



Discovery Physics at the LHC Kruger Workshop 2014

tH analysis: the feasibility of lifting top Yukawa degeneracies with ATLAS

Discovery Physics at the LHC Kruger Workshop 2014

Claire Antel

as part of my Masters
supervisor and co-supervisor:

Dr Andrew Hamilton and Dr Heather Gray



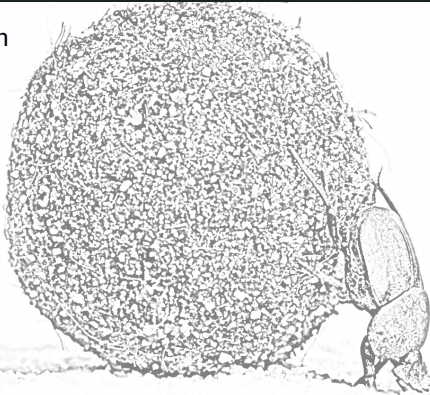
University of Cape Town

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Introduction



With the LHC's first run, enough data has been accumulated to uncover the predicted Higgs boson. Focus of searches has shifted to measuring its properties – is it THE Standard Model Higgs?

Thus far, all measurements have indicated a compliance with the Standard Model. However, there remain non-SM areas that have not been entirely excluded.

These are possible windows to new physics.



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fermionic and bosonic couplings

One of the Higgs properties that are of interest are its *coupling* strengths to other SM particles.

Using all available data at 7 and 8 TeV, ATLAS and CMS have presented constraints on the fermion and vector boson couplings.

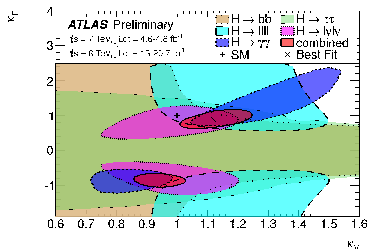


Figure : Correlation between coupling scale factors c_F (fermionic) and c_V (vector bosonic) with 68% CL contours from individual channels at ATLAS.¹

¹ATLAS collab. Combined coupling measurements of the Higgs-like boson with the ATLAS detector using up to 25 fb⁻¹ of proton-proton collision data. ATLAS-CONF-2013-03



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$$\kappa_V \equiv g_V/g_V^{SM} \quad \kappa_F \equiv g_f/g_f^{SM}$$

...essentially quantify the deviation from g^{SM} .

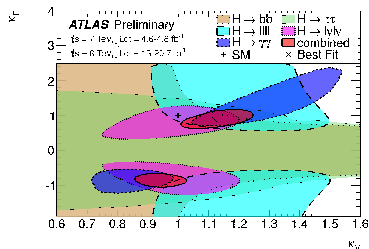


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...essentially quantify the deviation from g^{SM} .

But: There are **two** minima in the fits – one complying with SM (1:1); the other in a region of new physics – at negative values.

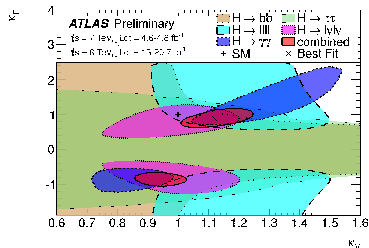
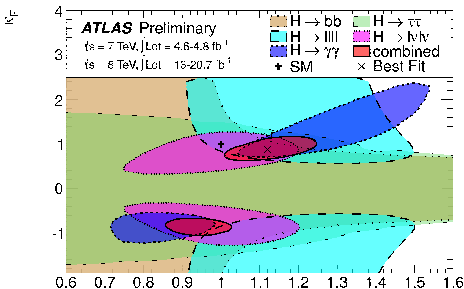


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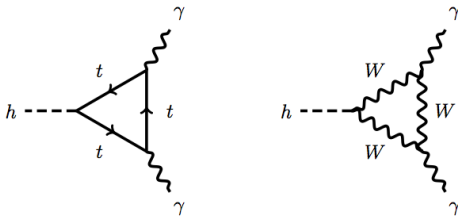


We need more sensitivity in the *relative sign* of κ_F ...



most channels not sensitive to the relative sign between κ_V and κ_F , as

$$\mathcal{N} \propto \frac{\kappa_x^2 \kappa_y^2}{\kappa_H} \quad (1)$$



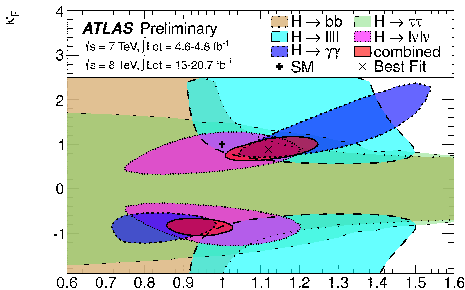
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However, $H \rightarrow \gamma\gamma$ **loop-induced** $\rightarrow \kappa_\gamma$ function of κ_F and κ_V – scales as $|\alpha\kappa_V + \beta\kappa_F|^2$.



degeneracies in fermionic coupling fits



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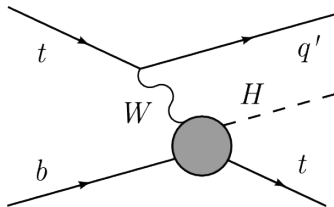
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However, $H \rightarrow \gamma\gamma$ loop-induced $\rightarrow \kappa_\gamma$ function of κ_F and κ_V – scales as $|\alpha\kappa_V + \beta\kappa_F|^2$.

...interference term $\alpha\beta\kappa_F\kappa_V$ emerges, resulting in fits also favouring negative signs.



Resolving the degeneracies: the tH process

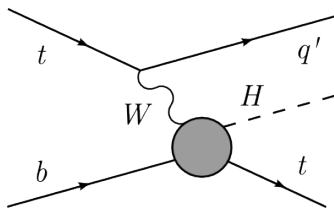


The *single top production* in association with the Higgs, tH , offers a unique opportunity to study the *relative sign* of the Yukawa coupling.

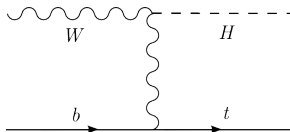
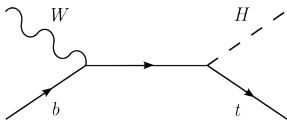


Resolving the degeneracies: the tH process

This is due to an interference...



...between bosonic and fermionic processes...





Due to the interference the cross-section of the tH process is **enhanced** for non-SM values...

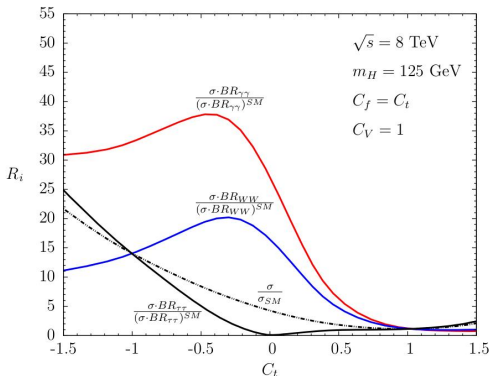
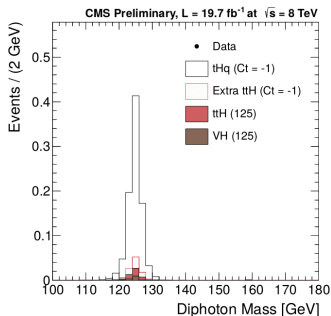


Figure : Cross section ratio as function of the top yukawa coupling, $c_t = c_F$, $c_V = 1$.²

²Sanjoy Biswas, Emidio Gabrielli, Fabrizio Margaroli, and Barbara Mele. Direct constraints on the top-Higgs coupling from the 8 TeV LHC data. JHEP, 07:073, 2013.



An update on measurements of the Single Top channel



- CMS: $H \rightarrow \gamma\gamma$ at 8 TeV/ 19.6 fb^{-1} - reported zero events passing acceptance cuts and 95% CL on $H \rightarrow \gamma\gamma$ event yield of 4.1 times the expectation.³
- ATLAS: $H \rightarrow \gamma\gamma$ cross-sections of $t\bar{t}H$ and tH channels combined at 7 and 8 TeV - strongly constrained negative Yukawa sector; 95% CL for c_F at -1.3 and +8.0.⁴


³CMS-PAS-HIG-14-001

⁴CERN-PH-EP-2014-179



This is a **feasibility study** based on

Lifting degeneracies in Higgs couplings using single top production in association with a Higgs boson

Marco Farina^a, Christophe Grojean^{b,c}, Fabio Maltoni^d,
Ennio Salvioni^{c,e} and Andrea Thamm^{c,f} 

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taking $H \rightarrow b\bar{b}$.

The analysis is done at **truth level** using 8 TeV samples and assuming 21fb^{-1} of integrated luminosity.

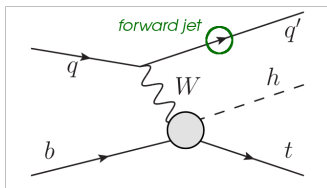
⁵Marco Farina, Christophe Grojean, Fabio Maltoni, Ennio Salvioni, and Andrea Thamm. Lifting degeneracies in Higgs couplings using single top production in association with a Higgs boson. JHEP, 1305:022, 2013.



Signal signature

Seek

- $H \rightarrow b\bar{b}$ channel (58%, dominant decay at this mass range)
- $t \rightarrow Wb$ (99.8%)
- $W \rightarrow l\nu$ (20%)
- a *forward jet* - characteristic high p_T jet, produced at Higgs and top production



Thus,

Signature

$$tHj \rightarrow 3 \text{ } b\text{-jets} + \ell(e^\pm/\mu^\pm) + j + MET$$



MC samples

Signal

Monte Carlo samples of tH process generated for $c_F = +1$ and $c_F = -1$ ($c_V = 1$).

Backgrounds

- $t\bar{b}j$ (*irreducible*)
- $tZj, Z \rightarrow b\bar{b}$ (*irreducible*)
- $t\bar{t}$, where one $Wb \rightarrow c\bar{s}b$, the c is mistagged and the s is taken as a forward jet (*reducible*)
- $t\bar{t}j$, where one $Wb \rightarrow c\bar{s}b$, the c or s is mistagged (*reducible*)

There is (at least) one leptonically decaying top in every event.

cross-sections

Process	Cross-Section (fb ⁻¹)	Generator
$c_F = +1$ tH	1.66	MADGRAPH + PYTHIA
$c_F = -1$ tH	27.81	MADGRAPH + PYTHIA
$t\bar{b}j$	11.28	MADGRAPH + PYTHIA
tZj	6.27	MADGRAPH + PYTHIA
$t\bar{t}(j)$	210.85×10^3	POWHEG+ PYTHIA

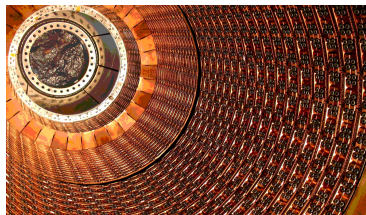


Event Selection

A simple event count is done by implementing certain selection criteria for the signal and background.

Acceptance cuts are made from the outset, motivated by

- suppression of low-momentum pileup and underlying event particles
- keeping the acceptance area to regions where the detector is of optimum efficiency.



Acceptance cuts are $p_T > 25$ GeV and $|\eta| < 2.5$ on all jets and leptons (except for forward jet).



summary of selection criteria

- $p_T^b, p_T^\ell, p_T^j > 25 \text{ GeV}$
- $\eta^\ell, \eta^b < 2.5$
- 1 lepton: acts as trigger
- 1 forward jet \rightarrow forward jet selection:
if 1 of top 2 highest p_T non-b-tagged jets at $> 2.5\eta$.
- ≥ 3 b-jets
- a b-jet pair where $65 \text{ GeV} < m_{bb} < 145 \text{ GeV}$





electron/muon no. (trigger)

no. of $\ell = 1$ cut

no. of leptons/event	with acceptance cuts		
	0	1	2
$c_F = +1$:	48%	52%	0%
$c_F = -1$:	45%	55%	0%
$tbbj$	25%	75%	0%
$tZbbj$	50%	50%	0%
$t\bar{t}$	44%	51%	5%

Sensitivities

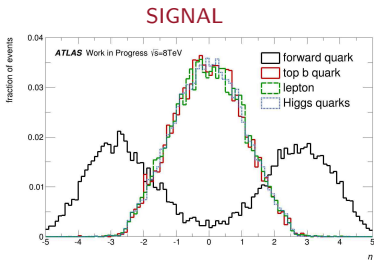
- $c_F = +1$: 0.01

- $c_F = -1$: 0.22

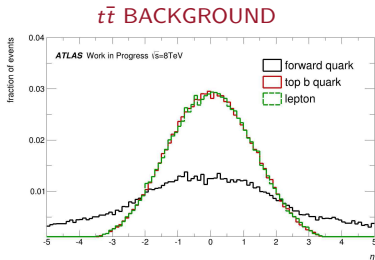


η comparisons

A fine distinction between signal and $t\bar{t}$ bg lies in the η distributions...



(a) η distribution for $c_F = -1$



(b) η distribution for $t\bar{t}$

- important characteristic of signal: high p_T jets at high η



η comparisons

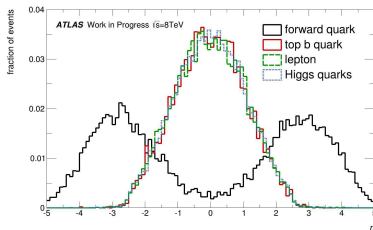


Figure : η distribution for $c_F = -1$

→ set requirement: one of top two highest p_T jets at $|\eta| > 2.5$

Sensitivities

- $c_F = +1$: 0.01
- $c_F = -1$: 0.27



b jet no.

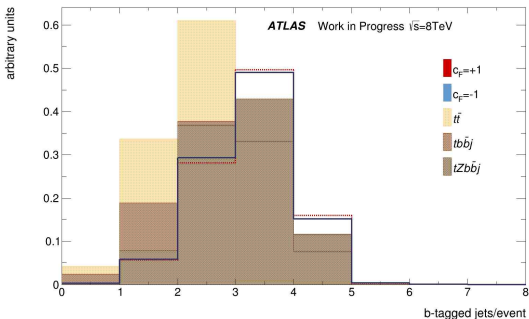


Figure : no. of b jets/event for all samples.

- $p_T > 25\text{GeV}$, $|\eta| < 2.5$ cuts applied.

Sensitivities

- $c_F = +1$: 0.08
- $c_F = -1$: 1.63

constraint on no. of b-jets (≥ 3)

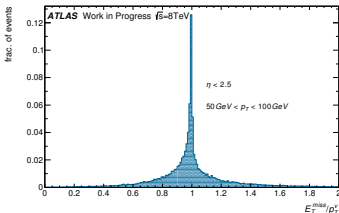
- 99% cut on $t\bar{t}$
- $\sim 34\%$ cut on signal



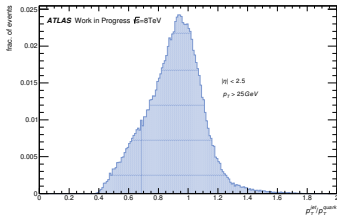
reconstructed Higgs mass

Salient difference between signal and all backgrounds is presence of a **Higgs boson**.

- In order to reduce b-jet combinations, first reconstructed W and top quark...
 - work with E_T^{miss} : estimate p_z^ν by setting a W mass constraint.
 - reconstruct top quark through W+b.



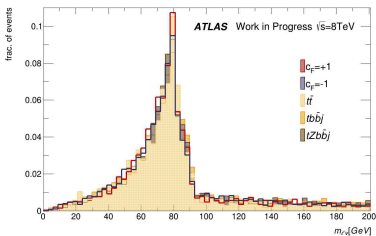
energy ratio between E_T^{miss} and truth neutrino



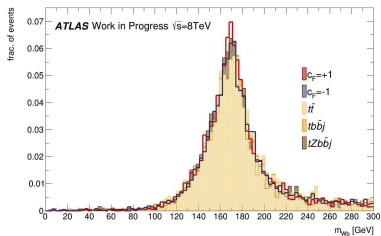
p_T ratio between truth b jet and truth b quark



reconstructed W boson and top quark



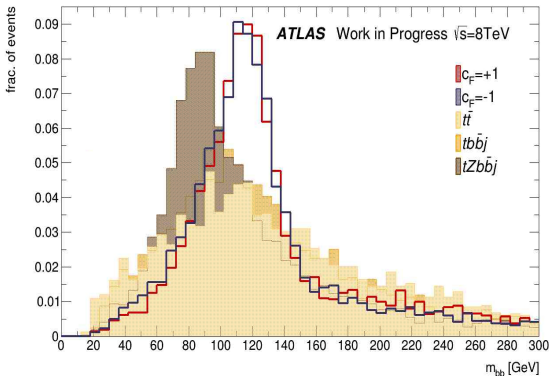
reconstructed W boson mass for all samples



reconstructed top mass for all samples



reconstructed Higgs



m_{bb} constraint

- $65 \text{ GeV} < m_{bb} < 145 \text{ GeV}$

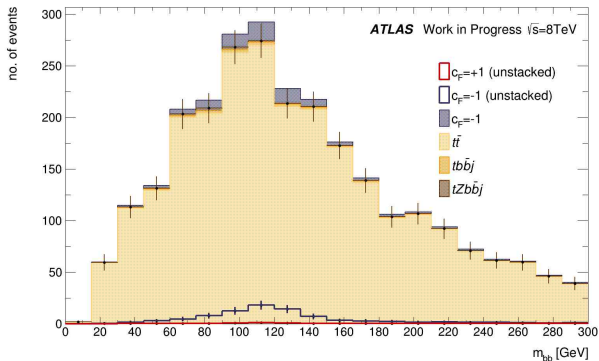
Sensitivities

- $C_F = +1$: 0.08
- $C_F = -1$: 1.74

- The Z boson is clearly evident at $\sim 90 \text{ GeV}$.
- The $t\bar{t}$ and $tbbj$ events are widely distributed.



Final expected distribution...



- an excess of 62 ± 1 over 1274 ± 16 events is expected in the region $65 \text{ GeV} < m_{bb} < 145 \text{ GeV}$



cut flow table

table of number of scaled events + significance of signal after each cut:

cuts	$\kappa_F = +1$	$\kappa_F = -1$	tbbj	tZbbj	$t\bar{t}$	$\Sigma(\text{bg})$	$\frac{S}{\sqrt{B}}$ (SM)	$\frac{S}{\sqrt{B}}$ (non-SM)
initial	35	592	240	134	4491105	4491479	0.02	0.28
1 lepton	18	325	180	67	2217352	2217599	0.01	0.22
1 fwd jet	7	136	58	24	258299	258381	0.01	0.27
≥ 3 central jets	6	122	44	20	193120	193184	0.01	0.28
≥ 2 b-jets	6	119	41	20	134949	135009	0.02	0.32
≥ 3 b-jets	4	88	25	13	2903	2940	0.08	1.63
m_{bb}	2	62	12	8	1254	1274	0.08	1.74



considering effects of systematic uncertainties

- What is the effect of systematic uncertainties on the significance value?

Summary of significance calculations in the SM and non-SM case, assuming systematic uncertainties of 0, 5, 10 and 20%

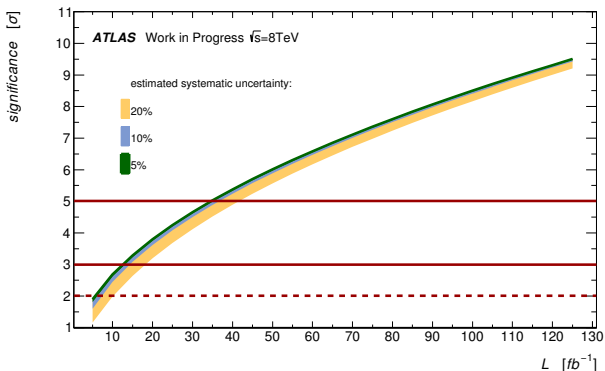
tH signal	S / \sqrt{B}	$S / \sqrt{B + (0.05B)^2}$	$S / \sqrt{B + (0.1B)^2}$	$S / \sqrt{B + (0.2B)^2}$
$c_F = +1$:	0.08	0.03	0.02	0.01
$c_F = -1$:	1.74	0.85	0.47	0.24



projection to 14 TeV energies

The situation is much improved when considering higher centre-of-mass energies + boosted cross-sections (RUN II)

Process	8 TeV [fb]	14 TeV [fb]
tH ($c_F = -1$)	234	982
$t\bar{t}(j)$	240×10^3	833×10^3



Conclusion

- At truth-level and 8 TeV, a sensitivity of $\sim 1.7\sigma$ is reached for anomalous coupling $c_F = -1$.
- This is significantly reduced when including estimates of systematic uncertainty → meaningful exclusion of the $c_F = -1$ cross-section unlikely.
- Results of the feasibility study were used to project the significance to 14 TeV, in anticipation of the LHC's upcoming run...
 - Cross-sections were rescaled to match new centre-of-mass energy.
 - Even with a systematic uncertainty of 20% in background, already at $\sim 20\text{fb}^{-1}$, degeneracy may be resolved.
 - Full-fledged analysis at 13/14 TeV with reconstructed data a worthwhile endeavour.





An update on tH measurements in $H \rightarrow b\bar{b}$ channel

Recently, CMS published⁶

Search for $H \rightarrow b\bar{b}$ in association with single top quarks as
a test of Higgs boson couplings

The CMS Collaboration

in which they drew a 95% CL upper limit at 7.62 times the rate expected for the
anomalous signal (expected sensitivity of 5.22).

⁶CMS PAS HIG-14-015



Potential extensions to analysis: lepton charge asymmetry

- **charge-asymmetric** initial state pp at the LHC → single top yields more ℓ^+ ⁷
- however, $t\bar{t}$ channel **symmetric** to LO ...@ NLO, t more centrally produced than \bar{t} → slightly more ℓ^- produced

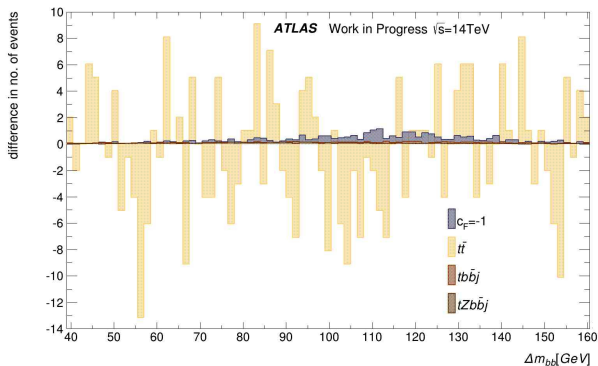
	frac. pos. leptons	frac. neg. leptons
sample $c_F = -1$	66%	33%
sample $t\bar{t}$	49.6%	50.4%

This asymmetry may prove useful: taking the difference in ℓ^+ & ℓ^- events means eliminating $t\bar{t}$ background, save for fluctuations.

⁷Matthew T. Bowen. Using charge asymmetries to measure single top quark production at the LHC. 558 Phys.Rev., D73:097501, 2006.



Potential extensions to analysis: lepton charge asymmetry



- Unfortunately, very little statistics (1 event or less per bin) remain.



THANK YOU!





systematic uncertainties

$$N = L \times \sigma \times \textit{acceptance} \times \textit{efficiency}$$

This analysis is rough: uncertainties are largely ignored as we are only looking into the feasibility of the work.

Uncertainties that will have to be considered for an extended analysis are uncertainties in

- jet energy scale
- jet energy resolution – broadens m_{bb} distribution
- lepton ID, reconstruction and trigger
- b, c and light jet tagging (increased 3 b-jet events in $t\bar{t}$ background, missed signal events...)
- background and signal modelling (PDF distributions, renormalisation, factorisation, calculations of cross-sections, mass measurements...)
- pile-up effects: affects jet energy scales and resolutions – more prominent for 14 TeV run.



Amplitude equation for tH process

The core process is $Wb \rightarrow tH$. The amplitude in the high-energy, hard-scattering regime for this is

$$\mathcal{A} = \frac{g}{\sqrt{2}} \left[(c_F - c_V) \frac{m_t \sqrt{s}}{m_W v} A \left(\frac{t}{s}, \varphi; \xi_t, \xi_b \right) + \left(c_V \frac{2m_W s}{v t} + (2c_F - c_V) \frac{m_t^2}{m_W v} \right) B \left(\frac{t}{s}, \varphi; \xi_t, \xi_b \right) \right]$$

We see that

- when $c_F = c_V$, the first term cancels
- when $c_F \neq c_V$, the term remains and grows $\propto \sqrt{s}$.



Significance formula

In order to get an estimate for the significance, do a simple event selection and count for each sample.

The p -value for an observation of $N = S + B$ events in a *background-only hypothesis* is

$$p = 1 - \Phi\left(\frac{N - B}{\sqrt{B}}\right) \quad (2)$$

where $\Phi\left(\frac{N - B}{\sqrt{B}}\right)$ is the cumulative distribution function for background-only.

The z -value is the significance,

$$\begin{aligned} Z &= \Phi^{-1}(1 - p) \\ &= \frac{N - B}{\sqrt{B}} \\ &= \frac{S}{\sqrt{B}}. \end{aligned}$$

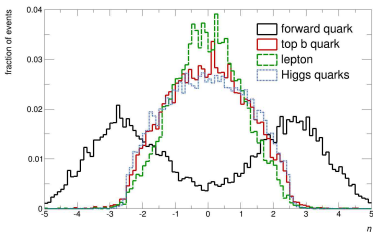
This holds for $S \ll B$.

As a general rule:

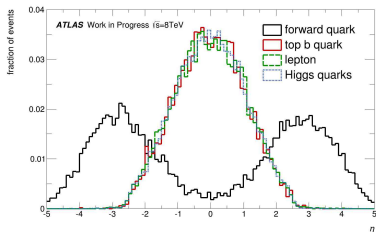
- $> 3\sigma$: “evidence”
- $> 5\sigma$: “observation”



η comparisons



(a) η distribution for $c_F = +1$

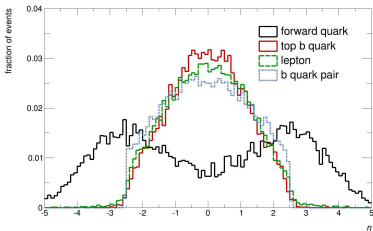


(b) η distribution for $c_F = -1$

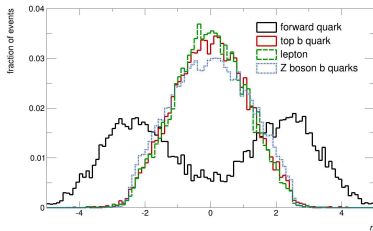
- wider η distribution for b quarks in $c_F = +1$ MC



η comparisons



(a) η distribution for bg tbbj

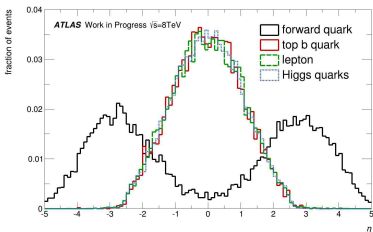


(b) η distribution for bg tZbbj

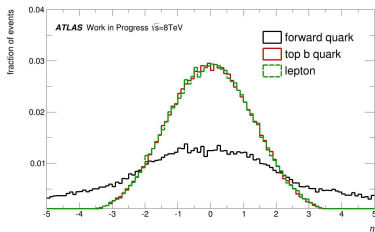
- backgrounds mimic signal



η comparisons



(a) η distribution for $C_F = -1$



(b) η distribution for $t\bar{t}$

- important characteristic of signal: non b-tagged jets at high η



η comparisons

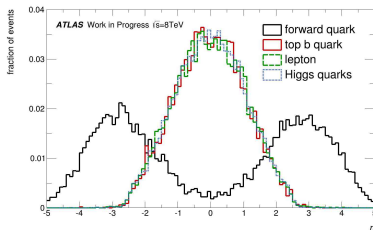


Figure : η distribution for $c_F = -1$

→ set requirement high p_T jet at $|\eta| > 2.5$

Sensitivities

- $c_F = +1$: 0.081
- $c_F = -1$: 1.612



Jets will have lower p_T values than the matched truth quarks.

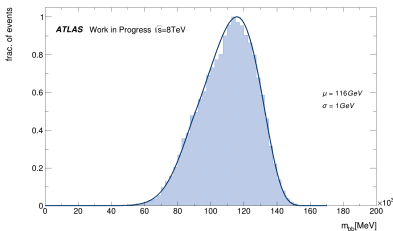


Figure shows the mean invariant mass of Higgs b jet pairs – jet p_T scale is shifted to a lower mean value which motivates a window of constraint on Higgs mass off-centre of the measured 125 GeV value.

$m_{bb} cut$

- $65 \text{ GeV} < m_{bb} < 140 \text{ GeV}$

Figure : reconstructed higgs using only jets matched to truth quarks, fitted with a Bukin function.

Jets used in reconstruction:

- Only jets matched to truth b quarks from Higgs decay
- basic $p_T > 25 \text{ GeV}$ and $\eta > 2.5$ cuts applied.



Reconstructing the W Boson

The neutrino z-component was calculated by using the W mass and four-momentum conservation,

$$p_z^\nu = \frac{1}{2(E_\ell^2 - p_z^{\ell^2})} \langle p_z^\ell [2(p_x^\ell E_T^{\text{miss}} \cos\phi + p_y^\ell E_T^{\text{miss}} \sin\phi) - m_\ell^2 + M_W^2] \rangle$$

$$\pm \sqrt{E_\ell^2 [(2p_x^\ell E_T^{\text{miss}} \cos\phi + 2p_y^\ell E_T^{\text{miss}} \sin\phi - m_\ell^2 + M_W^2)^2 - 4E_T^{\text{miss}2} (E_T^{\text{miss}2} - p_z^{\ell^2})]}$$

The minimum value for p_z^ν is selected. If the value within the square root is negative, E_T^{miss} is recalculated so that discriminant is 0,

$$E_T^{\text{miss}} = \frac{1}{2[E_\ell^2 - p_z^{\ell^2} - (p_x^\ell \cos\phi + p_y^\ell \sin\phi)^2]} \langle -(p_x^\ell \cos\phi + p_y^\ell \sin\phi) \times (m_\ell^2 - M_W^2) \pm \sqrt{(m_\ell^2 - M_W^2)^2} \rangle$$

(4)

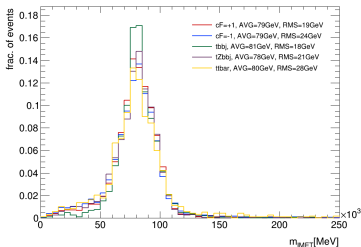
and choosing the value that results in a valid value for eq. 3.



Higgs reconstruction using top information

W boson is reconstructed using $\ell + MET \rightarrow$ top is reconstructed using W + best matched b jet \rightarrow Higgs is reconstructed using remaining b jets (if a 4 b jet event, pair is chosen with invariant mass closest to truth Higgs mass):

reconstructed W:



reconstructed top:

