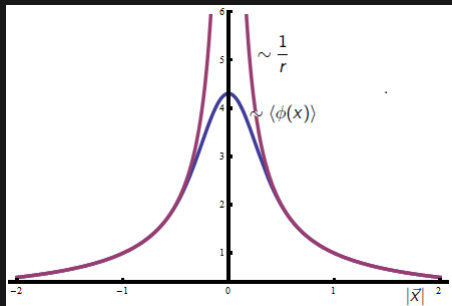
with initial condition

$$\phi(x, 0) = \frac{e^{-m_\phi |\vec{x}|}}{|\vec{x}|}.$$



Approximation: Calculate  $\langle \phi(x) \rangle$  and Plug into  $T_{\mu\nu}$

$$\langle \phi(x) \rangle = \frac{g}{m_\psi} \int d^3 z \left( \frac{e^{-m_\phi |\vec{z}|}}{|\vec{z}|} \right) \frac{e^{-\frac{(\vec{x}-\vec{z})^2}{\alpha}}}{\sqrt{4\pi\alpha}^3}$$



# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state

# Conditional Expectation Value

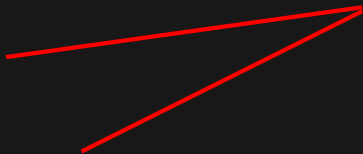
- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state

# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state

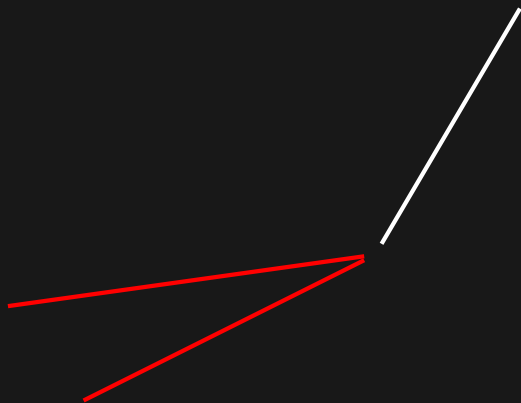
# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state



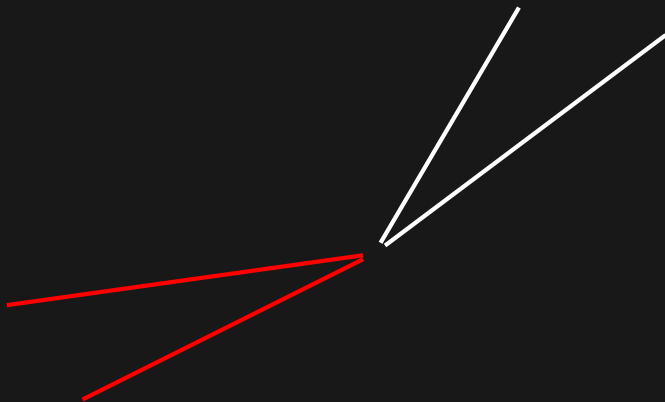
# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state



# Conditional Expectation Value

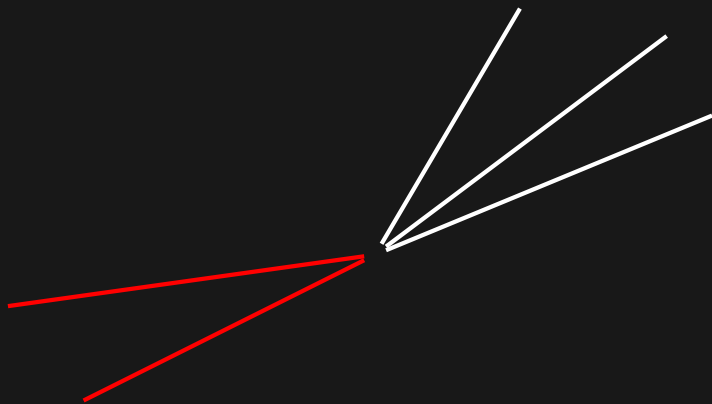
- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state





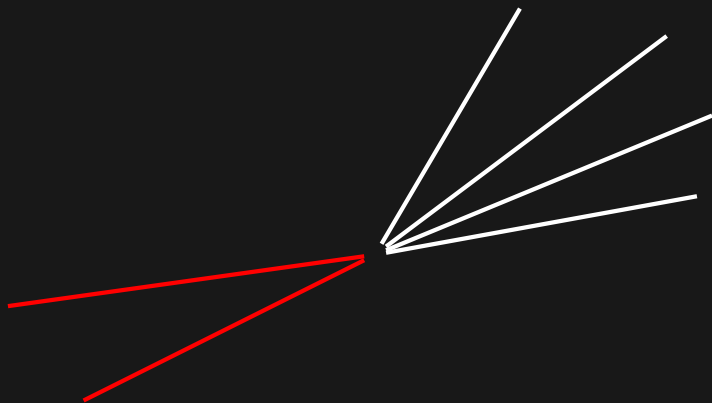
# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state



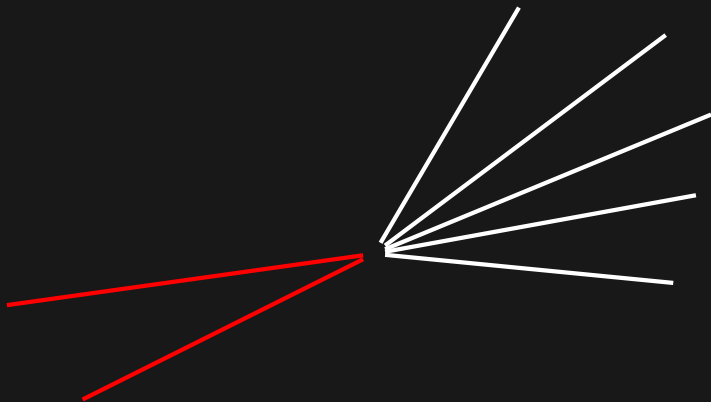
# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state



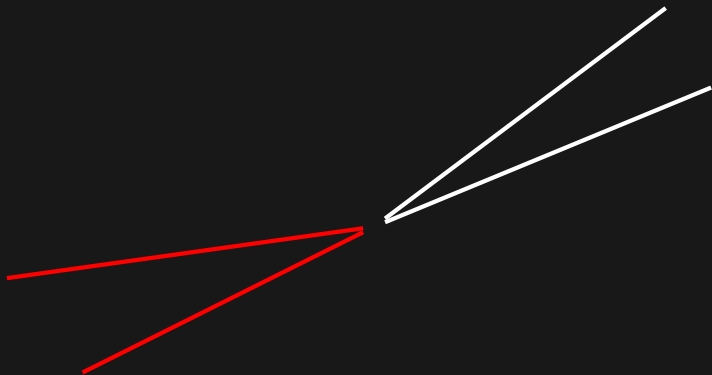
# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state



# Conditional Expectation Value

- ▶ Specifying only the initial state gave us a Yukawa
- ▶ We want to refine our result by specifying the final state



# Conditional Expectation Value

We define the **Conditional Expectation Value**

# Conditional Expectation Value

We define the **Conditional Expectation Value**

$$\begin{aligned} E[\hat{O}(x) | |\text{in}\rangle, |\text{out}\rangle] &= \sum_i O_i P(|q_i\rangle | |\text{in}\rangle, |\text{out}\rangle) \\ &= \frac{\langle \text{in} | \hat{\Theta}_M \hat{O}(x) \hat{\Theta}_M | \text{in} \rangle}{\langle \text{in} | \hat{\Theta}_M | \text{in} \rangle} \end{aligned}$$

where  $\hat{\Theta}_M = |\text{out}\rangle\langle \text{out}|$  is the **projection operator** built from the out states.

# Conditional Expectation Value

We define the **Conditional Expectation Value**

$$\begin{aligned} E[\hat{O}(x) | |\text{in}\rangle, |\text{out}\rangle] &= \sum_i O_i P(|q_i\rangle | |\text{in}\rangle, |\text{out}\rangle) \\ &= \frac{\langle \text{in} | \hat{\Theta}_M \hat{O}(x) \hat{\Theta}_M | \text{in} \rangle}{\langle \text{in} | \hat{\Theta}_M | \text{in} \rangle} \end{aligned}$$

where  $\hat{\Theta}_M = |\text{out}\rangle\langle \text{out}|$  is the **projection operator** built from the out states.

This result is actually a generalization of Baye's Theorem.

# Conclusion



# Conclusion

- ▶ We've explored the possibility of finite time QFT calculations

# Conclusion

- ▶ We've explored the possibility of finite time QFT calculations
- ▶ We found an expression for  $T_{\mu\nu}$  through  $\langle\phi(x)\rangle$

# Conclusion

- ▶ We've explored the possibility of finite time QFT calculations
- ▶ We found an expression for  $T_{\mu\nu}$  through  $\langle\phi(x)\rangle$
- ▶ We defined a concept of Conditional Expectation Values in QFT

## Conclusion

- ▶ We've explored the possibility of finite time QFT calculations
- ▶ We found an expression for  $T_{\mu\nu}$  through  $\langle\phi(x)\rangle$
- ▶ We defined a concept of Conditional Expectation Values in QFT

The End.

# Back up Slides

# Data for QGP

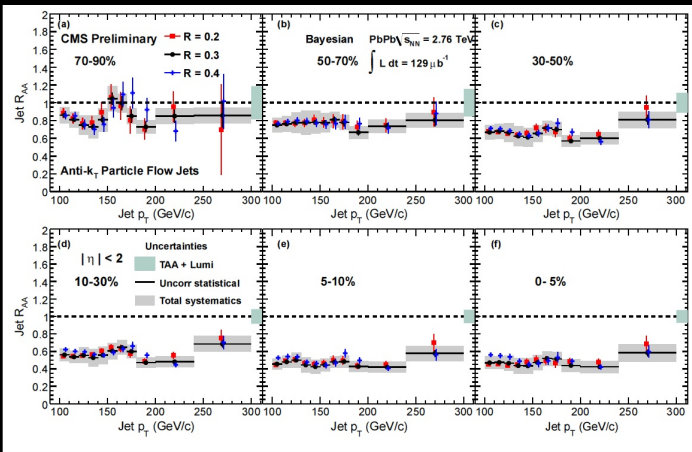


Figure: CMS Preliminary data for Jet  $R_{AA}$ . Taken from (1409.7545)

# Conditional Expectation Value: Example

For  $|\text{in}\rangle = |\psi\rangle$ ,  $|\text{out}\rangle = |\psi'\rangle$  we find

$$E[\phi(x) | |\psi\rangle, |\psi'\rangle] = \langle\psi' | \phi(x) | \psi'\rangle + 2i \text{Im} \left( \frac{\langle\psi' | \phi(x) | \psi\rangle}{|\langle\psi' | \psi\rangle|} \right)$$

(At least to first order).

# Conditional Expectation Value: Example

For  $|\text{in}\rangle = |\psi\rangle$ ,  $|\text{out}\rangle = |\psi'\rangle$  we find

$$E[\phi(x) | |\psi\rangle, |\psi'\rangle] = \langle\psi' | \phi(x) | \psi'\rangle + 2i \text{Im} \left( \frac{\langle\psi' | \phi(x) | \psi\rangle}{|\langle\psi' | \psi\rangle|} \right)$$

(At least to first order).

This says that if  $|\psi\rangle \neq |\psi'\rangle$ , the expectation of the field is complex.